

# Analysis of $\Xi(1620)$ resonance with chiral unitary approach

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T.Nishibuchi and T.Hyodo, arXiv:2305.10753 [hep-ph]

# Result of Belle

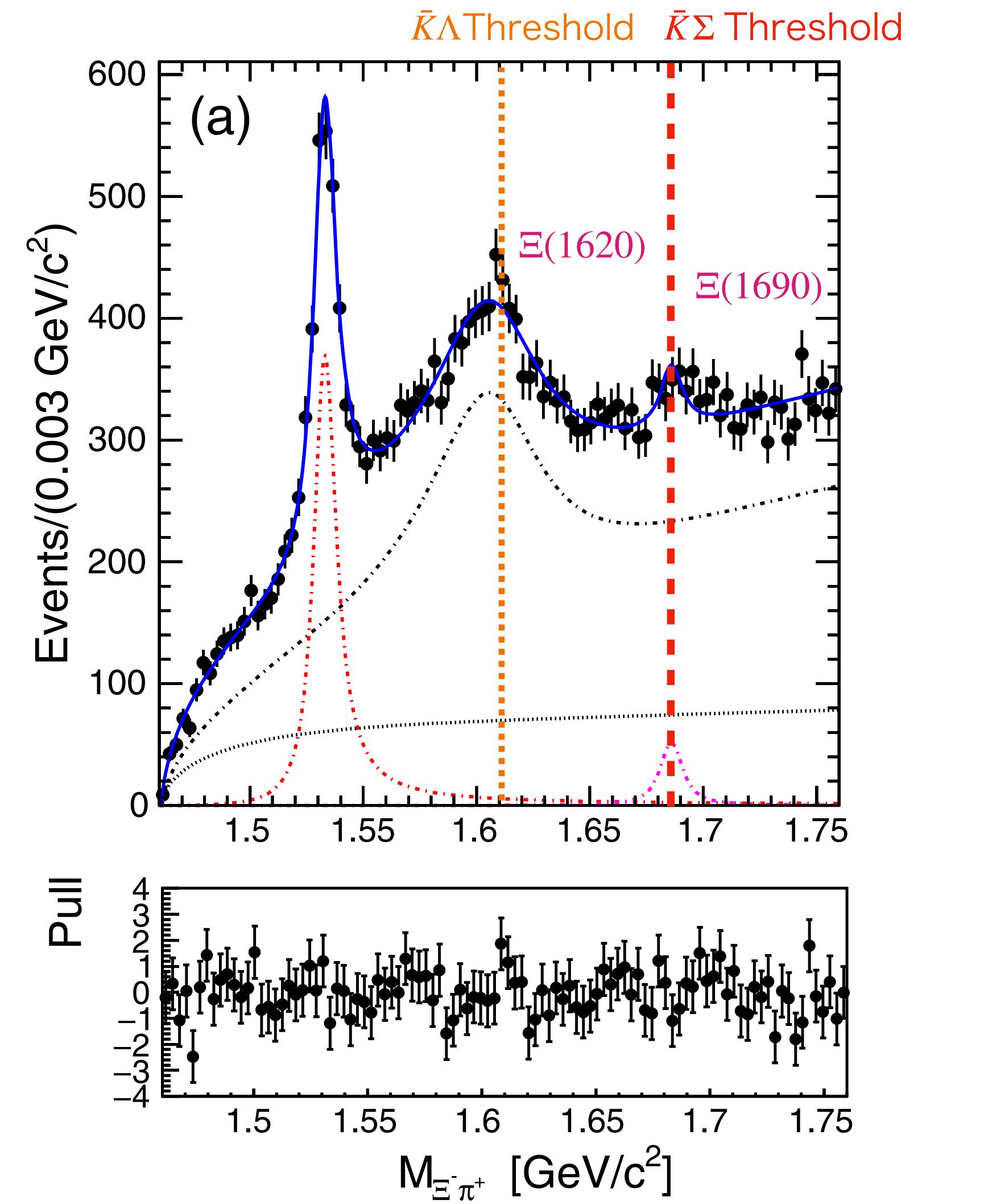
## Belle experiment of $\Xi_c \rightarrow \pi\pi\Xi$ (2019)[1]

- $\Xi(1620)$  and  $\Xi(1690)$  peaks are observed in the invariant mass distribution of  $\pi^+\Xi^-$ .
- The mass  $M_R$  and width  $\Gamma_R$  of  $\Xi(1620)$  are reported as follows.

$$M_R = 1610.4 \pm 6.0(\text{stat.})^{+6.1}_{-4.2}(\text{syst.}) \text{ MeV}$$

$$\Gamma_R = 59.9 \pm 4.8(\text{stat.})^{+2.8}_{-7.1}(\text{syst.}) \text{ MeV}$$

- Mass and width do not agree with values in previous theoretical studies.
- Peaks are close to thresholds of  $\bar{K}\Lambda$  and  $\bar{K}\Sigma$ ?



Invariant mass distribution of  $\pi^+\Xi^-$  in the  $\Xi_c \rightarrow \pi\pi\Xi$  decay [1].

[1]Belle collaboration, M. Sumihama et al., Phys. Rev. Lett. **122**, 072501 (2019).

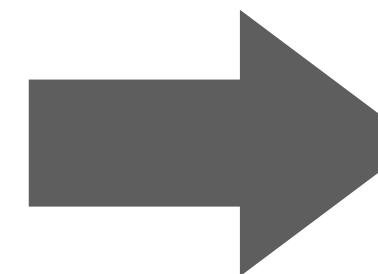
# Result of ALICE

## ALICE experiment(2021)[2]

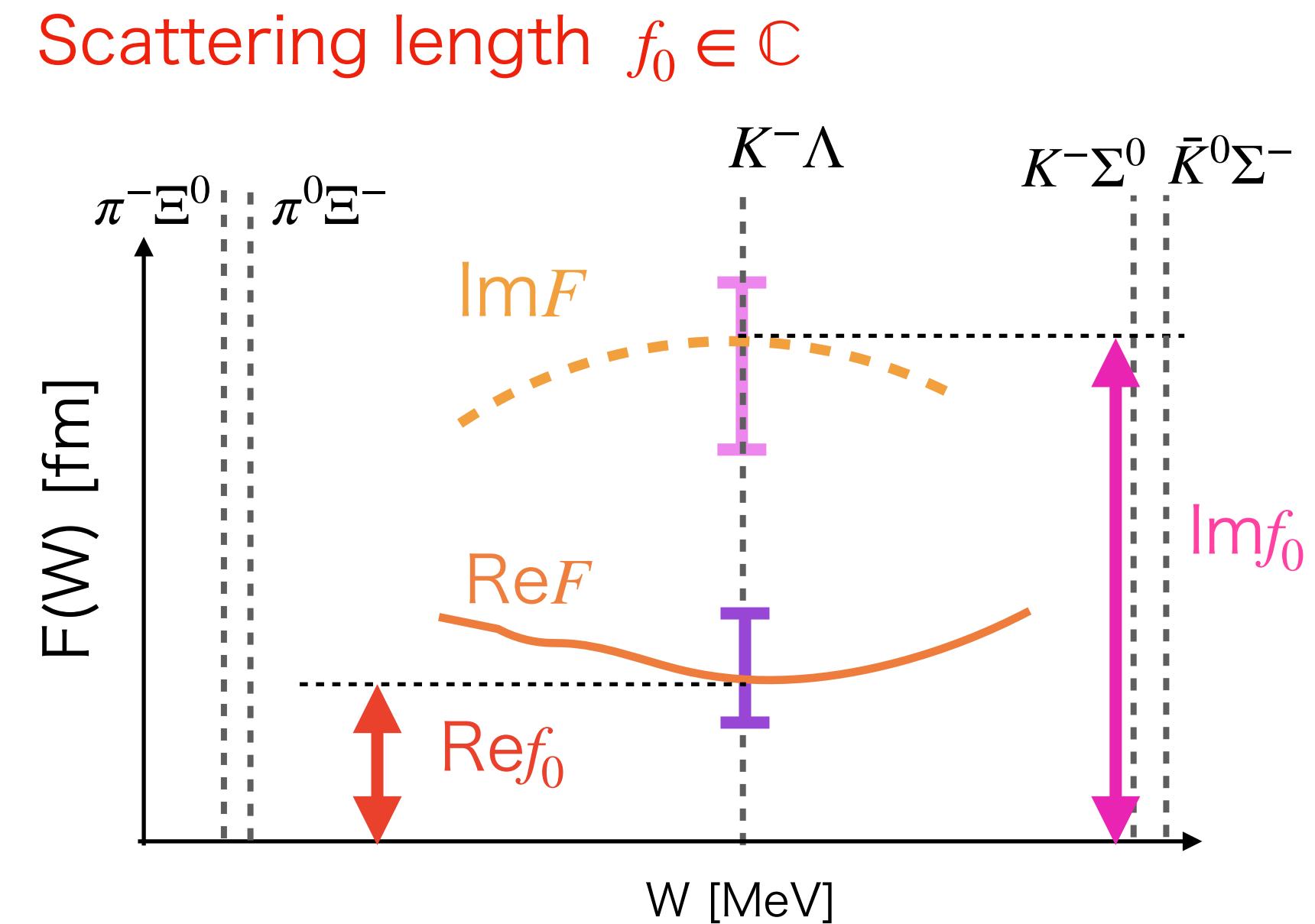
The scattering length  $f_0$  of  $K^-\Lambda$  was determined with femtoscopy in Pb-Pb collisions as follows.

$$\text{Re}f_0 = 0.27 \pm 0.12(\text{stat.}) \pm 0.07(\text{syst.}) \text{ fm}$$
$$\text{Im}f_0 = 0.40 \pm 0.11(\text{stat.}) \pm 0.07(\text{syst.}) \text{ fm}$$

$f_0$  constrains real and imaginary parts of scattering amplitude at threshold.



More quantitative comparison is possible than the spectrum fit.



Schematic figure of the scattering length

The error bar of  $\text{Re}f_0$

The error bar of  $\text{Im}f_0$

# New studies for $\Xi$ excited states

LHCb Collaboration(2021)[6]

- $\Xi^-(1690)$  and  $\Xi^-(1820)$  are observed in  $\Xi_b^- \rightarrow J/\psi \Lambda K^-$  decay.
- Mass  $M_R$  and width  $\Gamma_R$  of  $\Xi^-(1690)$  are reported as follows.

$$M_R = 1692.0 \pm 1.3(\text{stat.})^{+1.2}_{-0.4}(\text{syst.}) \text{ MeV}$$

$$\Gamma_R = 25.9 \pm 9.5(\text{stat.})^{+14.0}_{-13.5}(\text{syst.}) \text{ MeV}$$

**New theoretical analysis of  $\Xi(1620)$  and  $\Xi(1690)$  (2023)[7]**

The study based on chiral unitary approach with the Born and NLO terms.

It can generate dynamically both  $\Xi(1620)$  and  $\Xi(1690)$  with  $a_i$  in the range

$[-3.5, -1]$ .

$$\Xi(1620) \quad M_R = 1599.95 \text{ MeV}, \Gamma_R = 158.88 \text{ MeV.}$$

$$\Xi(1690) \quad M_R = 1683.04 \text{ MeV}, \Gamma_R = 11.51 \text{ MeV.}$$

$$M_R = 1608.51 \text{ MeV}, \Gamma_R = 170.00 \text{ MeV.}$$

$$M_R = 1686.17 \text{ MeV}, \Gamma_R = 29.72 \text{ MeV.}$$

[6]R. Aaij, et al., Sci. Bull. 66 (2021) 1278–1287. [7]Feijoo, A., Valcarce, V. and Magas, V. K., Phys.Lett.B 841 (2023) 137927.

# Strategy of this study

## Strategy of model construction

- Model of previous study[3] Set 1

Mass  $M_R$  and width  $\Gamma_R$  of  $\Xi(1620)$

$$M_R = 1607 \text{ MeV}, \Gamma_R = 280 \text{ MeV}.$$

We construct theoretical models which reproduce new experimental results.

- Model for Belle result (Model 1)

Assume pole position from mass  $M_R$  and width  $\Gamma_R$  reported by Belle

- Model for ALICE result (Model 2)

Reproduce the  $K^- \Lambda$  scattering length determined by ALICE

[3] A.Ramos, E.Oset and C.Bernhold, Phys. Rev. Lett. **89**.252001 (2002).

# Formulation

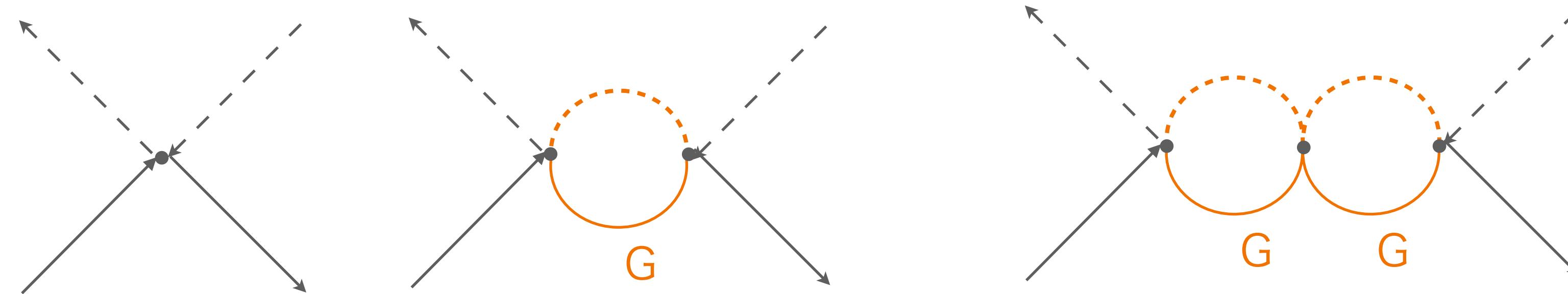
Coupled-channel meson-baryon scattering amplitude  $T_{ij}(W)$  at total energy  $W$ .

Scattering equation

$$T_{ij}(W) = V_{ij}(W) + V_{ik}(W)G_k(W)T_{kj}(W)$$

$V_{ij}(W)$ …Interaction kernel  
 $G_i(W)$ …Loop function

$$T_{ij}(W) = V_{ij}(W) + V_{ik}(W)G_k(W)V_{kj}(W) + V_{ik}(W)G_k(W)V_{kl}(W)G_l(W)V_{lj}(W) + \dots$$



Meson-baryon  
multiple scattering

The solution of the equation is obtained as

$$T_{ij}(W) = [[V(W)]^{-1} - G(W)]_{ij}^{-1}$$

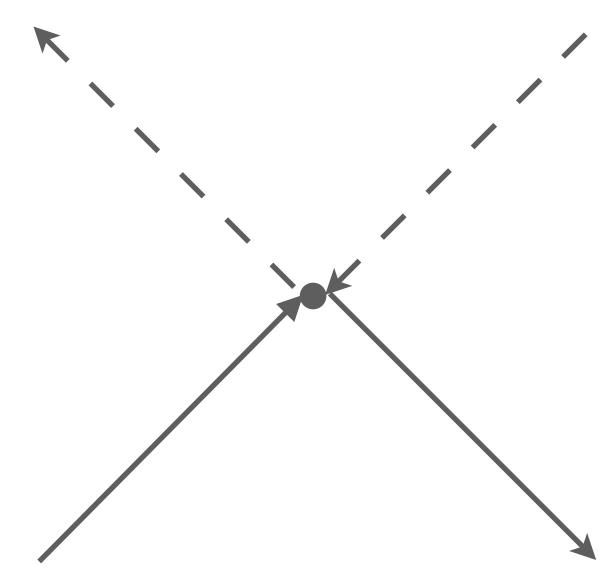
# Formulation

$V_{ij}(W)$ …Interaction kernel (Weinberg-Tomozawa term)  
s-wave interaction satisfying chiral low energy theorem.

$$V_{ij}(W) = -\frac{C_{ij}}{4f_i f_j} N_i N_j (2W - M_i - M_j)$$

$f_i$  : Meson decay constant,  $C_{ij}$  : Group theoretical coefficient,

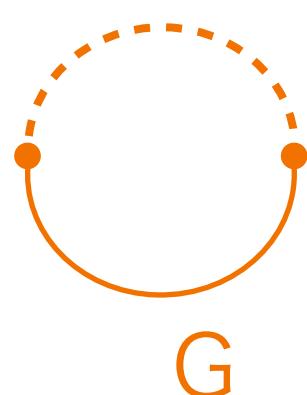
$M_i$  : Baryon Mass,  $N_i$  : Kinematical coefficient



$G_i(W, a_i)$ …Loop function  
(Divergence renormalized by dimensional regularization)

$$G_i(W) \rightarrow G_i(W, a_i)$$

$W$  : Total energy,  $a_i$  : Subtraction constant



# Model 1 construction

- Belle result :  $M_R = 1610 \text{ MeV}$ ,  $\Gamma_R = 60 \text{ MeV}$

- Based on the peak position, we define

$$z_{\text{ex}} = [1610 - 30i] \text{ MeV}.$$

- $z_{\text{th}}$  : Pole in theoretical model

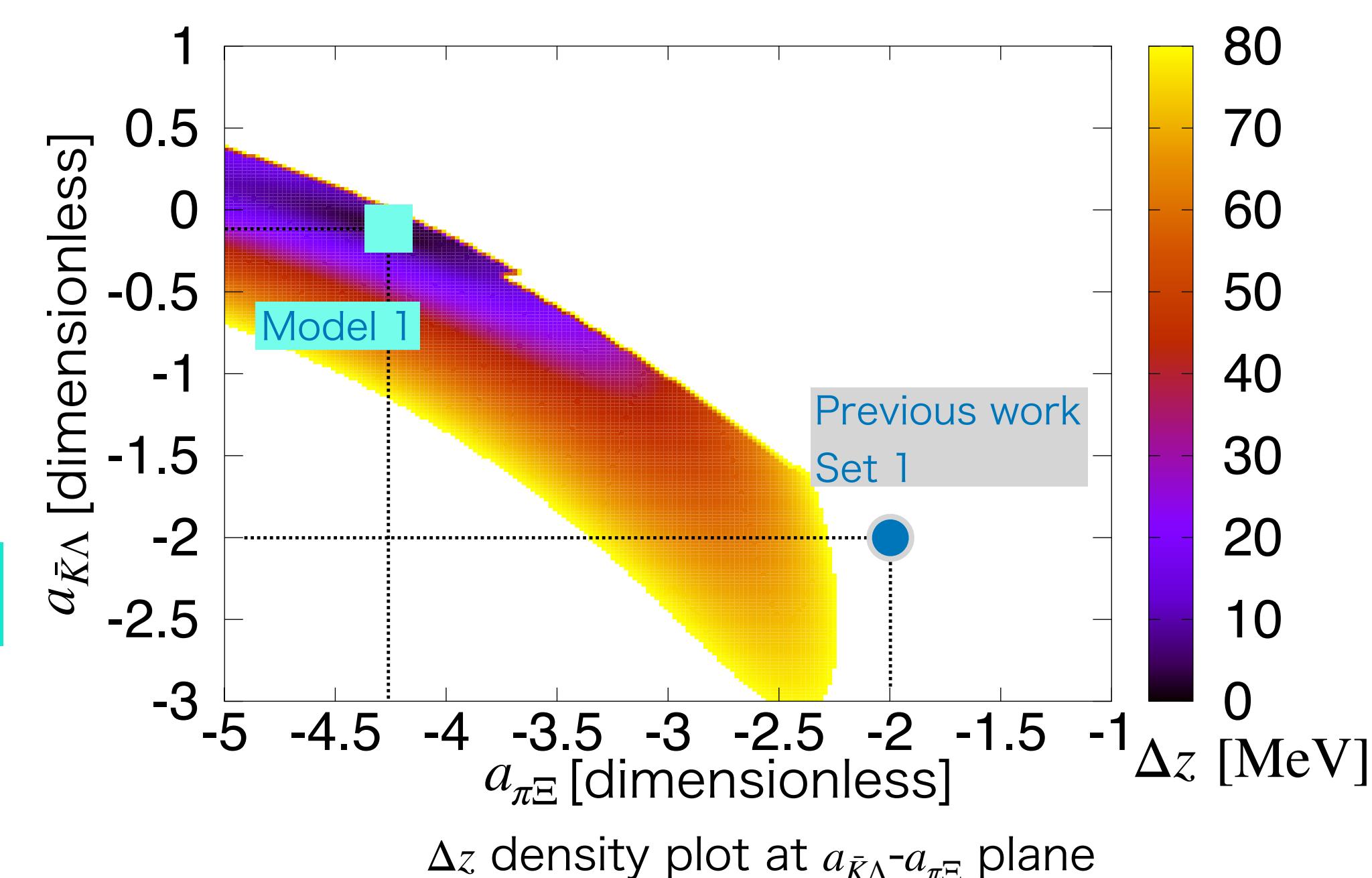
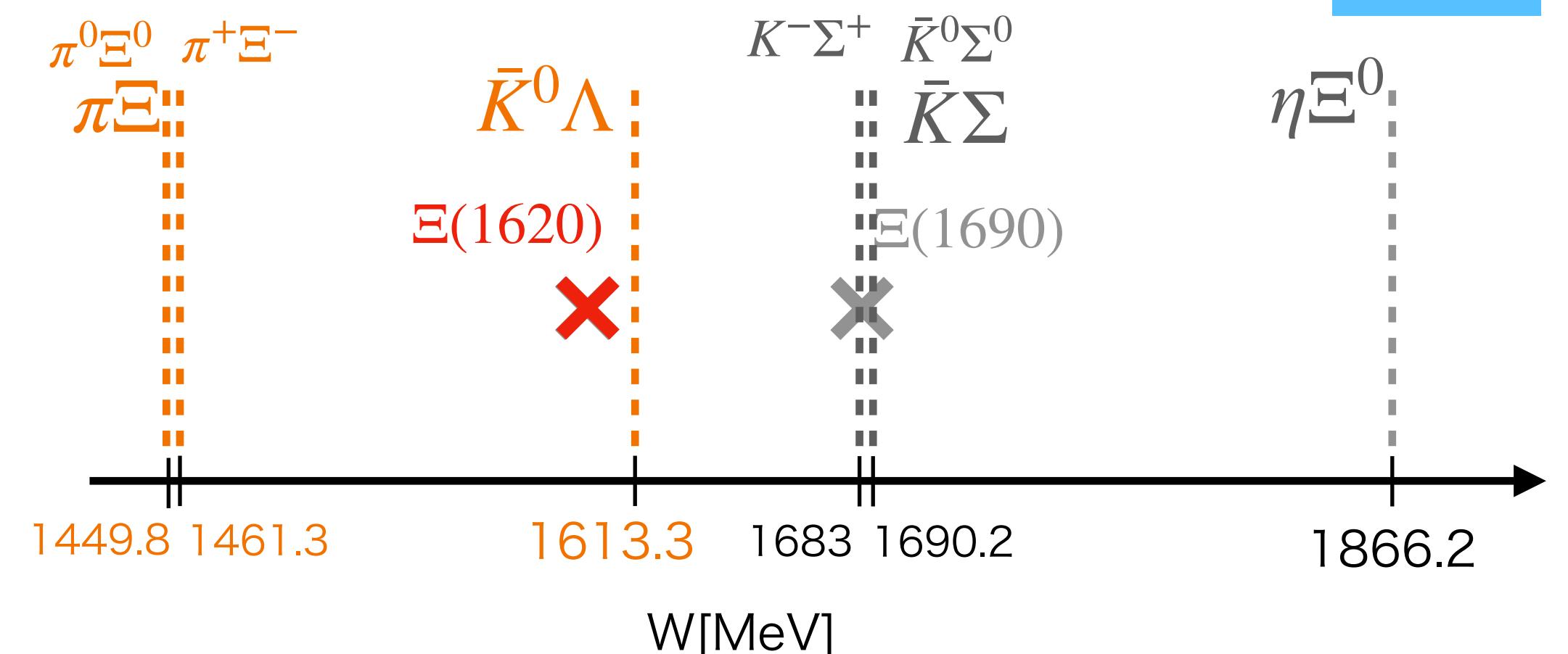
$$\Delta z = |z_{\text{th}} - z_{\text{ex}}|$$

- We minimize  $\Delta z$  by adjusting subtraction constants  $a_{\pi\Xi}$  and  $a_{\bar{K}\Lambda}$ .

$\Delta z \leq 1 \text{ MeV}$  is achieved

at  $a_{\pi\Xi} = -4.26$  and  $a_{\bar{K}\Lambda} = -0.12$ [4].

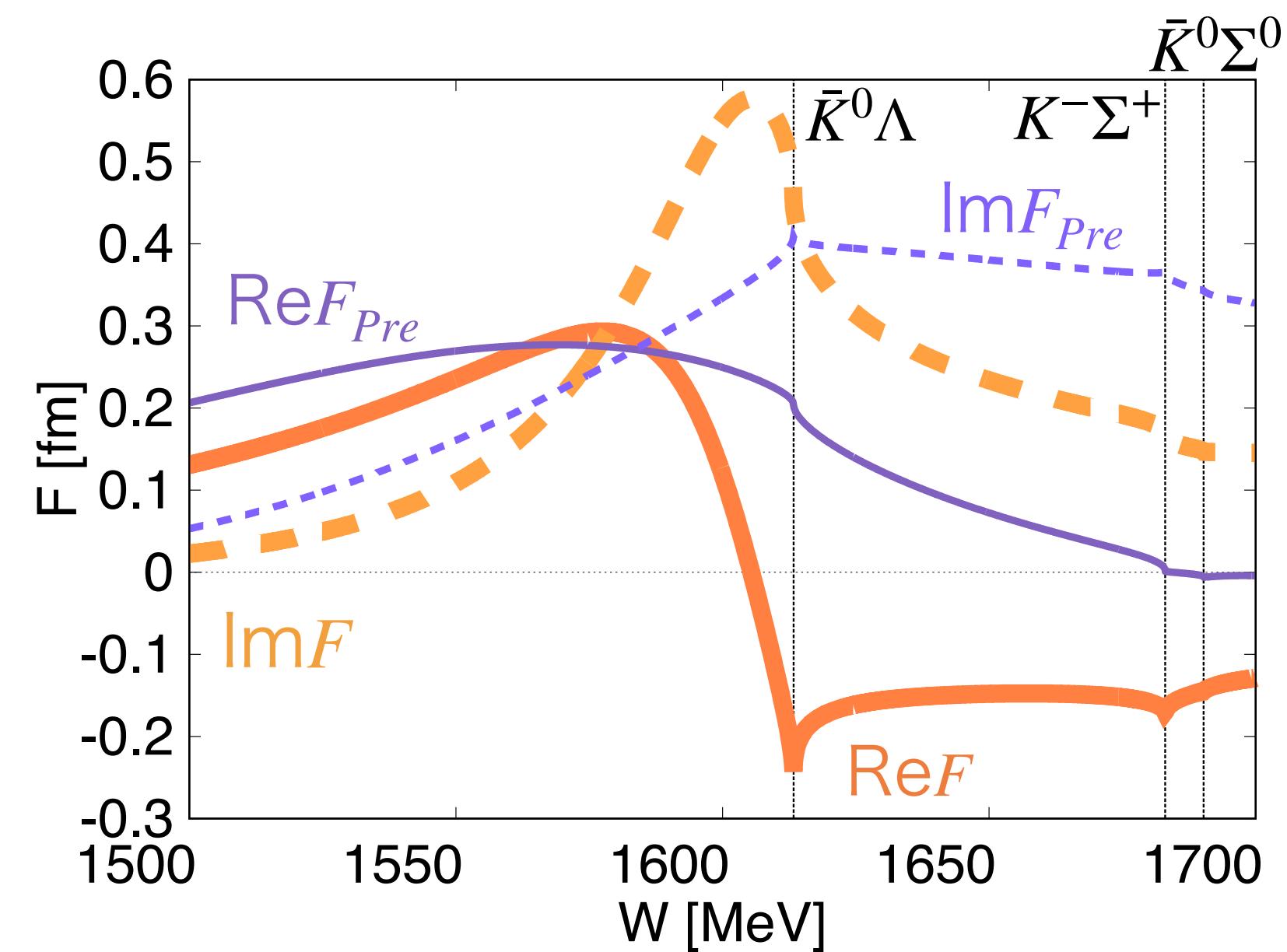
Named as **Model 1**



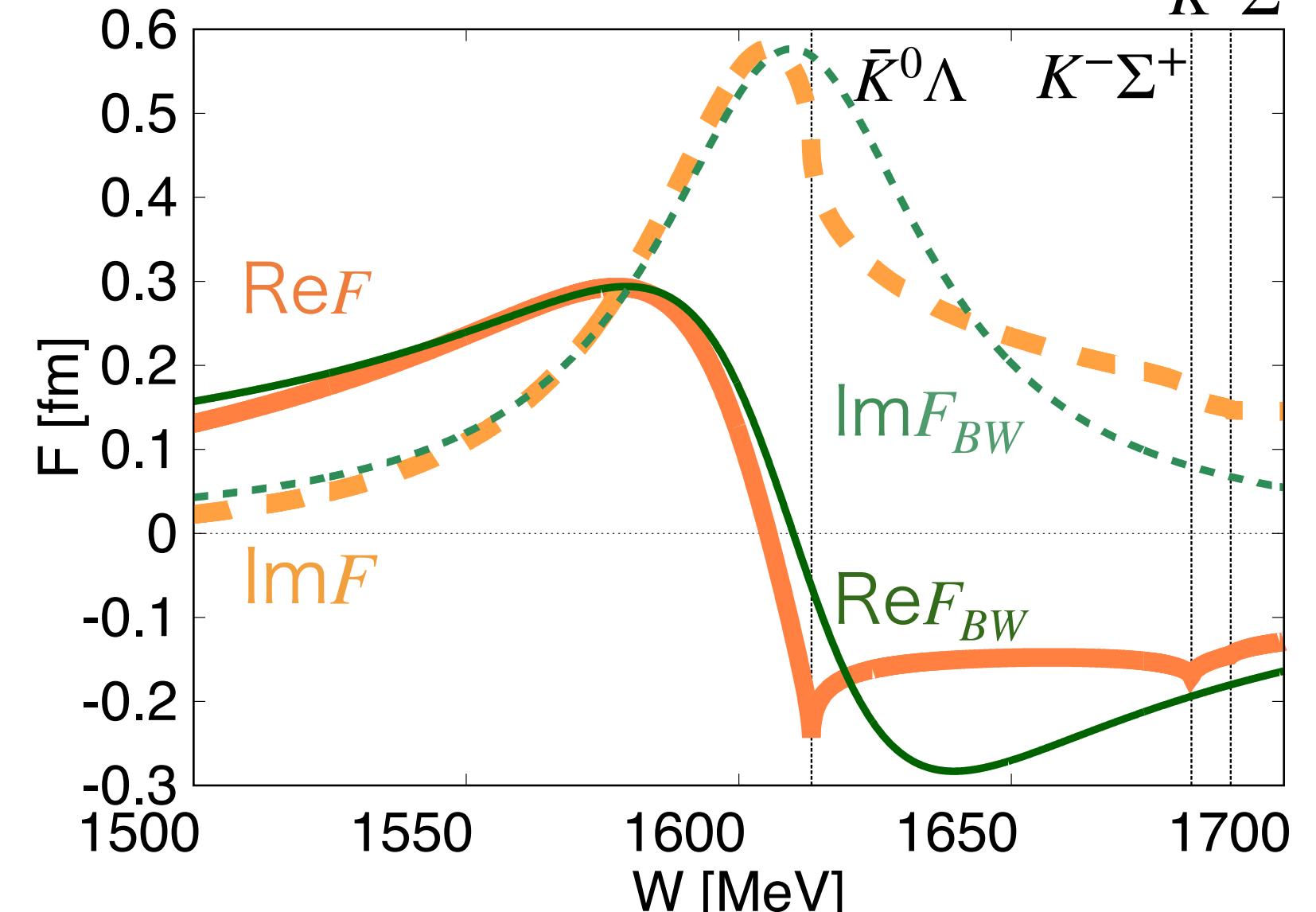
[4] T.Nishibuchi and T.Hyodo, EPJ Web of Conferences **271**, 10002 (2022); arXiv:2305.10753 [hep-ph]

# Result of Model 1

- $\pi^+\Xi^-$ -scattering amplitude of Model 1(Thin lines)
  - In comparison with previous study[3], distinct peak appears on real axis like Belle result.
  - In comparison with Breit-Wigner distribution, the peak position is shifted and the shape is distorted by the threshold effect.



▲comparison with previous study



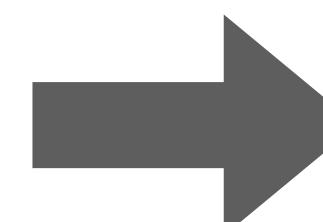
▲comparison with Breit-Wigner distribution

[3] A.Ramos, E.Oset and C.Bernhold Phys. Rev. Lett. **89**.252001 (2002).

# Comparison with previous study

Let compare this result  
with that of previous study

- . In previous study,  $a_i$  is changed in the range  $[-3.5, -1]$ .
- . The  $\Xi(1620)$  pole in previous study has broad decay width.



In this study, the  $\Xi(1620)$  pole with narrow decay width is reproduced with  $a_i$  out of the range  $[-3.5, -1]$ .

	This study $V = V_{WT}$	Previous study $V = V_{WT} + V_{Born} + V_{NLO}$	Model I	Model II
$a_{\pi\Xi}$	-4.26	-2.7981	-2.7228	
$a_{\bar{K}\Lambda}$	-0.12	-1.0071	-1.0000	
$a_{\bar{K}\Sigma}$	-2.00	-3.0938	-2.9381	
$a_{\eta\Xi}$	-2.00	-3.2665	-3.3984	
$\Xi(1620) M$ [MeV]	1610	1599.95	1608.51	
$\Gamma$ [MeV]	60	158.88	170.00	
$\Xi(1690) M$ [MeV]	-	1683.04	1686.17	
$\Gamma$ [MeV]	-	11.51	29.72	

Table1.Comparison the result of this study [5] with that of previous study [7].

[5]T.Nishibuchi and T.Hyodo, arXiv:2305.10753 [hep-ph]

[7]Feijoo, A., Valcarce, V. and Magas, V. K., Phys.Lett.B **841** (2023) 137927.

# Model 2 construction

- $f_0 \cdots K^- \Lambda$  scattering length

ALICE experiment:  $f_{\text{ALICE}} = 0.27 + 0.40i$  fm

Previous work(Set 1):  $f_0 = -0.07 + 0.21i$  fm

Model 1:  $f_0 = -0.75 + 0.93i$  fm

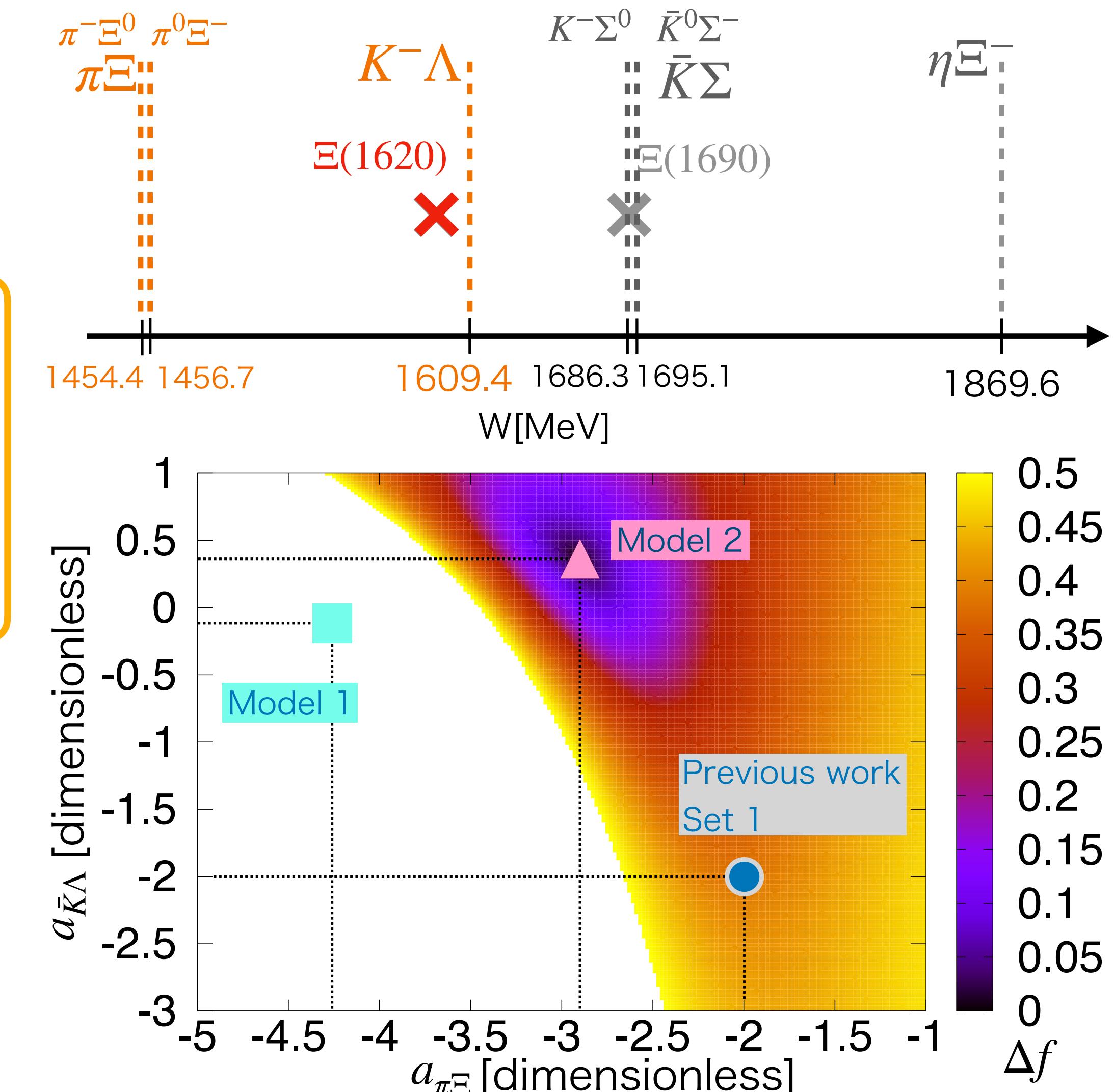
→ They do not reproduce  $f_{\text{ALICE}}$ .

- $f_{\text{th}}$ : scattering length in theoretical model

$$\Delta f = |f_{\text{th}} - f_{\text{ALICE}}|$$

- We minimize  $\Delta f$  by adjusting subtraction constants  $a_{\pi \Xi}$  and  $a_{\bar{K}\Lambda}$ .

$\Delta f = 0.00 + 0.00i$  fm is achieved at  $a_{\pi \Xi} = -2.90$  and  $a_{\bar{K}\Lambda} = 0.36$

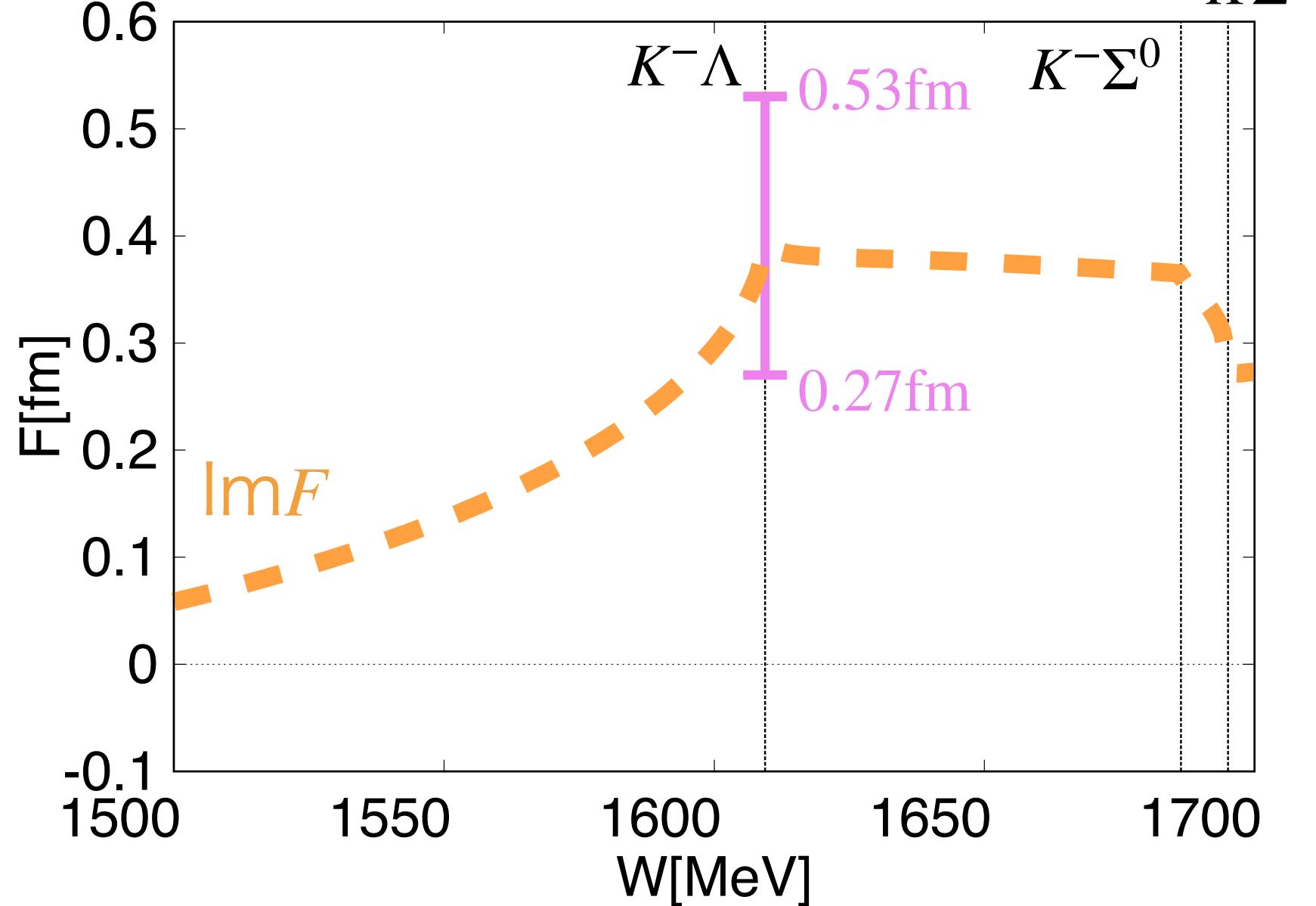
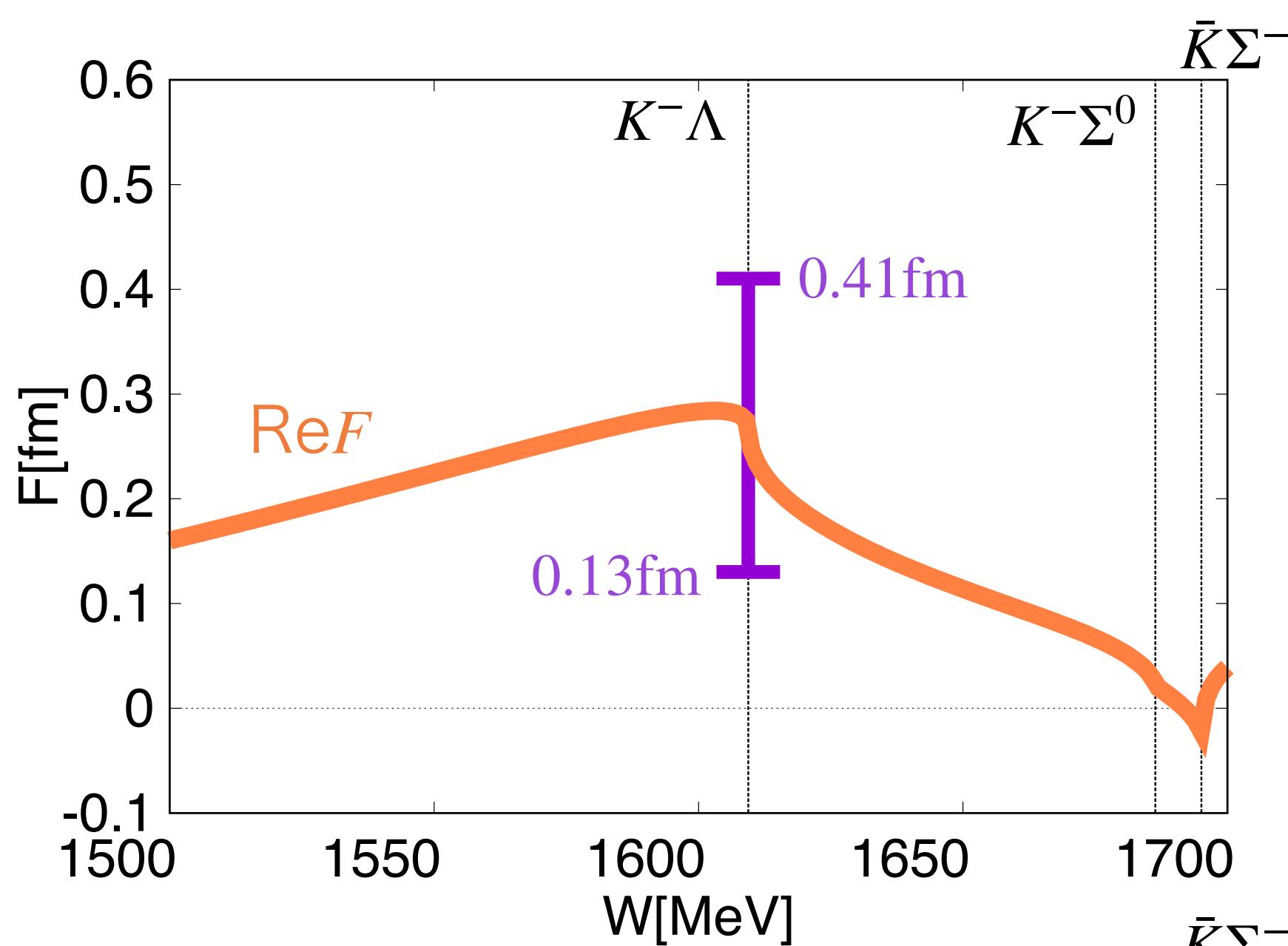


# Result of Model 2

- We plot the scattering amplitude with[5]  
 $a_{\pi\Sigma} = -2.90$ ,  $a_{\bar{K}\Lambda} = 0.36$  and  $f_{\text{th}} = 0.27 + 0.40i \text{ fm}$ .
- There are no peaks in the spectrum, but a cusp at the threshold.
- There are no poles on the physically relevant Riemann sheets.
- The pole locates at  $1725 + 76i$  [ttttt].

 The error bar of real part of  $f_{\text{ALICE}}$

 The error bar of imaginary part of  $f_{\text{ALICE}}$



[5]T.Nishibuchi and T.Hyodo, arXiv:2305.10753 [hep-ph]

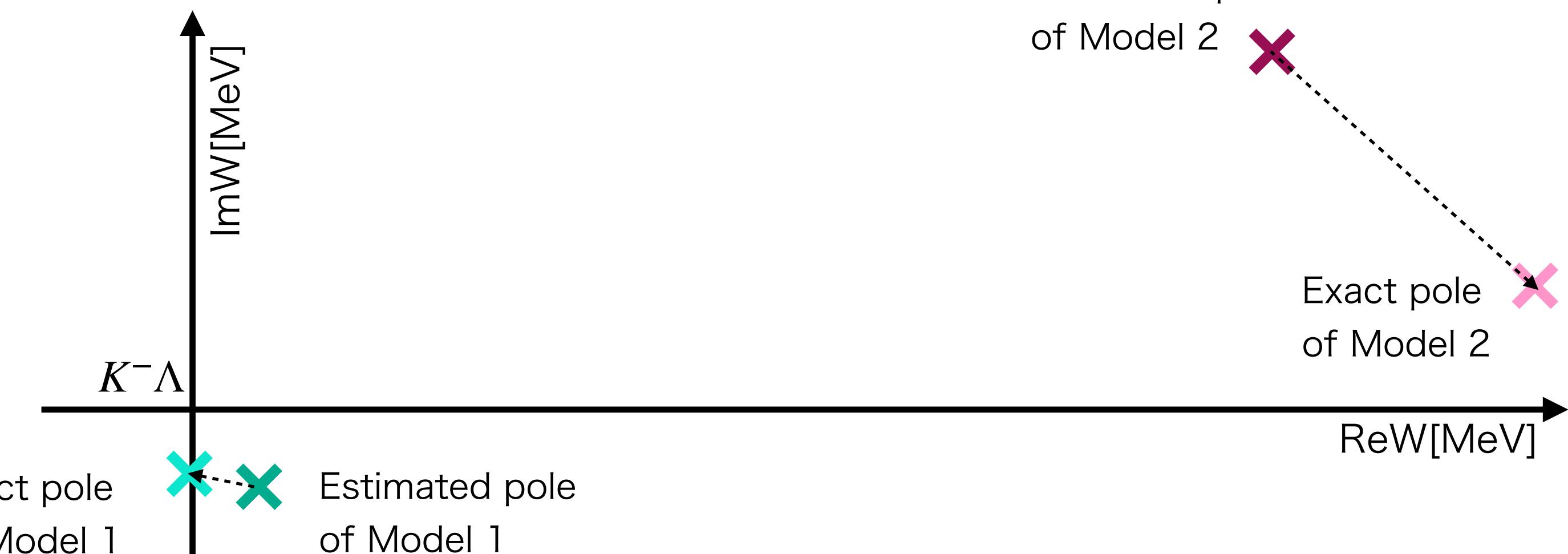
▲The scattering amplitudes of ALICE model

# Estimation of pole position

The pole position  $z$  can be estimated as follows.

$$z \sim \frac{-1}{2\mu_{K^-\Lambda} a_0^2} + M_\Lambda + m_{K^-} \quad (a)$$

$$f = (-1/a_o - ip)^{-1}$$



Model name	$a_0$ [fm]	$ a_0 $ [fm]	Pole position by eq(a) [MeV]	Exact pole position [MeV]	Distance [MeV]
Model 1	$0.80 - 0.92i$	1.21	$1615 - 38i$ [bbtttt]	$1607 - 29i$ [bbtttt]	12.04
Model 2	$-0.27 - 0.40i$	0.48	$1701 + 228i$ [ttbttt]	$1725 + 76i$ [ttbttt]	153.88

- Eq.(a) can predict the Riemann sheet that exact pole locates.
- The prediction for pole position is deviated especially for small  $|a_0|$ .
- The Riemann sheet can be correctly predicted even for small  $|a_0|$ .

# Consistency of ALICE and Belle

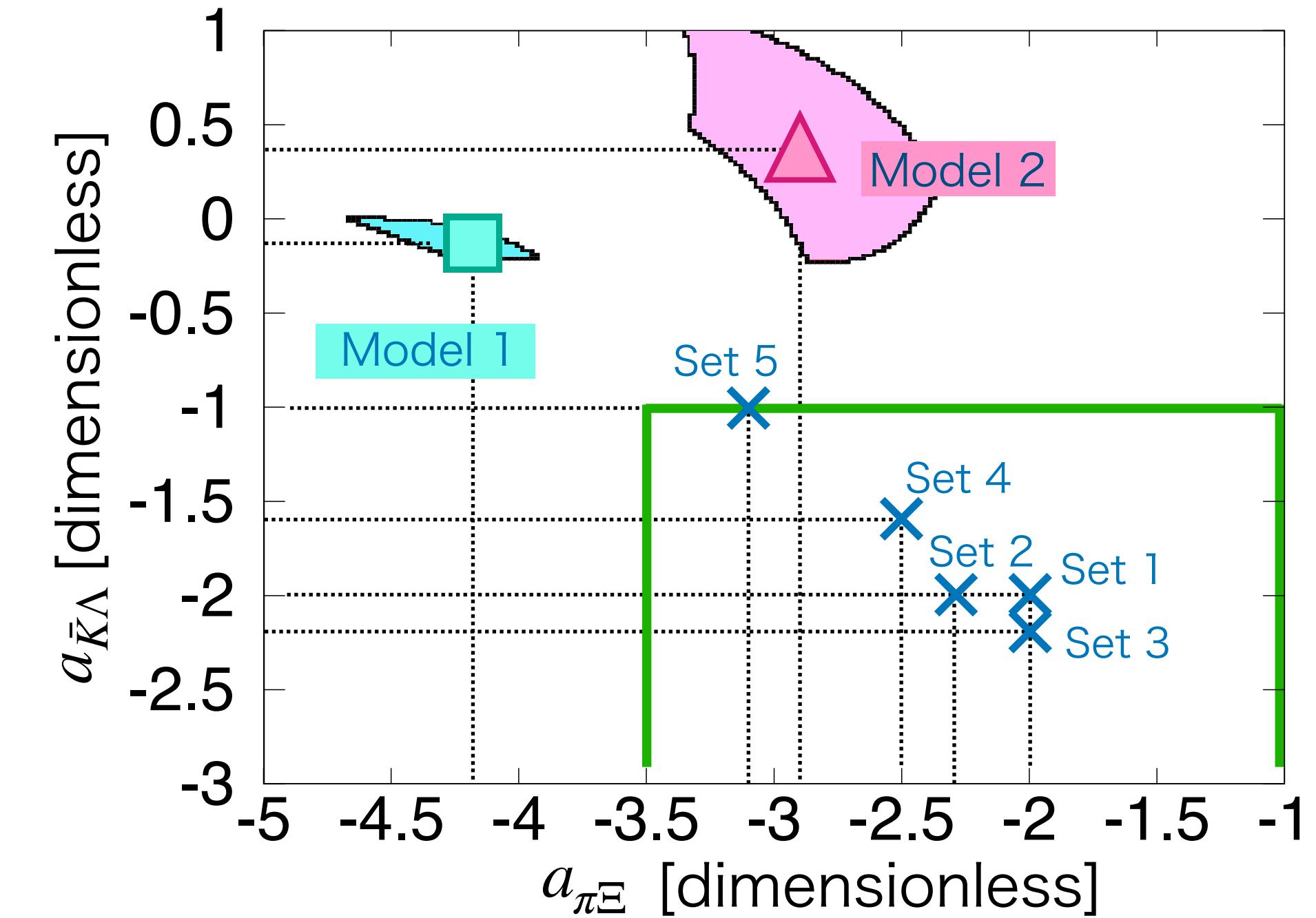
Is there a model which satisfies both Belle and ALICE?

→ We consider the error of each experiment.

$$M_R \simeq 1610.4^{+6.1}_{-7.3} \text{ MeV}, \Gamma_R \simeq 59.9^{+5.6}_{-8.5} \text{ MeV}$$

$$\text{Re}f_0 \simeq 0.27 \pm 0.14 \text{ fm}, \text{Im}f_0 \simeq 0.40 \pm 0.13 \text{ fm}$$

- Figure shows
  - Set 1,2,...,5 in previous study[5],
  - the variable range in Ref. [7],
  - the region which satisfies assumption of pole at  $M_R - i\Gamma_R/2$ ,
  - the region which reproduce  $K^-\Lambda$  scattering length.

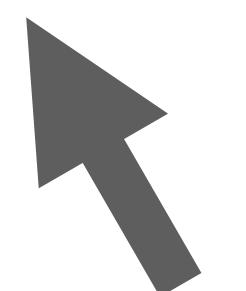


[7]Feijoo, A., Valcarce, V. and Magas, V. K., Phys.Lett.B 841 (2023) 137927.

# Consistency of ALICE and Belle

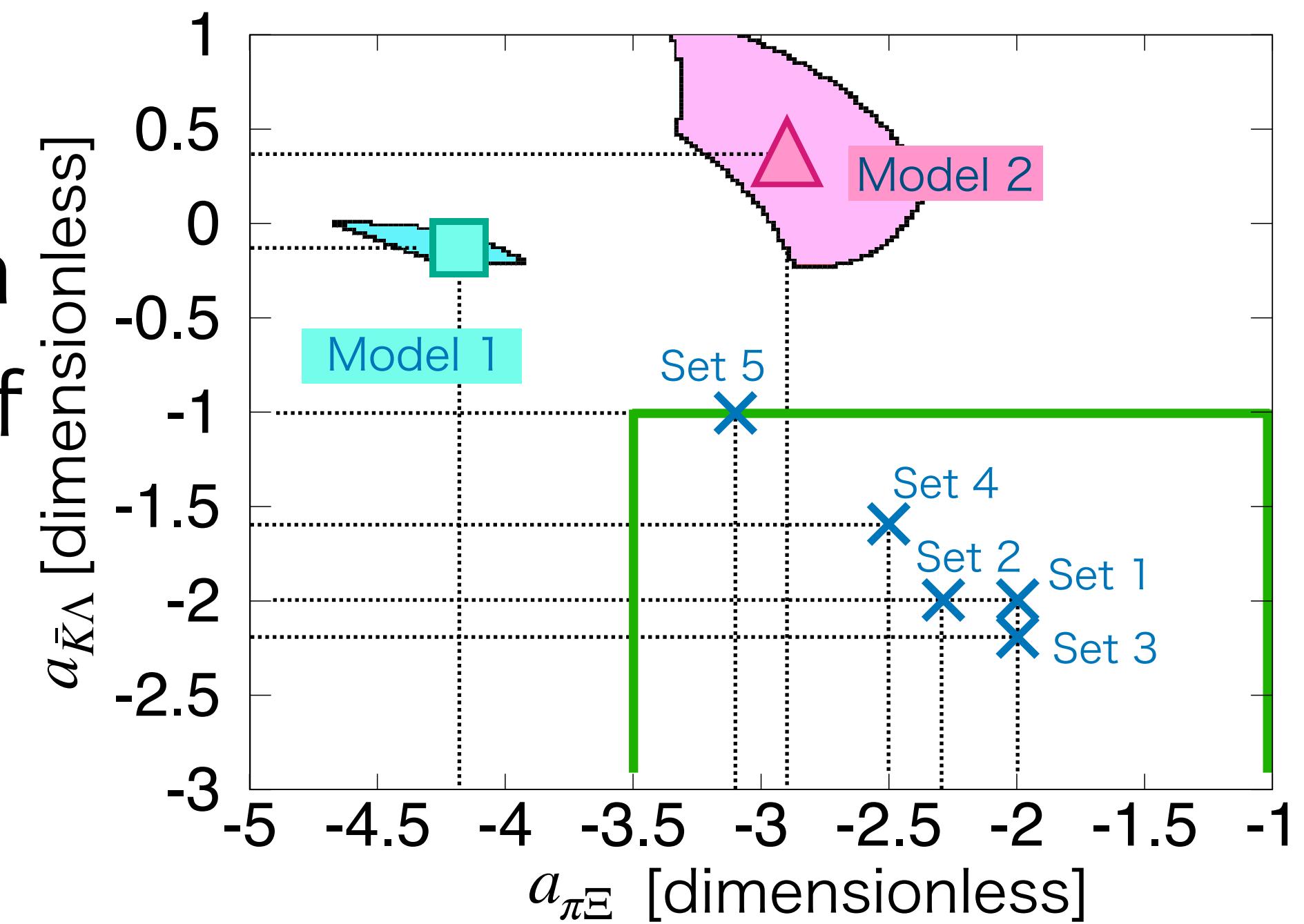
Is there a model which satisfies both Belle and ALICE?

- There is no overlap region which satisfies both ALICE scattering length and the assumption of subthreshold pole.
- The assumption of Model 1 may not be adequate.



Belle used Breit-Wigner fit to the spectrum, and they did not determine the pole position.

- To compare with Belle data, it is better to use the  $\pi\Xi$  spectrum directly.



[5] T.Nishibuchi and T.Hyodo, arXiv:2305.10753 [hep-ph]

# Conclusion

- We construct the models to reproduce the Belle data of the  $\pi\Xi$  spectrum (Model 1) and the  $K^-\Lambda$  scattering length by ALICE (Model 2).
- In Model 1, we find that the near-threshold resonance peak is distorted by the threshold effect.
- In Model 2, the scattering amplitude shows the cusp at  $K^-\Lambda$  threshold. There are no pole of  $\Xi(1620)$  in physically relevant Riemann sheets.
- There is no parameter region which satisfies both ALICE scattering length and the assumption of QB state (subthreshold pole) near the  $K^-\Lambda$  threshold.
- Future plan: study of  $\Xi(1690)$ , calculation of  $\Xi_c \rightarrow \pi\pi\Xi$  decay.

[5]T.Nishibuchi and T.Hyodo, arXiv:2305.10753 [hep-ph]