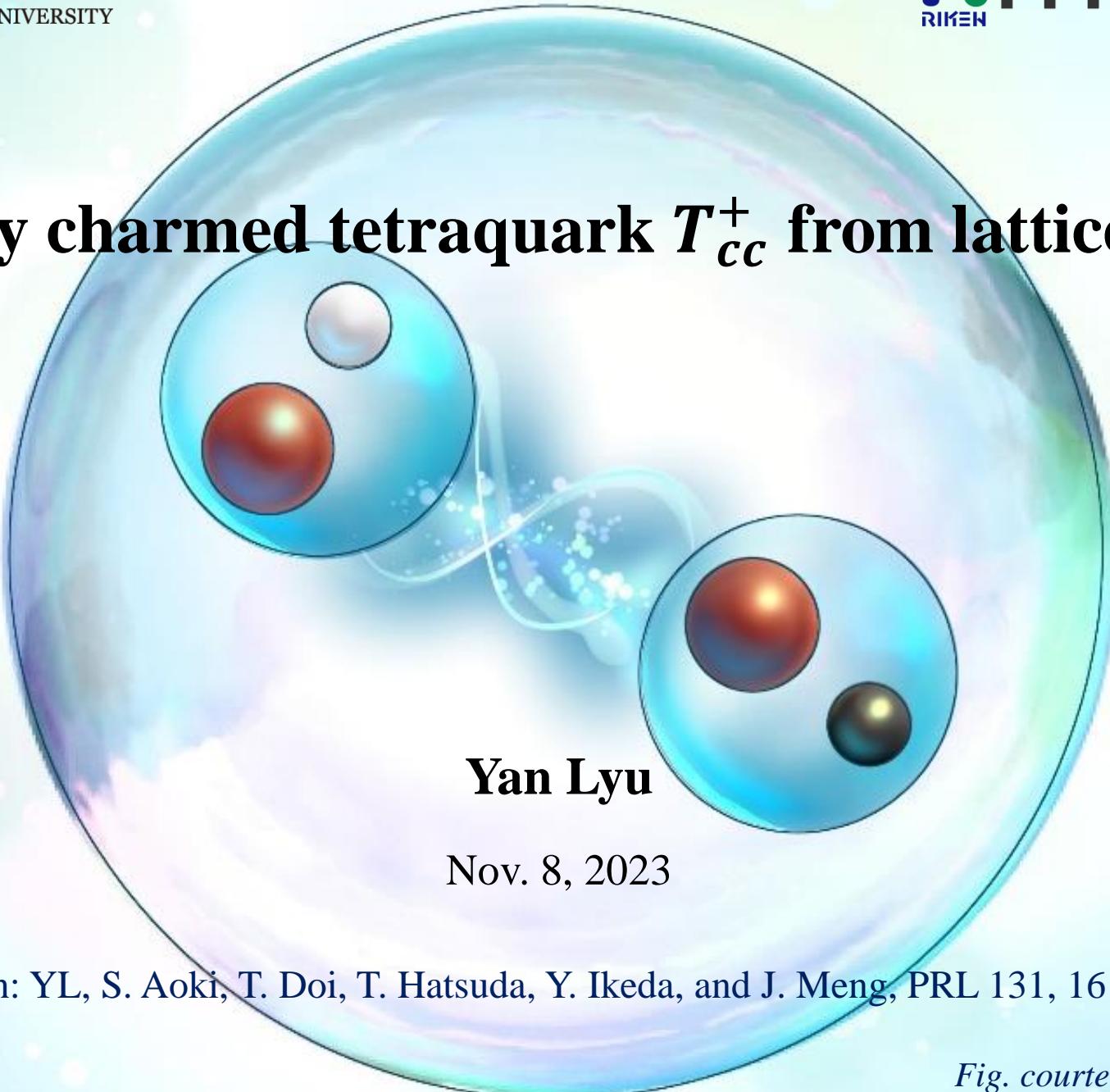




Doubly charmed tetraquark T_{cc}^+ from lattice QCD



A large, semi-transparent circular graphic serves as the background. Inside, two clusters of particles are shown. The left cluster consists of a red sphere and a smaller white sphere, both surrounded by a blue glow. A curved, light-blue trajectory with small white dots connects this cluster to a second cluster on the right. The right cluster features a red sphere and a black sphere, also within a blue glow. The overall effect is one of particles interacting or forming a bound state.

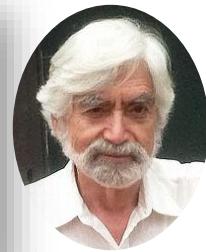
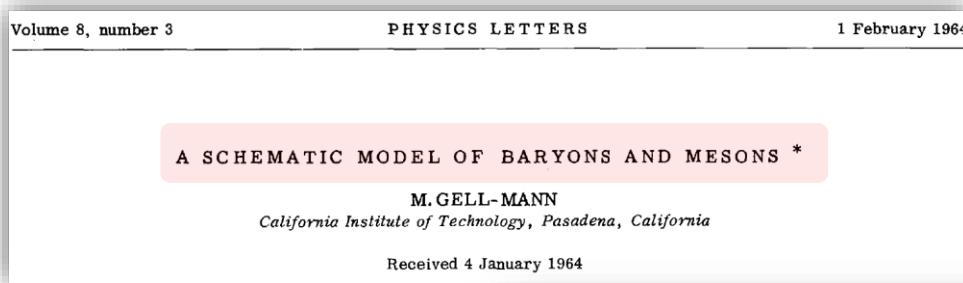
Yan Lyu

Nov. 8, 2023

Based on: YL, S. Aoki, T. Doi, T. Hatsuda, Y. Ikeda, and J. Meng, PRL 131, 161901 (2023)

Fig. courtesy of K. Murano

Conventional hadrons

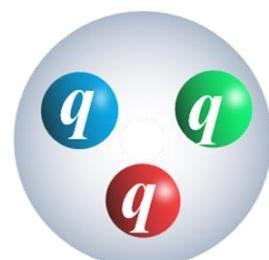


AN SU_3 MODEL FOR STRONG INTERACTION SYMMETRY AND ITS BREAKING

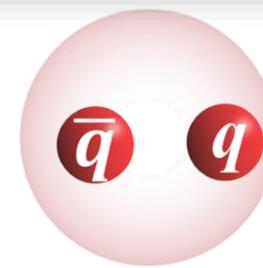
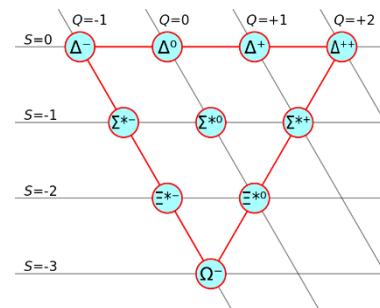
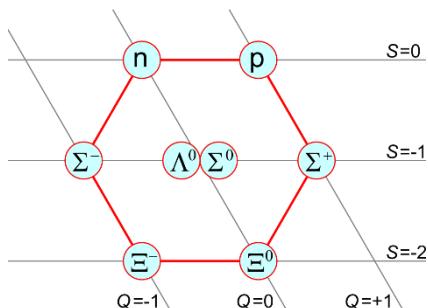
CERN LIBRARIES, GENEVA

G. Zweig *)

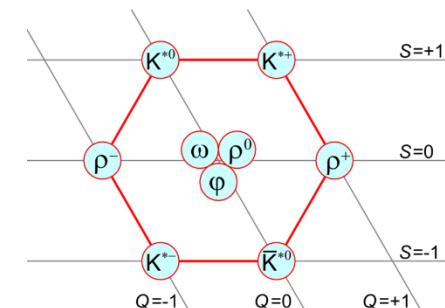
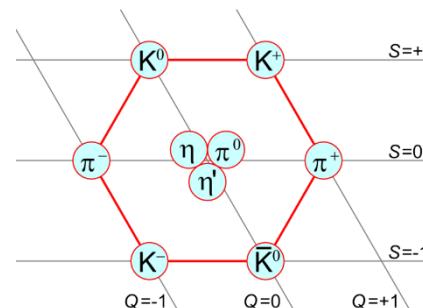
CERN - Geneva



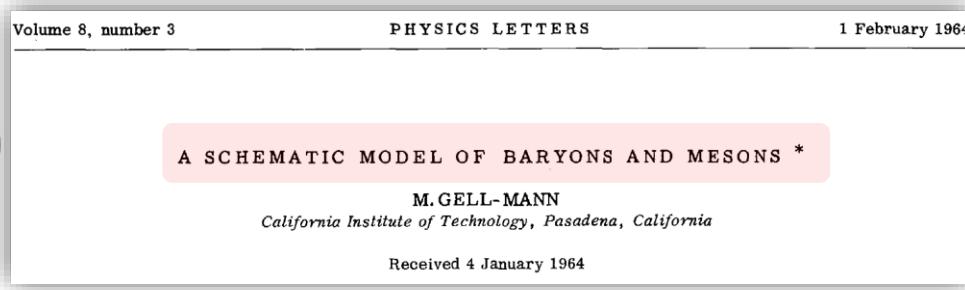
baryon



meson



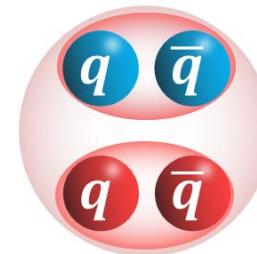
Exotic hadrons



“Baryons can now be constructed from quarks by using the combinations (qqq) , $(qqqq\bar{q})$, etc., while mesons are made out of $(q\bar{q})$, $(q\bar{q}\bar{q}\bar{q})$, etc.”



pentaquark



tetraquark

Hidden-charm exotic candidates

$P_c(4380)$

$P_c(4440)$

$P_c(4312)$

$P_c(4440)$

$P_c(4457)$

● 2003

$\chi_{c1}(3872)$

$\psi(4360)$
 $\psi(4660)$

$Z_c(4430)$
 $X(4050)$
 $X(4250)$

$\chi_{c1}(4140)$
 $\chi_{c1}(4274)$

$Z_c(3900)$

$X(4020)$
 $Z_c(4200)$
 $R_{c0}(4240)$

● 2013

● 2015

● 2020

$QQ\bar{q}\bar{q}'$ exotics

➤ Intriguing aspects on $QQ\bar{q}\bar{q}'$

- Open flavor, once observed its minimal quark content contains four quarks
- Likely to be bound in the limit of $m_Q \rightarrow \infty$

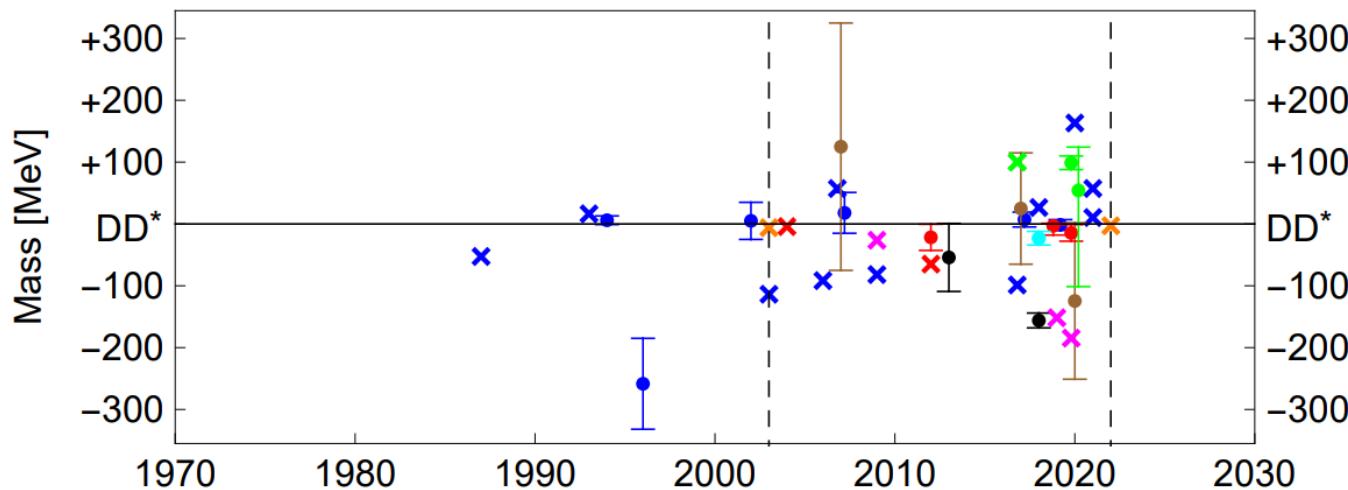
A. Manohar and M. Wise, Nucl. Phys. B 339, 17 (1993)

$bb\bar{q}\bar{q}'$ (↙)

$cc\bar{q}\bar{q}'$ (?)

$ss\bar{q}\bar{q}'$ (✗)

➤ A long history of theoretical prediction on $cc\bar{u}\bar{d}$ ($IJ^P = 01^+$)



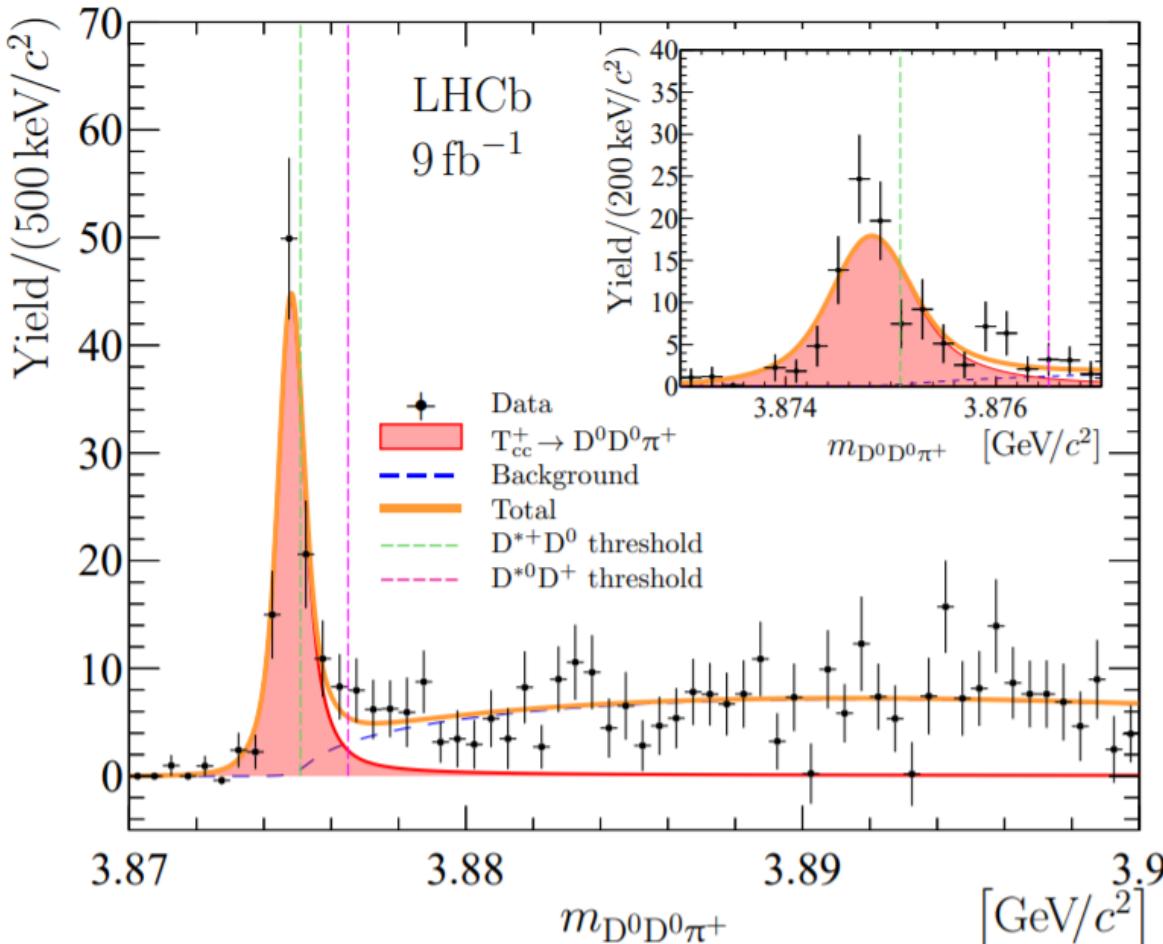
H.-X. Chen *et al.*, Rept. Prog. Phys. 86, 026201 (2023)

- Mass varies within ± 300 MeV with respect to $m_{D^*} + m_D$

First doubly charmed tetraquark T_{cc}^+

- 2022, LHCb discovered T_{cc}^+ in the $D^0 D^0 \pi^+$ spectrum

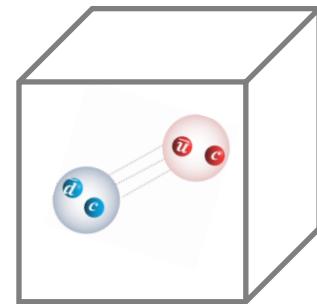
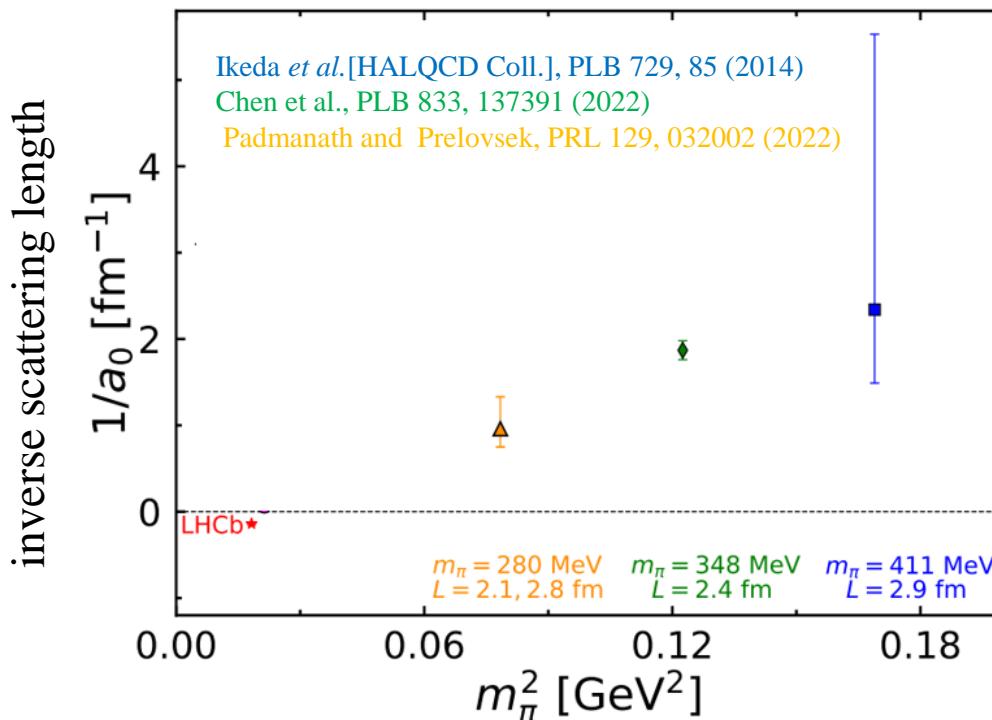
LHCb Coll., Nature Phys. 18, 751 (2022); Nature Comm. 13, 3351 (2022)



$I(J^P)$	δm_{pole}	Γ_{pole}	$\text{Re}(a_0)$	$\text{Im}(a_0)$
$0(1^+)$	$-360 \pm 40 \text{ keV}$	$48 \pm 2 \text{ keV}$	$7.16 \pm 0.51 \text{ fm}$	$-1.85 \pm 0.28 \text{ fm}$

T_{cc}^+ from first-principle lattice QCD

- Limited to heavy quark masses ($m_\pi \geq 280$ MeV)



- Purpose of this talk
 1. present the latest lattice results with (nearly) physical quark masses
 2. directly compare theoretical and experimental $DD\pi$ mass spectrum

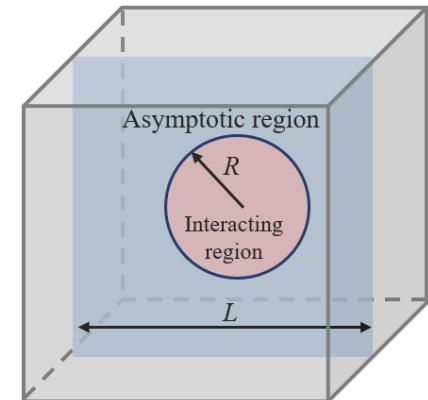
HAL QCD method

➤ Nambu-Bethe-Salpeter (NBS) amplitude

$$\psi^k(\mathbf{r})e^{-Et} = \langle 0 | \hat{D}^*(\mathbf{r}, t) \hat{D}(\mathbf{0}, t) | D^*(\mathbf{k}) D(-\mathbf{k}); E \rangle$$

- Asymptotic region: $\psi^k(r) \simeq A \frac{\sin(kr - l\pi/2 + \delta(\mathbf{k}))}{kr}$
- Interacting region: define potential

$$(\nabla^2 + k^2)\psi^k(\mathbf{r}) = 2\mu \int d\mathbf{r}' \mathbf{U}(\mathbf{r}, \mathbf{r}') \psi^k(\mathbf{r}')$$



➤ R correlator (superposition of NBS amplitudes)

$$R(\mathbf{r}, t) = \sum_n a_n \psi_n(\mathbf{r}) e^{-\Delta E_n t}$$

$$\left(\frac{1}{8\mu} \frac{\partial^2}{\partial t^2} - \frac{\partial}{\partial t} + \frac{\nabla^2}{2\mu} \right) R(\mathbf{r}, t) = \int d\mathbf{r}' \mathbf{U}(\mathbf{r}, \mathbf{r}') R(\mathbf{r}', t)$$

- Derivative expansion: $U(\mathbf{r}, \mathbf{r}') = \sum V_i(\mathbf{r}) \nabla^i \delta(\mathbf{r} - \mathbf{r}')$

$$V(r) = R(\mathbf{r}, t)^{-1} \left(\frac{1}{8\mu} \frac{\partial^2}{\partial t^2} - \frac{\partial}{\partial t} + \frac{\nabla^2}{2\mu} \right) R(\mathbf{r}, t)$$

N. Ishii, S. Aoki and T. Hatsuda, Phys. Rev. Lett. 99, 022001 (2007)
 N. Ishii, *et al.* [HAL QCD Coll.], Phys. Lett. B 712, 437 (2012)

Lattice setup

➤ (2+1)-flavor configuration

- Iwasaki gauge action
- $O(a)$ -improved Wilson quark action for uds quark
- Relativistic heavy quark action for c quark

K.-I. Ishikawa *et al.* [PACS Coll.], *Proc. Sci., LATTICE2015* 075 (2016)

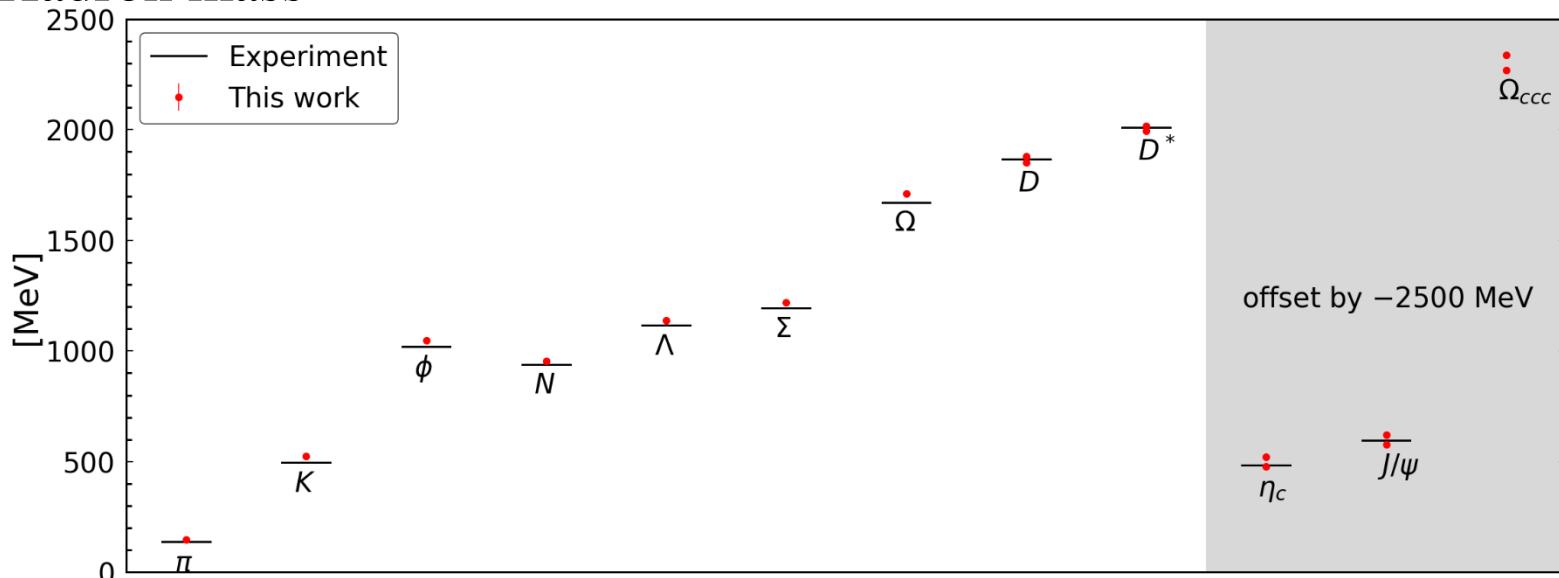
Y. Namekawa *et al.* [PACS Coll.], *Proc. Sci., LATTICE2016* 125 (2017)



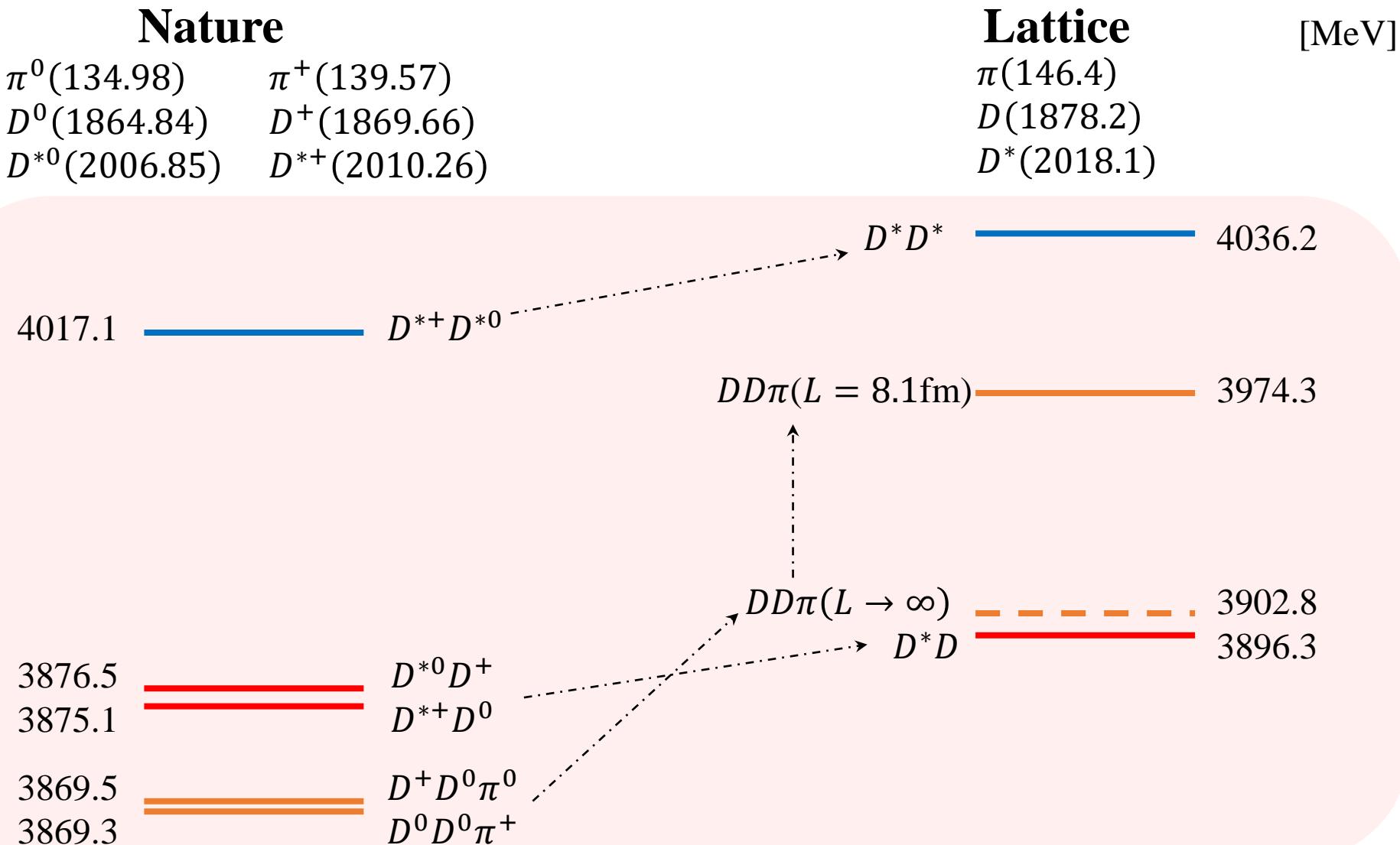
Fugaku supercomputer (440 PFlops)

$L^3 \times T$	a [fm]	La [fm]	m_π [MeV]	m_K [MeV]
$96^3 \times 96$	0.0846	8.1	146	525

➤ Hadron mass



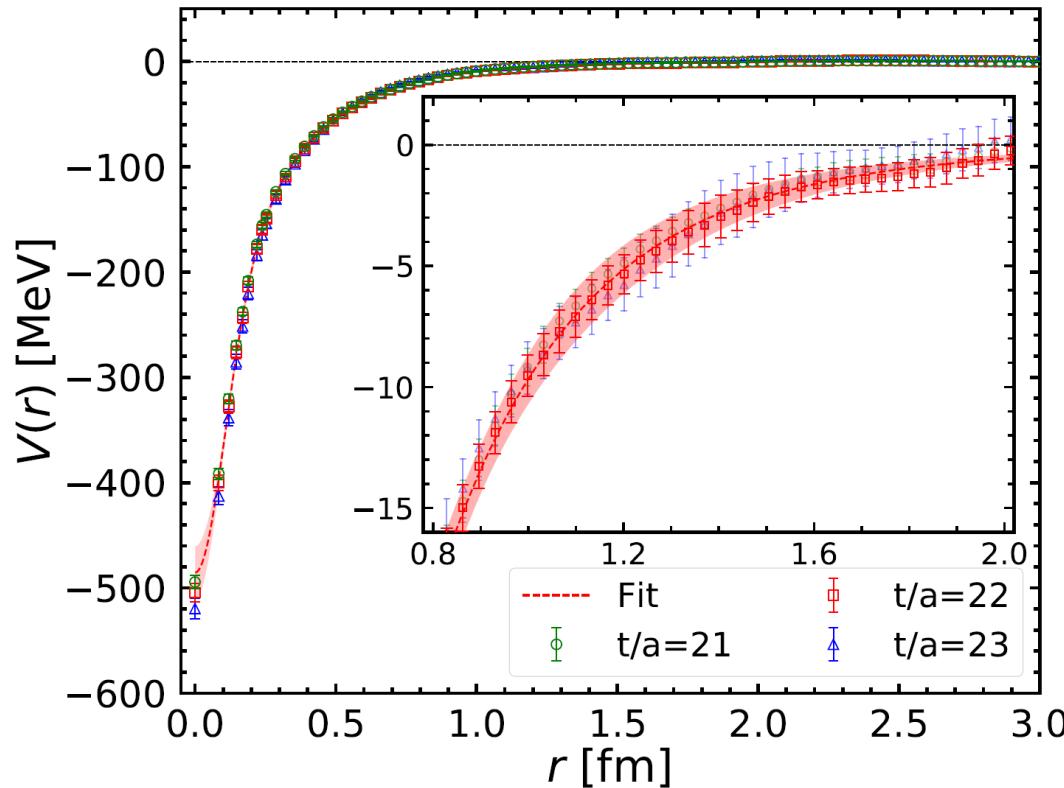
Energy levels



- The lowest energy level of $DD\pi$ (D^*D^*) is around 78 (140) MeV above on the lattice

D^*D interaction

- D^*D potential in the $(I,J^P) = (0,1^+)$ channel



- Short range: antiquark-diquark $[\bar{u}d]_{3_c, I=J=0} [cc]_{\bar{3}_c, J=1}$
R. Jaffe and F. Wilczek, Phys. Rev. Lett. 91 232003 (2003)
- Long range: attraction from pion-exchange interaction

Long-range potential

➤ One-pion exchange

S. Ohkoda *et. al.*, Phys. Rev. D 86, 034019 (2012)
N. Li, *et. al.*, Phys. Rev. D 88, 114008 (2013)

$$V(r) = -\alpha \frac{e^{-\mu r}}{r}, \quad \mu = m_\pi \text{ or } \sqrt{(m_{D^*} - m_D)^2 - m_\pi^2}$$

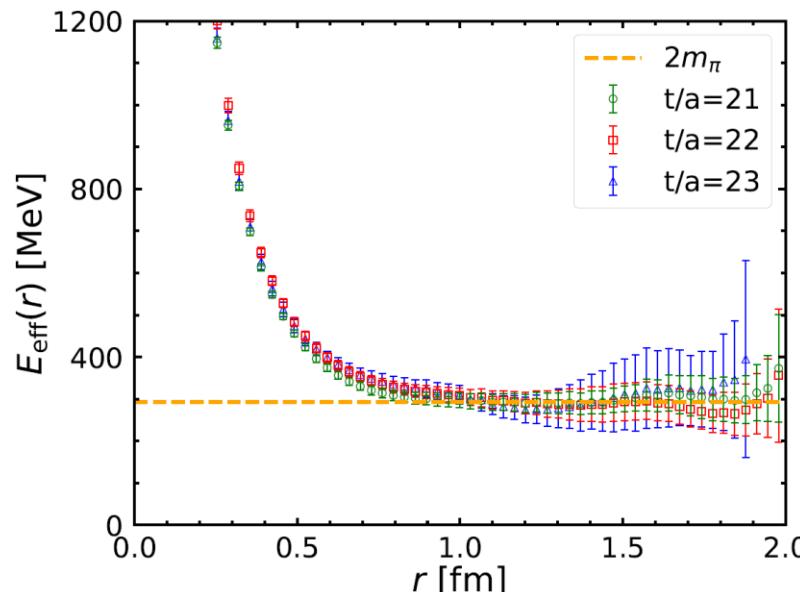
- Fail to describe long-range potential (why?)

➤ Two-pion exchange

$$V(r) = -\alpha \frac{e^{-2m_\pi r}}{r^2}$$



$$E_{\text{eff}}(r) = -\frac{\ln[-V(r)r^2/\alpha]}{r}$$



- Long range potential ($1 \leq r \leq 2$ fm) is consistent with two-pion exchange

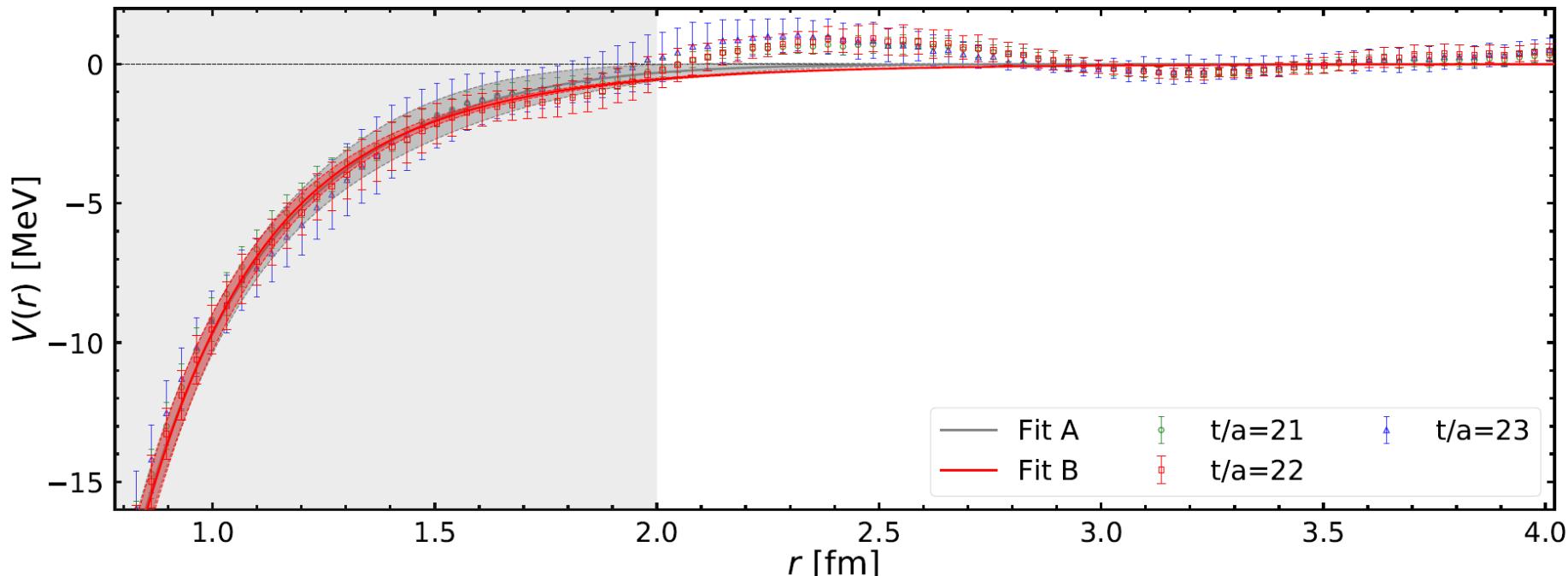
Fit

- Fit A: purely phenomenological fit ($\chi^2/\text{d.o.f.} = 1.01$)

$$V_{\text{fit}}(r) = \sum_{i=1, \dots, 4} a_i e^{-(r/b_i)^2}$$

- Fit B: TPE-motivated fit ($\chi^2/\text{d.o.f.} = 0.96$)

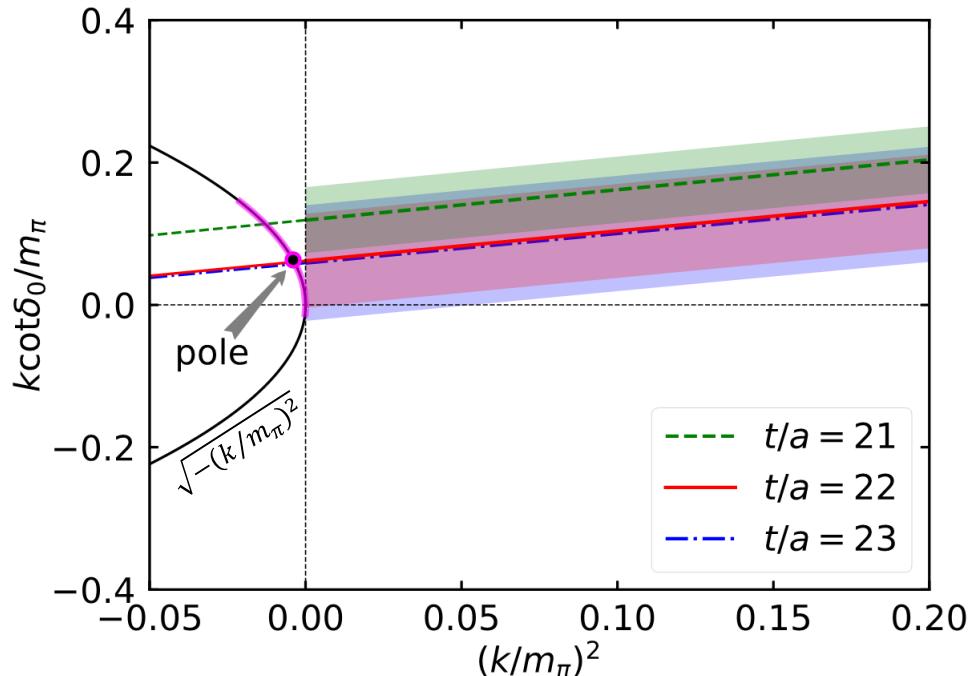
$$V_{\text{fit}}(r; m_\pi) = \sum_{i=1,2} a_i e^{-(r/b_i)^2} + a_3 (1 - e^{-(r/b_3)^2})^2 \frac{e^{-2m_\pi r}}{r^2}$$



Scattering properties

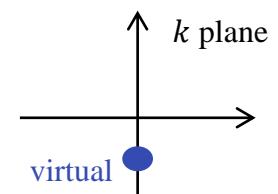
➤ Scattering phase shift

$$k \cot \delta_0 = \frac{1}{a_0} + \frac{1}{2} r_{\text{eff}} k^2$$
$$f^{-1}(k) \sim k \cot \delta_0 - ik$$



➤ Scattering parameters and pole singularities

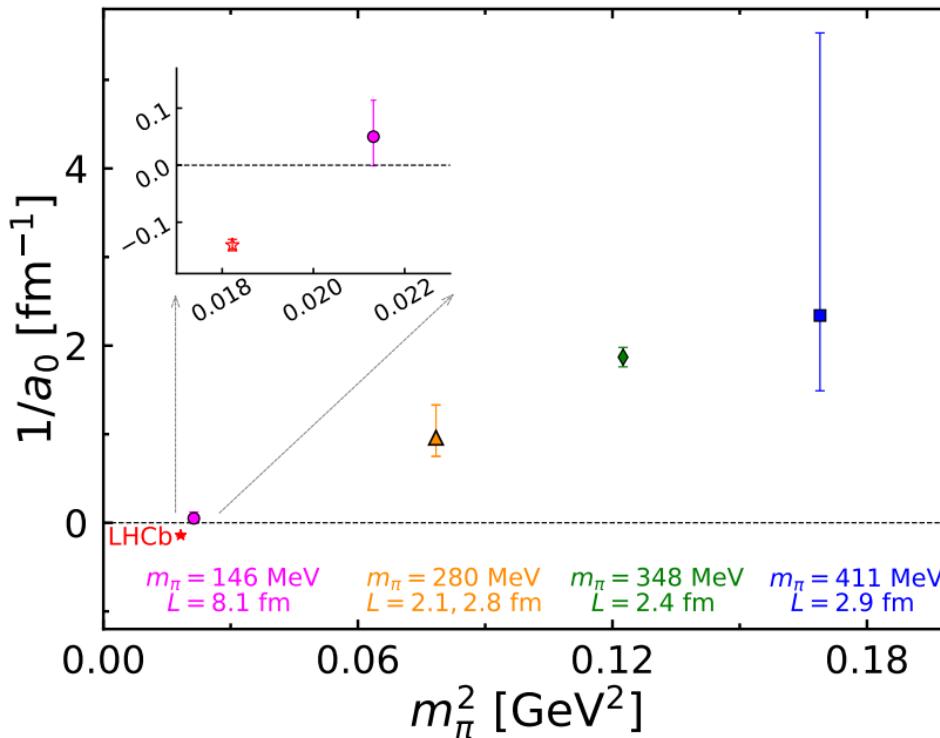
m_π (MeV)	146.4
$1/a_0$ (fm $^{-1}$)	$0.05(5)(^{+2}_{-2})$
r_{eff} (fm)	$1.12(3)(^{+3}_{-8})$
$k = ik_{\text{pole}}$ κ_{pole} (MeV)	$-8(8)(^{+3}_{-5})$
E_{pole} (keV)	$-59(^{+53}_{-99})(^{+2}_{-67})$



- T_{cc}^+ appears as a virtual state at $m_\pi = 146.4$ MeV

Comparison

➤ $1/a_0$



Ikeda *et al.* [HALQCD Coll.], Phys. Lett. B 729, 85 (2014)

Chen et al., Phys. Lett. B 833, 137391 (2022)

Padmanath and Prelovsek, Phys. Rev. Lett. 129, 032002 (2022)

- As m_π decreases, LQCD results approach to the experimental data

Extrapolate to physical point based on TPE

➤ Extrapolation

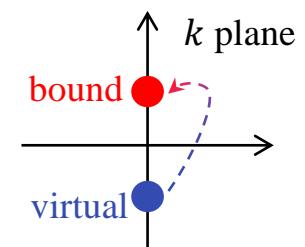
- Extrapolate TPE interaction to physical point

$$V_{\text{fit}}(r; m_\pi = 146 \rightarrow 135 \text{ MeV})$$

- Adopt physical values for $m_{D^{*+}}$ and m_{D^0}
- Do NOT consider isospin breaking nor opening of $DD\pi$ channel

➤ Scattering parameters and pole singularities

m_π (MeV)	146.4	135.0
$1/a_0$ (fm $^{-1}$)	0.05(5)($^{+2}_{-2}$)	-0.03(4)
r_{eff} (fm)	1.12(3)($^{+3}_{-8}$)	1.12(3)
$k = ik_{\text{pole}} \kappa_{\text{pole}}$ (MeV)	-8(8)($^{+3}_{-5}$)	+5(8)
E_{pole} (keV)	-59($^{+53}_{-99}$)($^{+2}_{-67}$)	-45($^{+41}_{-78}$)



- $m_\pi = 146 \rightarrow 135$ MeV, T_{cc}^+ evolves from a virtual state into a bound state

Extrapolation to physical point based on a_0

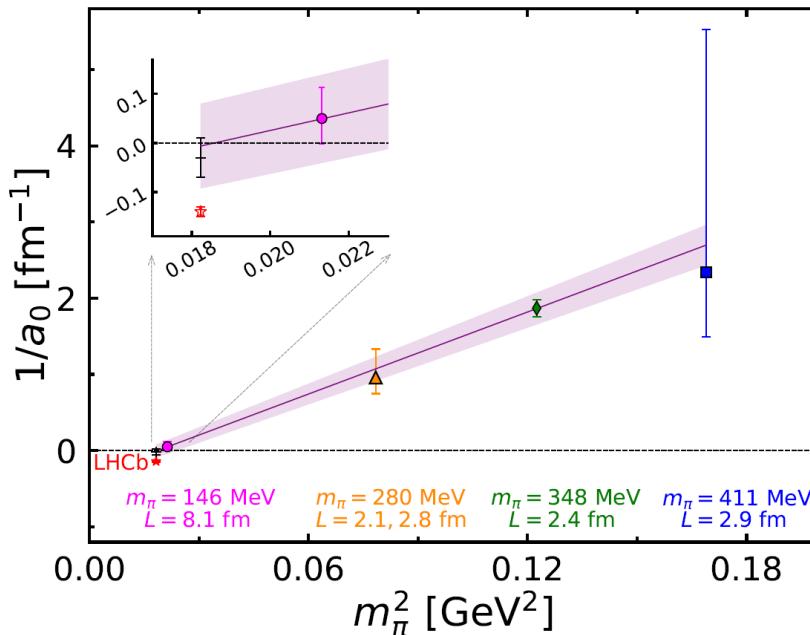
➤ Extrapolation

- A linear fit to four $1/a_0$ s from different m_π

$$1/a_0(m_\pi) = c + dm_\pi^2$$

- Four data from different calculations and posses different systematics

➤ $1/a_0$



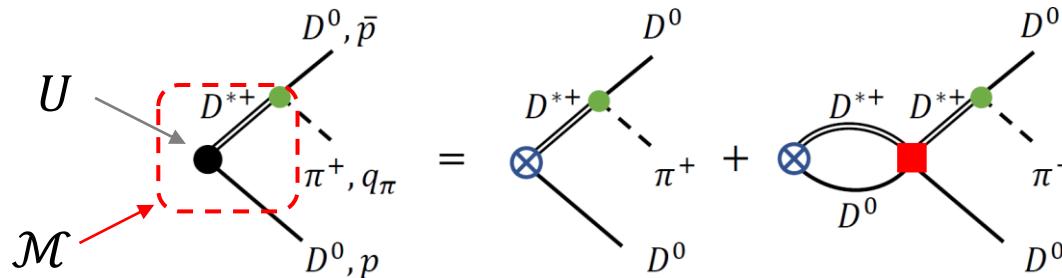
- $1/a_0$ from two extrapolations are consistent with each other

Construction of $D^0 D^0 \pi^+$ spectrum

- Production amplitude of $D^{*+}D^0$ from a source function P

$$U(M, p) = P + \int \frac{d^3 q}{(2\pi)^3} T(M, p, q) G(M, q) P$$

- For simplicity, consider a pointlike source (constant in p -space, $P = \mathcal{N}$)
- Only S -wave production at low energies



- Adopt experimental values for m_{D^{*+}, D^0, π^+} and $\Gamma_{D^{*+}}$ in the kinematics to keep the same phase space with the experiment
- Three-body mass spectrum for $D^0 D^0 \pi^+$

$$\mathcal{M}(U \rightarrow D^0 D^0 \pi^+) = U(M, p) G(M, p) q_\pi + U(M, \bar{p}) G(M, \bar{p}) \bar{q}_\pi$$

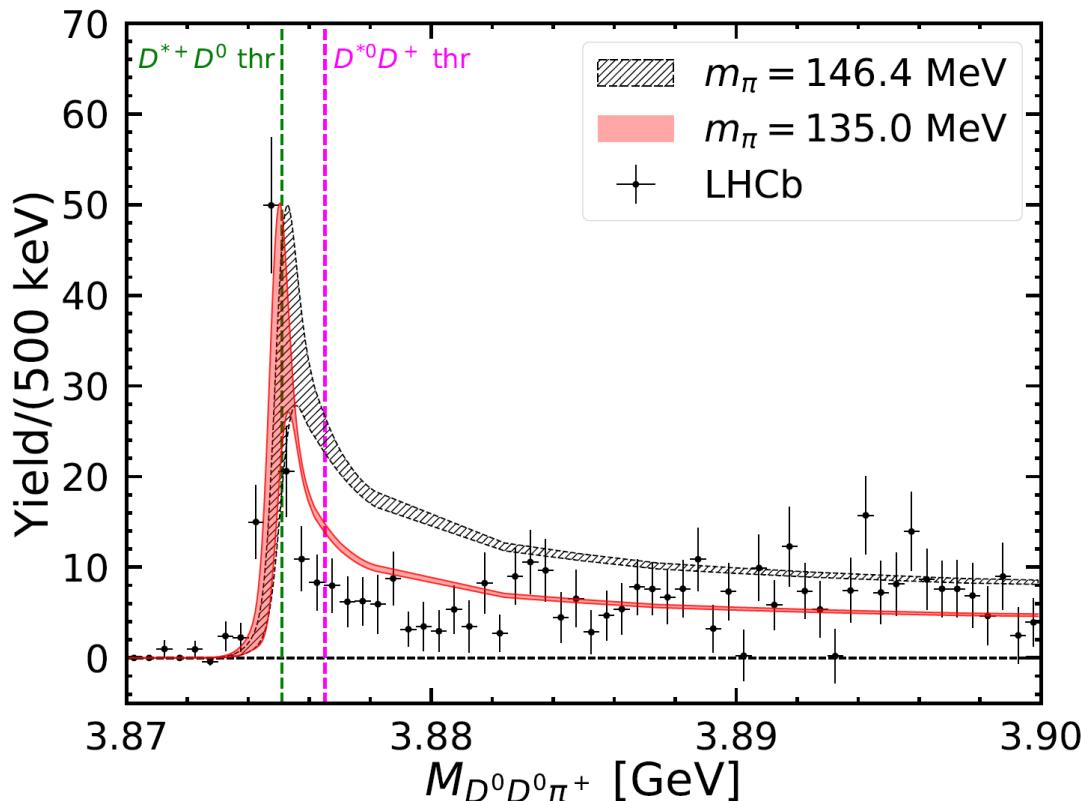
$$\frac{d\text{Br}}{dM} = \mathcal{N}' \int_0^{p_{\max}} pdp \int_{\bar{p}_{\min}}^{\bar{p}_{\max}} \bar{p} d\bar{p} |\mathcal{M}(U \rightarrow D^0 D^0 \pi^+)|^2$$

- A known energy resolution function needs to be considered for comparison w/ exp. data

LHCb Coll., Nature Comm. 13, 3351 (2022)

$D^0 D^0 \pi^+$ spectrum

➤ Results at different m_π

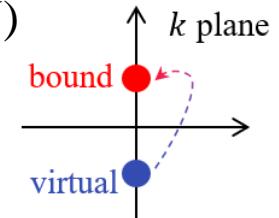


- A peak around $D^{*+} D^0$ threshold
- $m_\pi = 146$ MeV → 135 MeV, peak position shifts to the left, better description to LHCb data

Summary & Outlook

➤ Summary: present the scattering properties of D^*D system from LQCD with nearly physical $m_\pi=146$ MeV

- Attractive potential with two-pion exchange interaction at long distances
- The potential leads to a near-threshold virtual state
- Extrapolate the potential to physical point ($m_\pi = 146 \rightarrow 135$ MeV)
 - ✓ virtual state → bound state
 - ✓ better description to the $DD\pi$ spectrum of LHCb experiment



➤ Outlook

- Physical-point simulation
 - ✓ opening of three-body channel
 - ✓ isospin breaking effect (i.e., coupled channel calculation)
- Study one-pion exchange interaction and associated left-hand-cut from LQCD



Fig. courtesy of K. Murano