

Electromagnetic structure of light nuclei from chiral EFT

In collaboration with: [Arseniy Filin](#), Daniel Möller, Vadim Baru, Christopher Körber, Hermann Krebs, Andreas Nogga and Patrick Reinert

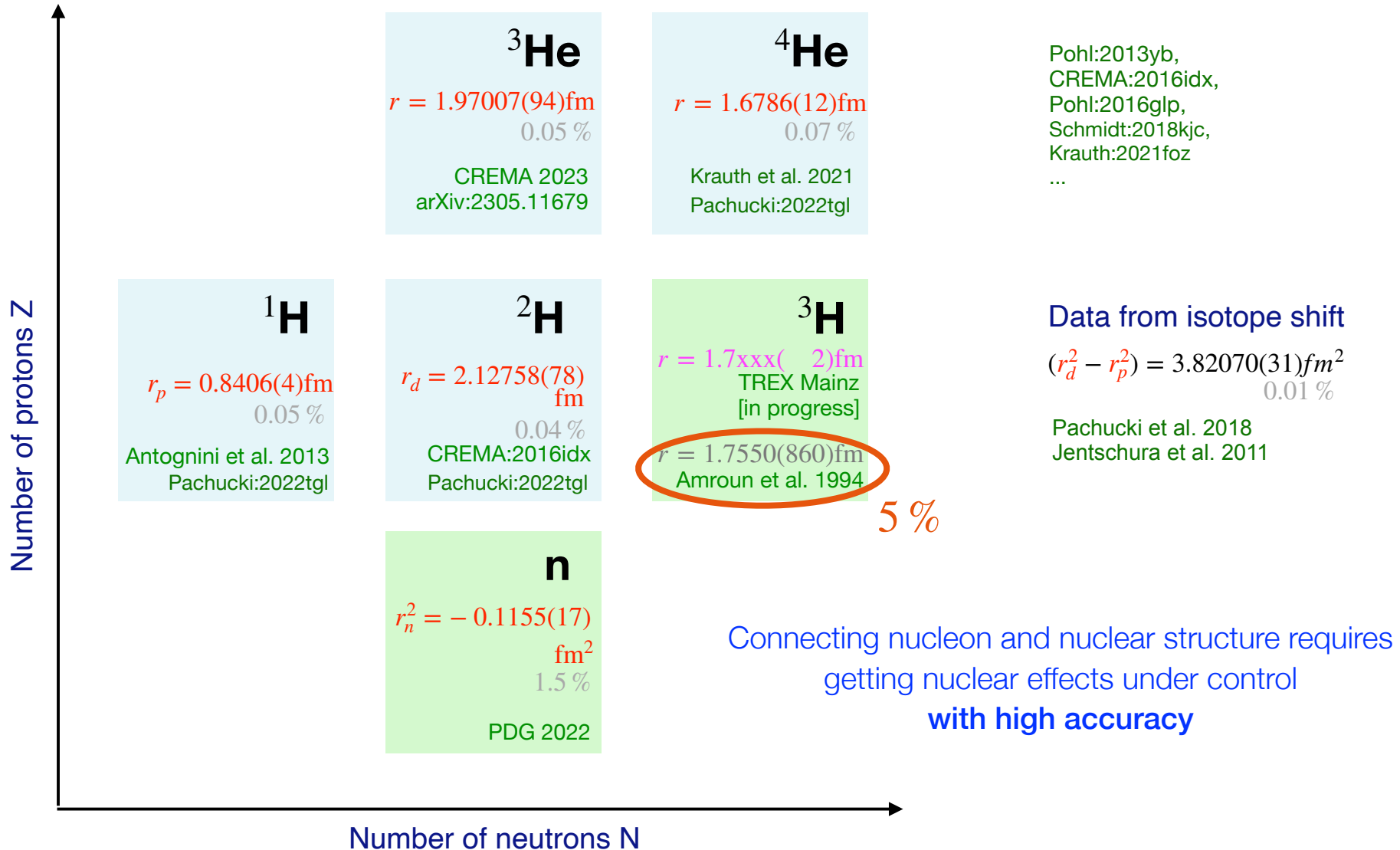


- Introduction
- Nuclear forces and e.m. currents in χ EFT
- E.m. structure of the deuteron
- Beyond the deuteron
- Summary

Updates since 2023:

- Error analysis
- Insights into scaling of MECs
- First results for the magnetic FFs

Experimental data



From QCD to nuclear physics

The Standard Model (QCD, ...)

Schwinger-Dyson , large- N_c , ...

Lattice QCD

Approximate chiral $SU(2)_L \times SU(2)_R$ symmetry

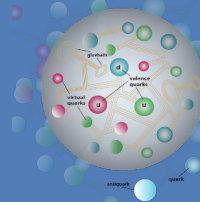
effective chiral Lagrangian $\mathcal{L}_{\text{eff}}(\pi, N)$

Chiral perturbation theory $Q \in \{M_\pi/\Lambda_b, |\vec{p}|/\Lambda_b\}$

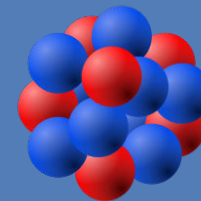
- S-matrix ($\pi\pi$, πN , $\pi\pi N$, ...)
- nuclear forces and currents

Few-body methods

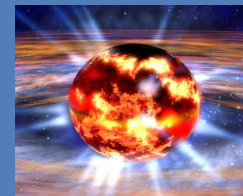
Hadron/nuclear structure and dynamics



proton



nuclei



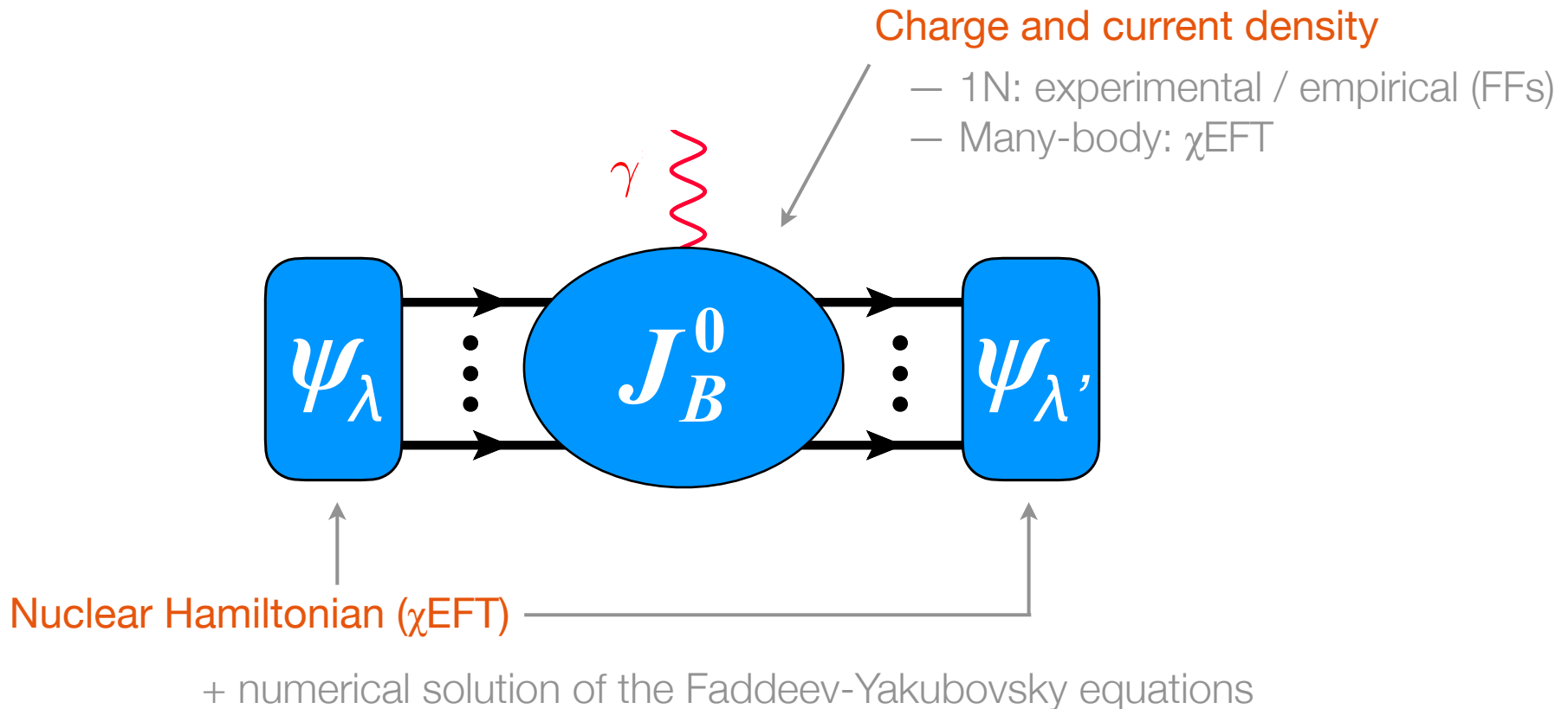
neutron stars

E_{FV}

EFT




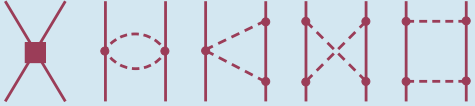





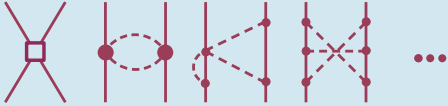
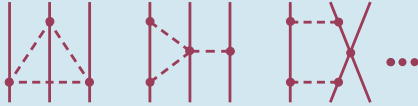

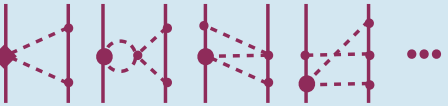


finite-volume methods

Anatomy of the calculation

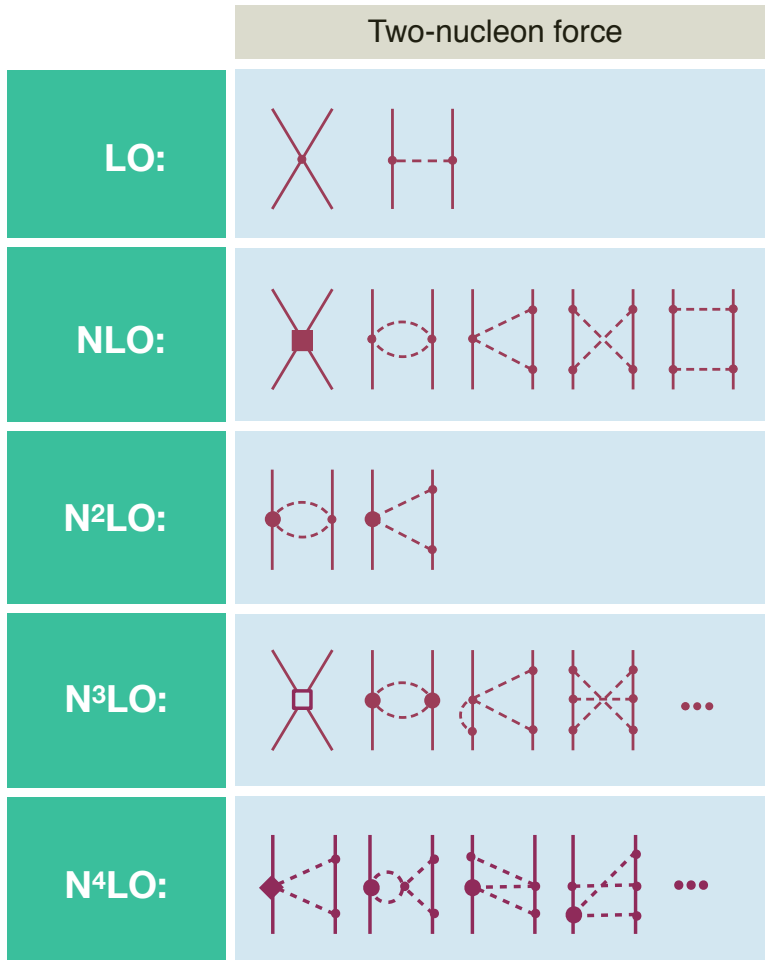


- data driven: Use np + pp data, $^3\text{H}(e)$ and ^4He BEs, up-to-date parametrizations of the nucleon FFs and ^2H , ^4He FF data (assuming 2γ effects are small...)
- no reliance on charge densities (everything expressed in terms of FFs)

Chiral expansion of nuclear forces

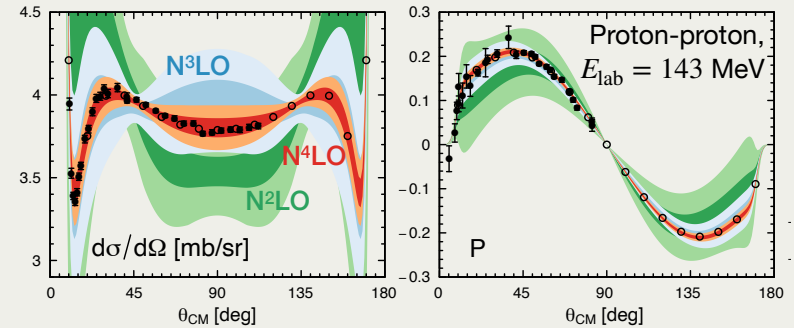
	Two-nucleon force	Three-nucleon force	Four-nucleon force
LO:			
NLO:			
N ² LO:			
N ³ LO:			
N ⁴ LO:			

Chiral expansion of nuclear forces



χ EFT as a precision tool in the 2N sector

- N⁴LO+: currently most accurate and precise NN interactions on the market
- clear evidence of the TPEP from NN data
- Bayesian truncation-error estimation



- almost no residual cutoff dependence

⇒ precision 2N physics from χ EFT




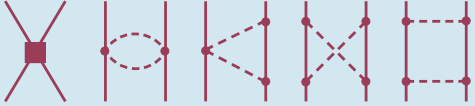





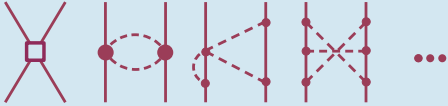
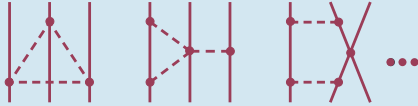

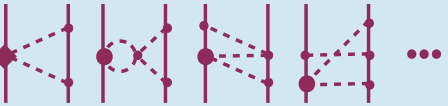


Semi-local regularization in momentum space Reinert, Krebs, EE, EPJA 54 (2018) 86; PRL 126 (2021) 092501

$$V_{1\pi}(q) = \frac{\alpha}{\vec{q}^2 + M_\pi^2} e^{-\frac{\vec{q}^2 + M_\pi^2}{\Lambda^2}} + \text{subtraction,}$$

$$V_{2\pi}(q) = \frac{2}{\pi} \int_{2M_\pi}^{\infty} d\mu \mu \frac{\rho(\mu)}{\vec{q}^2 + \mu^2} e^{-\frac{\vec{q}^2 + \mu^2}{2\Lambda^2}} + \text{subtractions}$$

+ nonlocal (Gaussian) cutoff for contacts

Chiral expansion of nuclear forces

	Two-nucleon force	Three-nucleon force	Four-nucleon force
LO:			
NLO:			
N ² LO:			
N ³ LO:			
N ⁴ LO:			

Chiral expansion of nuclear forces

	Two-nucleon force	Three-nucleon force	Four-nucleon force
LO:			
NLO:			
N ² LO:			
N ³ LO:			
N ⁴ LO:			

Loop diagrams are calculated using DimReg, but the Schrödinger equation is regularized with a cutoff...




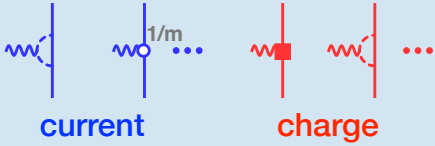


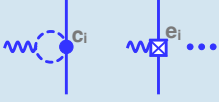
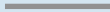


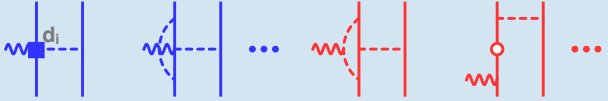



But: Mixing DimReg with CutoffReg violates χ -symmetry EE, Krebs, Reinert, Front. In Phys. 8 (20)

⇒ 3NF beyond N²LO & 4NF must be re-derived using invariant CutoffReg

like, e.g., Gradient Flow Krebs, EE, PRC 110 (2024) 04403; 04404

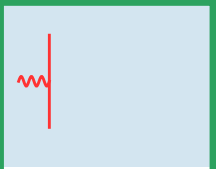

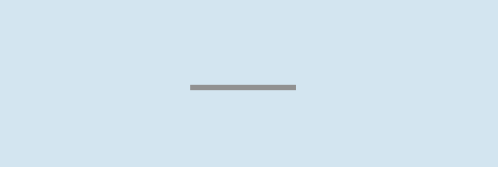
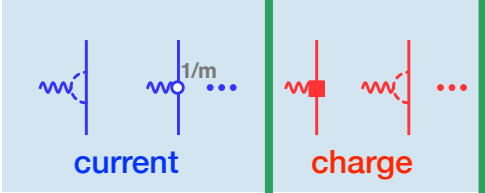
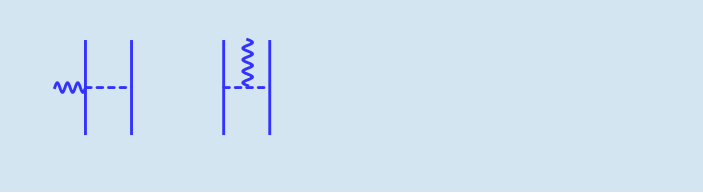
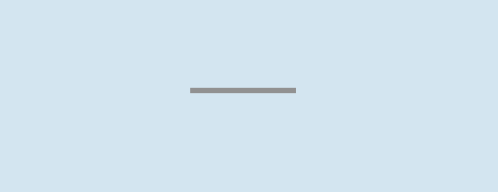
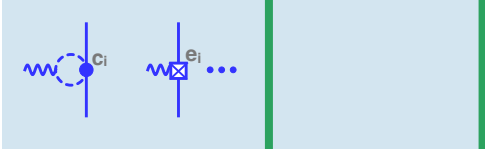


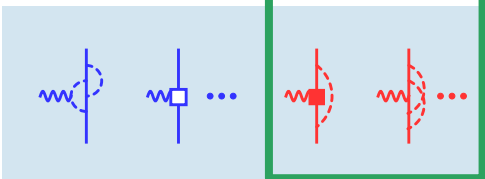
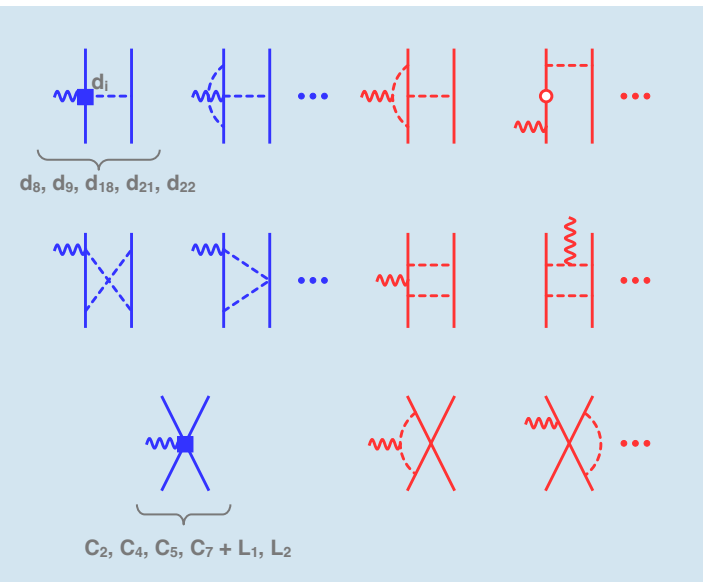
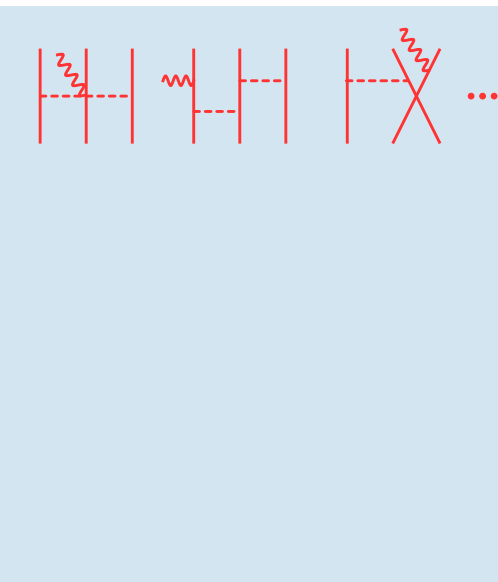
Electromagnetic currents

Kölling, EE, Krebs, Meißner, PRC 80 (09) 045502; PRC 86 (12) 047001; Krebs, EE, Meißner, FBS 60 (2019) 31; Krebs, EPJA 56 (00) 240

	Single-nucleon	Two-nucleon	Three-nucleon
LO			
NLO	 <p>current charge</p>		
N ² LO			
N ³ LO		 <p>$d_8, d_9, d_{18}, d_{21}, d_{22}$</p>   <p>$C_2, C_4, C_5, C_7 + L_1, L_2$</p>	

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Can be parametrized in terms of the nucleon FFs





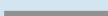
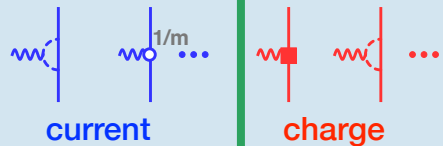


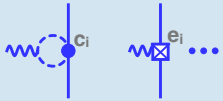
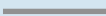
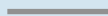
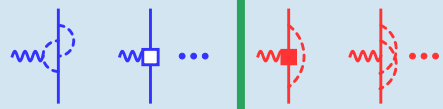

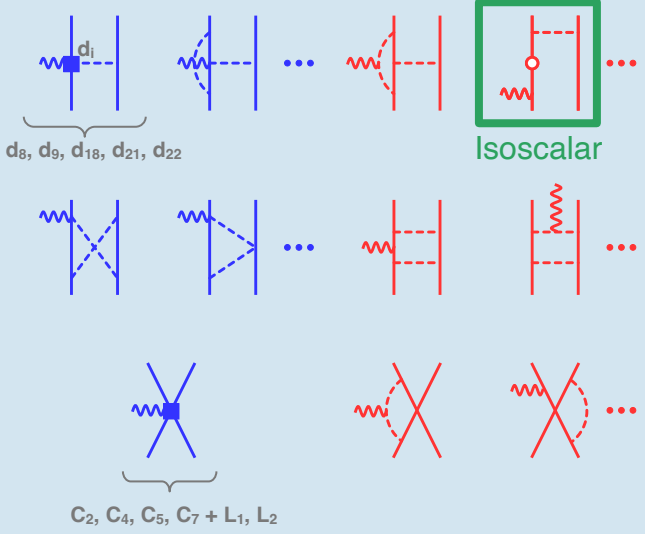
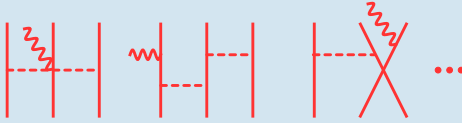
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	Single-nucleon	Two-nucleon	Three-nucleon
LO			
NLO	<p>current</p> <p>charge</p>		
N ² LO			
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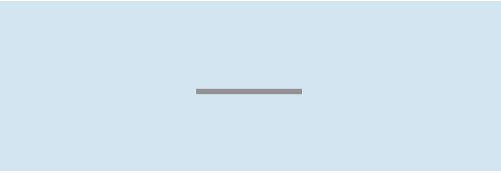
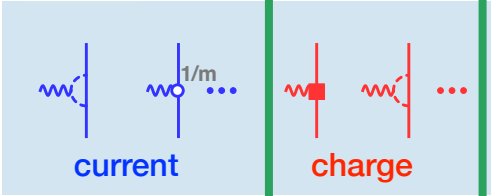

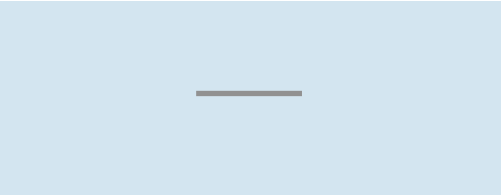
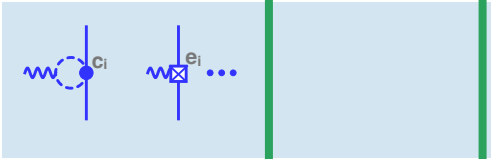
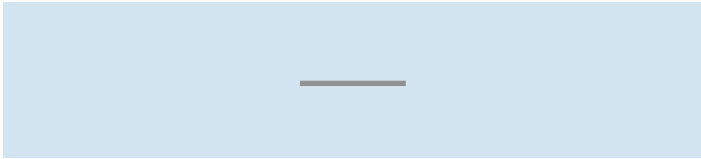
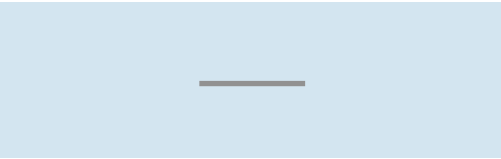
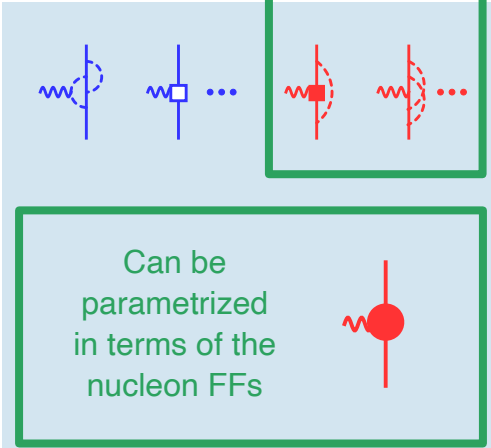
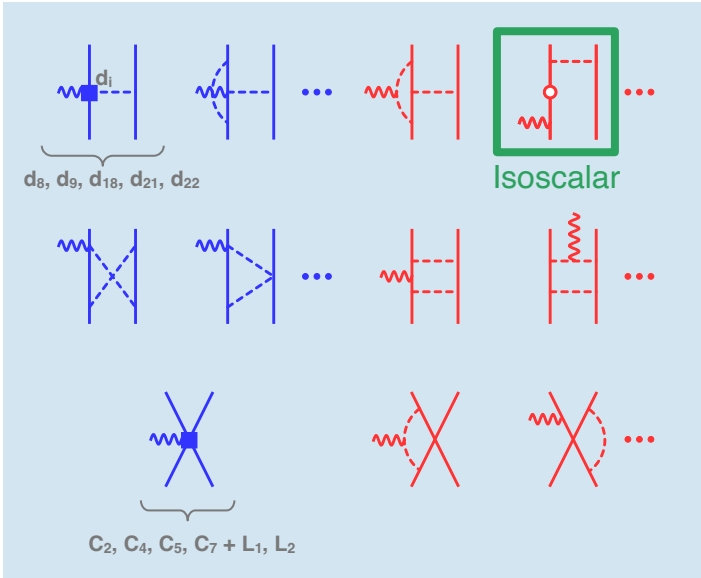
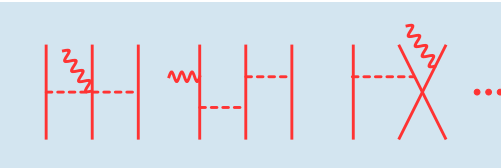
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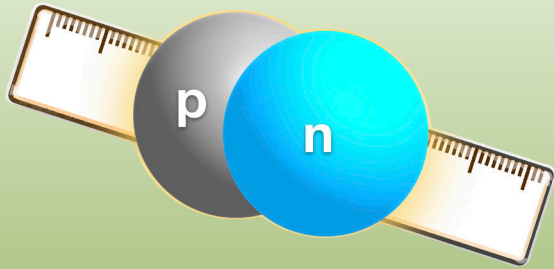
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The deuteron

Arseniy Filin, Vadim Baru, EE, Hermann Krebs, Daniel Möller, Patrick Reinert, Phys. Rev. Lett. 124 (2020) 082501;

Phys. Rev. C103 (2021) 024313

$$\rho_{1N}^{\text{DF}} = -e \frac{\mathbf{k}^2}{8m_N^2} G_E(\mathbf{k}^2)$$

$$G(Q^2) = G^{\text{Main}}(Q^2) + G^{\text{DF}}(Q^2) + G^{\text{SO}}(Q^2) + G^{\text{Boost}}(Q^2) + G^{1\pi}(Q^2) + G^{\text{Cont}}(Q^2)$$

$$\rho_{1N}^{\text{Main}} = e G_E(\mathbf{k}^2)$$

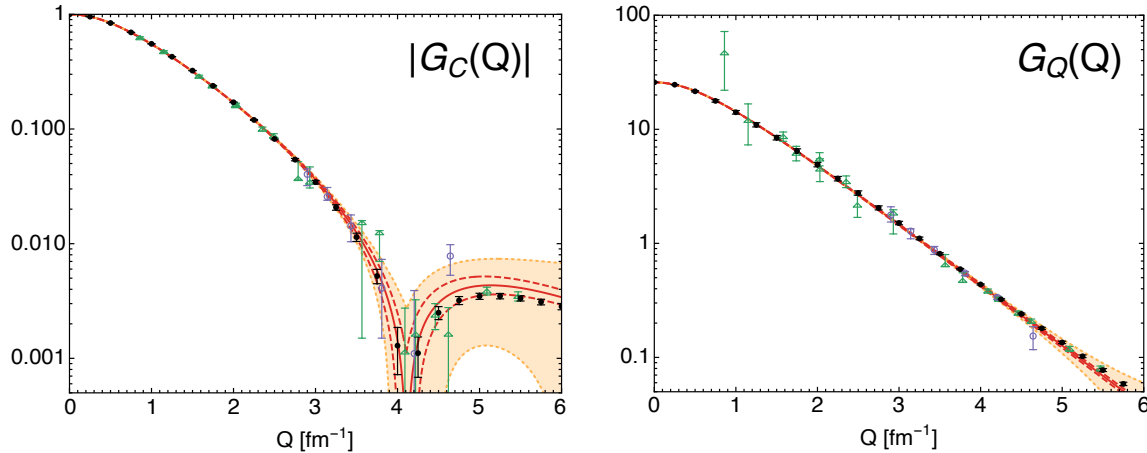
$$\rho_{1N}^{\text{SO}} = ie \frac{2G_M(\mathbf{k}^2) - G_E(\mathbf{k}^2)}{4m_N^2} \boldsymbol{\sigma} \cdot \mathbf{k} \times \mathbf{p}$$

- Both the nuclear force and the 2N charge density are available to N⁴LO
- Simple numerics

Electromagnetic FFs of the deuteron

Filin, Möller, Baru, EE, Krebs, Reinert, PRL 124 (2020) 082501; PRC 103 (2021) 024313

Charge and quadrupole form factors of the deuteron at N⁴LO



Extracted quadrupole moment:

$$Q_d = 0.2854^{+0.0038}_{-0.0017} \text{ fm}^2$$

EFT truncation, choice of fitting range,
NN, π N and γ NN LECs

to be compared with experiment

$$Q_d^{\text{exp}} = 0.285\,699(15)(18) \text{ fm}^2$$

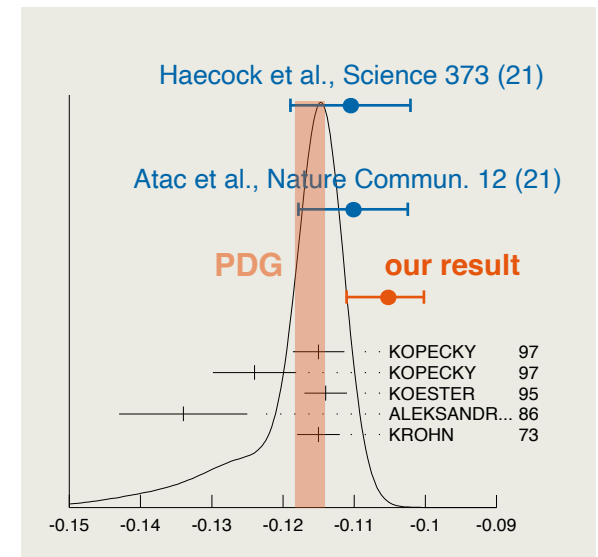
Puchalski et al., PRL 125 (2020)

The charge and structure radius:

$$r_d^2 = (-6) \frac{\partial G_C(Q^2)}{\partial Q^2} \Big|_{Q^2=0} = r_{\text{str}}^2 + r_p^2 + r_n^2 + \frac{3}{4m_p^2}$$

Combining our result $r_{\text{str}} = 1.9729^{+0.0015}_{-0.0012} \text{ fm}$ with very precise isotope-shift spectroscopy data for $r_d^2 - r_p^2$, we determine the neutron m.s. charge radius:

$$r_n^2 = -0.105^{+0.005}_{-0.006} \text{ fm}^2$$



4th moment of the charge distribution

Arseniy Filin et al., preliminary

The fourth-order moment $\langle r_d^4 \rangle := 60 G_C''(0)$ can be measured in the ULQ2 exp [Toshimi Suda et al.](#)

$$\langle r_d^4 \rangle = r_{\text{str}}^{(4)} + \frac{10}{3} r_{\text{str}}^{(2)} \left(r_n^{(2)} + r_p^{(2)} + \frac{3}{4m^2} \right) + \left(r_n^{(4)} + \frac{5}{2m^2} r_n^{(2)} \right) + \left(r_p^{(4)} + \frac{5}{2m^2} r_p^{(2)} + \frac{45}{16m^4} \right)$$

4th moment of the charge distribution

Arseniy Filin et al., preliminary

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$$\langle r_d^4 \rangle = r_{\text{str}}^{(4)} + \frac{10}{3} r_{\text{str}}^{(2)} \left(r_n^{(2)} + r_p^{(2)} + \frac{3}{4m^2} \right) + \left(r_n^{(4)} + \frac{5}{2m^2} r_n^{(2)} \right) + \left(r_p^{(4)} + \frac{5}{2m^2} r_p^{(2)} + \frac{45}{16m^4} \right)$$

Results for $\Lambda = 500$ MeV:
(preliminary, likely to change)

$$r_{\text{str}}^{(4)} = \underbrace{r_{\text{matter}}^{(4)}}_{55.442} + \underbrace{r_{\text{boost}}^{(4)}}_{0.215} + \underbrace{r_{\text{SO}}^{(4)}}_{-0.007} + \underbrace{r_{2\text{N,OPE}}^{(4)}}_{0.025} + \underbrace{r_{2\text{N,CT}}^{(4)}}_{0.008} = 55.68(5) \text{ fm}^4$$

Error budget (preliminary):

Truncation (N⁴LO): ± 0.035 , π N LECs: ± 0.005 , NN LECs: ± 0.04 , other errors negligible

4th moment of the charge distribution

Arseniy Filin et al., preliminary

The fourth-order moment $\langle r_d^4 \rangle := 60 G_C''(0)$ can be measured in the ULQ2 exp [Toshimi Suda et al.](#)

$$\langle r_d^4 \rangle = r_{\text{str}}^{(4)} + \frac{10}{3} r_{\text{str}}^{(2)} \left(r_n^{(2)} + r_p^{(2)} + \frac{3}{4m^2} \right) + \left(r_n^{(4)} + \frac{5}{2m^2} r_n^{(2)} \right) + \left(r_p^{(4)} + \frac{5}{2m^2} r_p^{(2)} + \frac{45}{16m^4} \right)$$

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As an application, one can predict $\langle r_d^4 \rangle = 64.62 \pm 0.05_{\chi\text{EFT}} \pm \underbrace{0.05}_{\text{CODATA + PDG}}_{r_N^{(2)}} \pm \underbrace{0.36}_{r_N^{(4)}} \text{ fm}^4$

(preliminary, likely to change)

CODATA + PDG

$$r_p^{(4)} = 1.3 \pm 0.3 \text{ fm}^4$$

$$r_n^{(4)} = -0.5 \pm 0.2 \text{ fm}^4$$

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Alternative determination of the nucleon isoscalar radius?

1% accuracy \Rightarrow 8% accuracy for $r_n^{(2)} + r_p^{(2)} \dots$

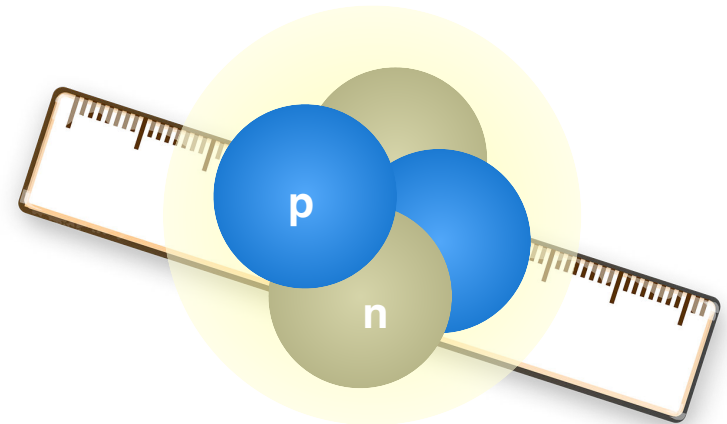
$$r_p^{(2)} + r_n^{(2)} = \frac{\left[\langle r_d^4 \rangle - r_p^{(4)} - r_n^{(4)} \right] - r_{\text{str}}^{(4)} - \frac{15}{16m^4}}{\frac{10}{3} r_{\text{str}}^{(2)} + \frac{5}{2m^2}} - \frac{3}{4m^2}$$

Generalization to $A = 3, 4$ systems

Arseniy Filin, Christopher Körber, Hermann Krebs, Daniel Möller, Andreas Nogga, Patrick Reinert, in preparation

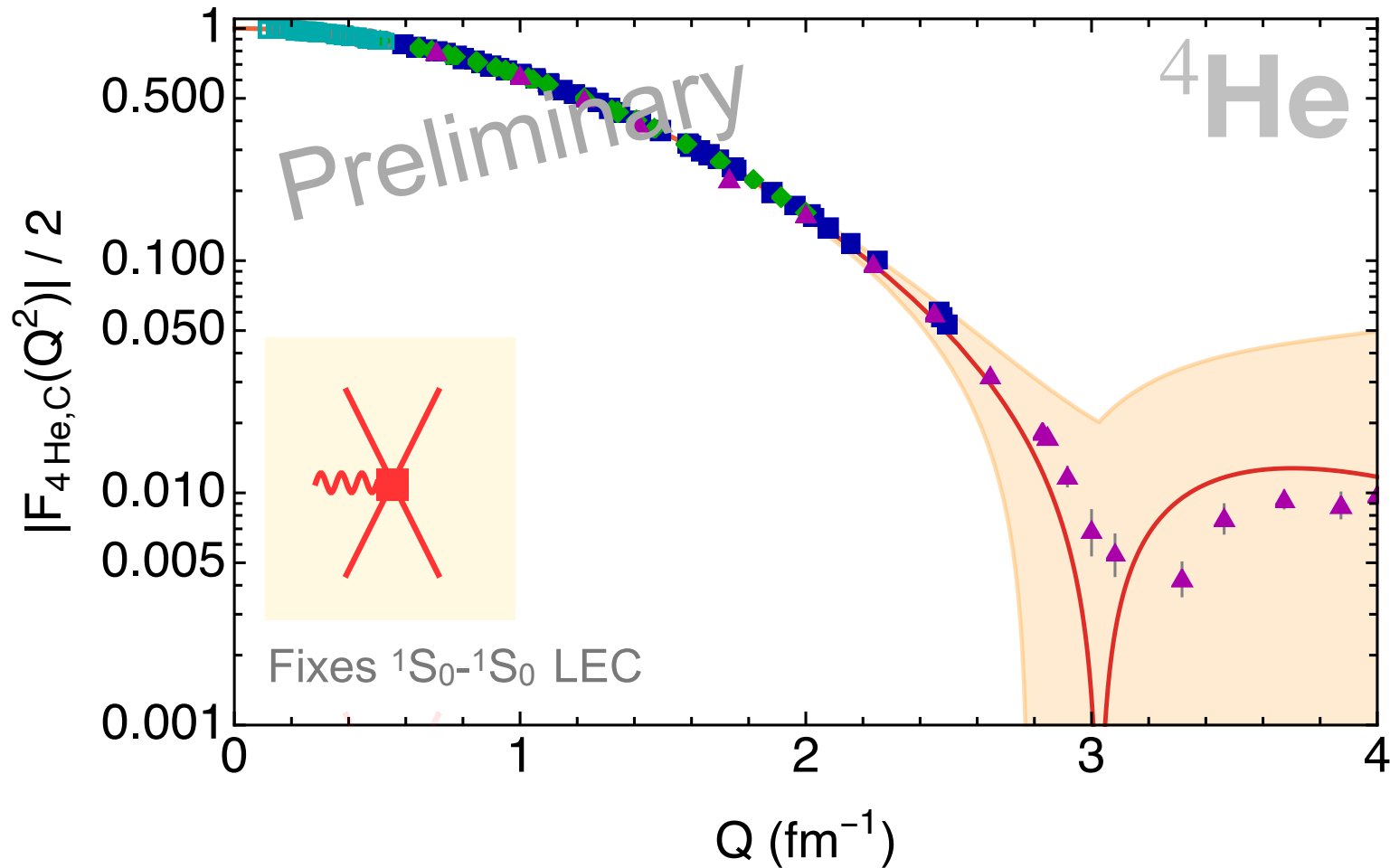
...is a difficult task:

- 3NF beyond N^2LO not yet available
 - use the strong correlations between BEs & matter radii to *implicitly* include 3NF effects
- difficult to account for relativistic corrections beyond 2H
 - irreps of the Poincare group on A -particle Hilbert spaces a-la Polyzou et al.
- iso-vector NN charge density not yet available beyond N^2LO
 - irrelevant for $4He$; for $A = 3$ focus on the iso-scalar linear combination only
- computational accuracy for evaluating the $A = 3, 4$ FFs
 - semi-analytical methods that exploit the local + separable form of the 2N charge density
- many sources of uncertainty
 - comprehensive error analysis



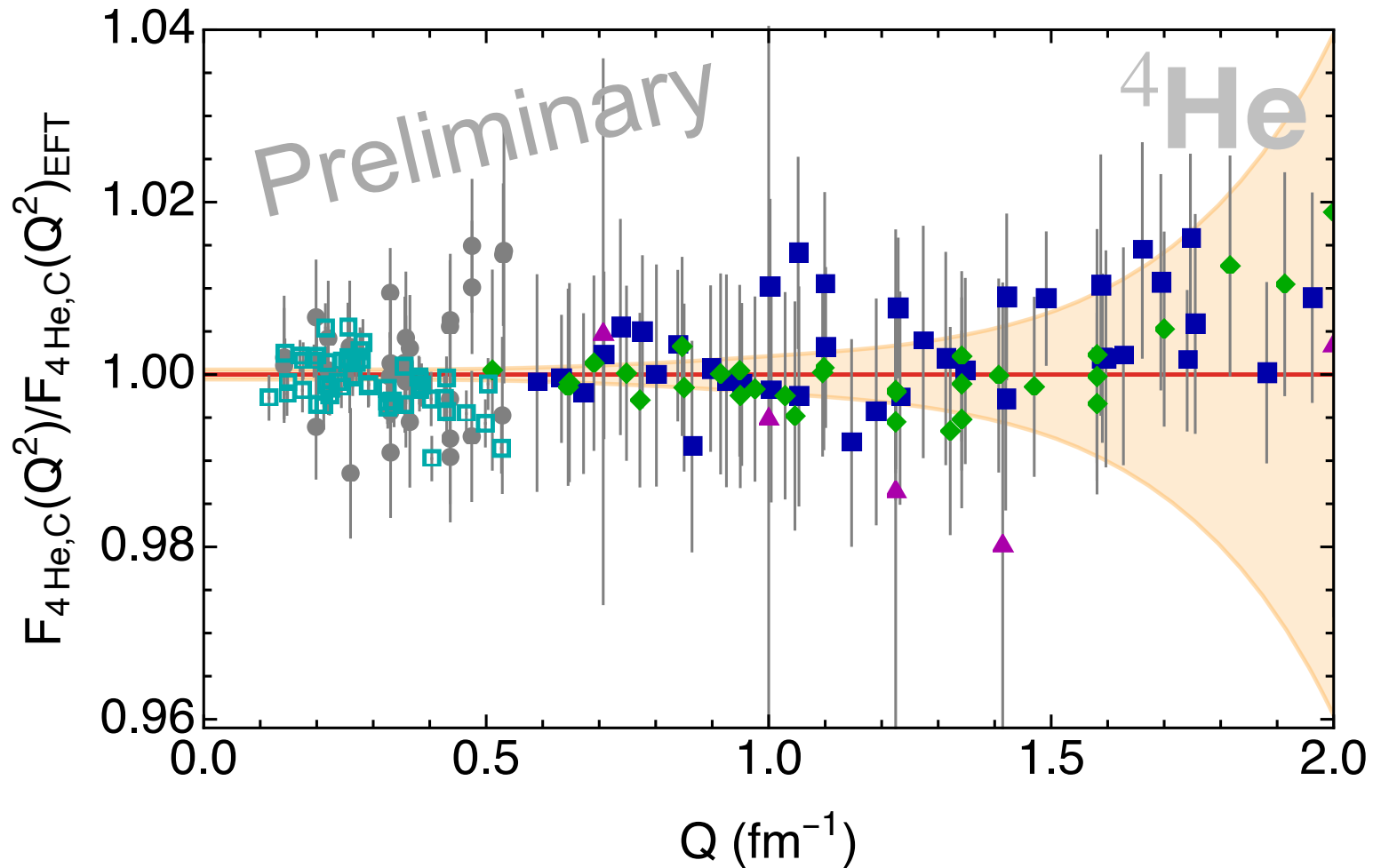
The charge FF of ${}^4\text{He}$

Arseniy Filin, Vadim Baru, EE, Hermann Krebs, Daniel Möller, Andreas Nogga, Patrick Reinert, in preparation



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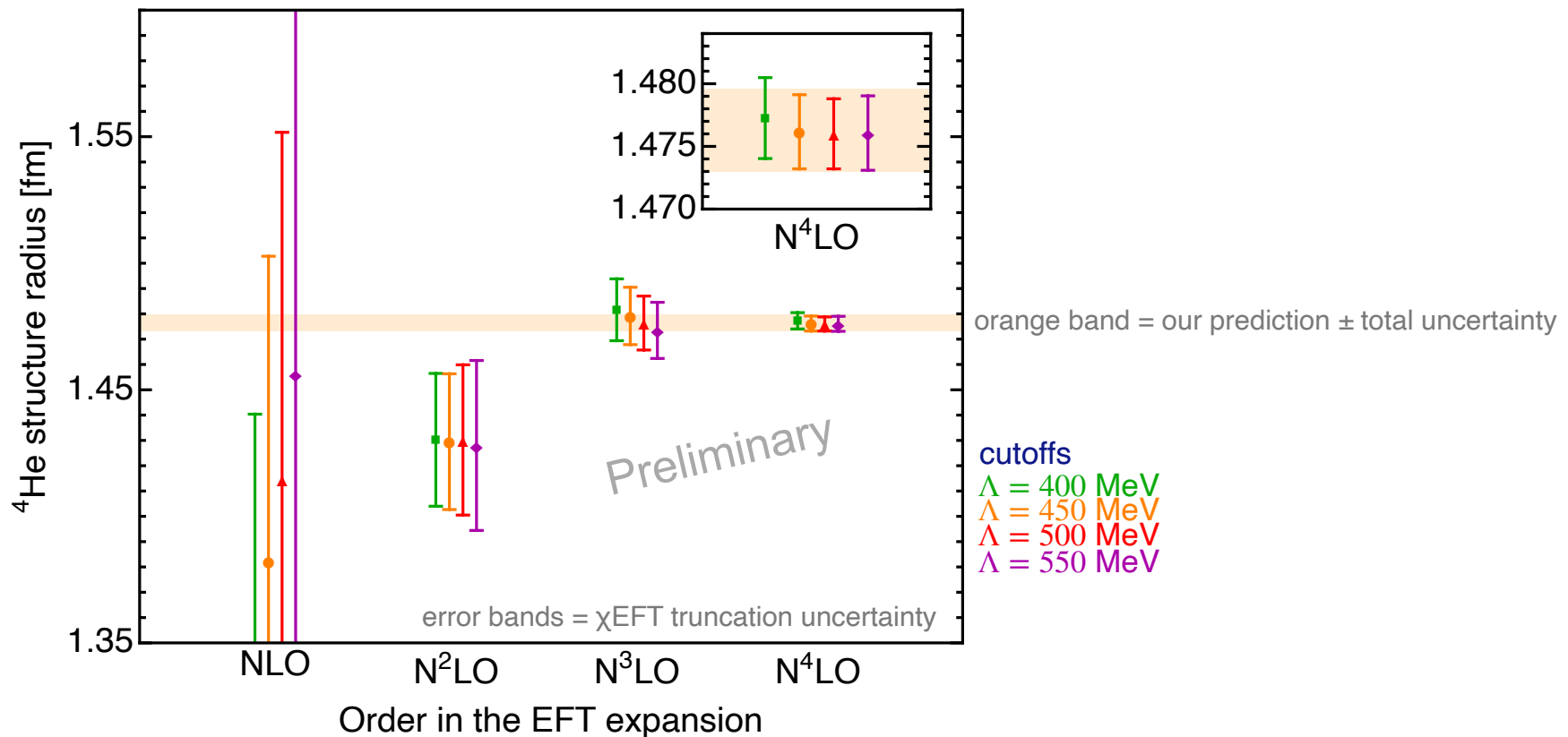
The structure radius of ${}^4\text{He}$

Arseniy Filin, Vadim Baru, EE, Hermann Krebs, Daniel Möller, Andreas Nogga, Patrick Reinert, in preparation

Preliminary result for the ${}^4\text{He}$ structure radius:

$$r_{\text{str}}({}^4\text{He}) = 1.4758 \pm 0.0028_{\text{trunc}} \pm 0.0011_{\text{stat}} \pm 0.0010_{\text{nucl-FF}} \text{ fm (Preliminary)}$$

Consistency check (residual cutoff dependence):



Nucleon size from the ^4He charge radius

Arseniy Filin, Vadim Baru, EE, Hermann Krebs, Daniel Möller, Andreas Nogga, Patrick Reinert, in preparation

Alternatively: Nucleon size from ^4He radius

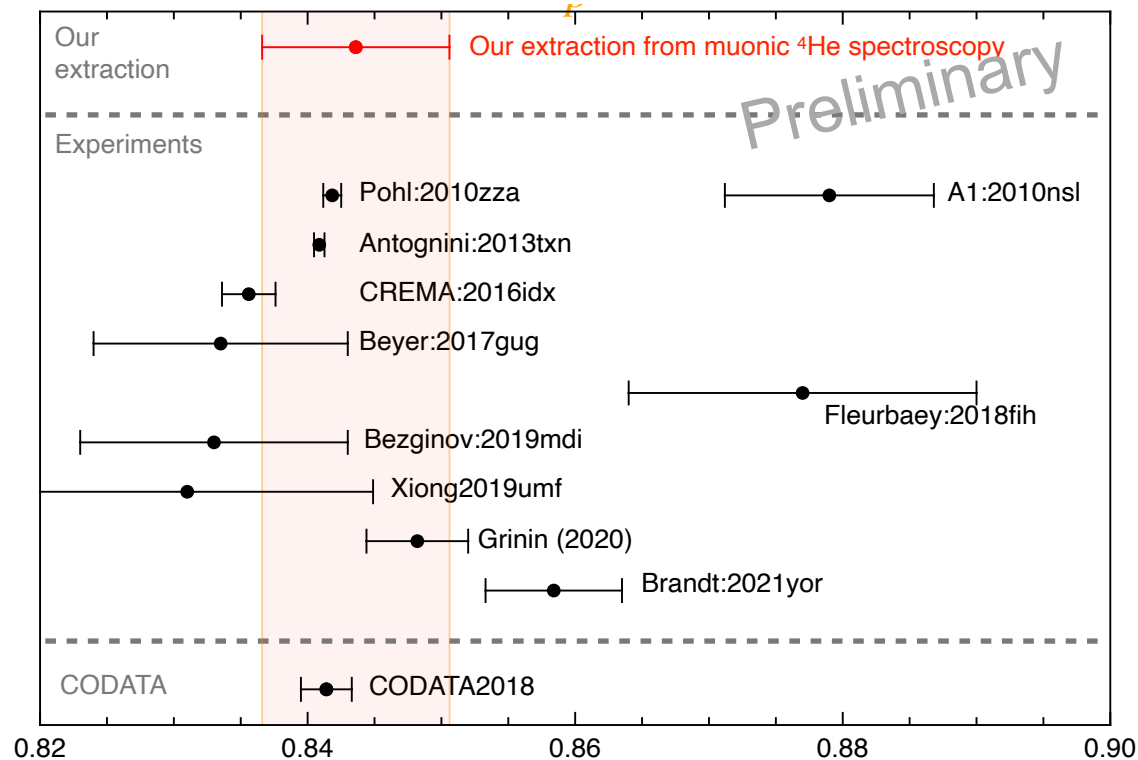
$$r_C^2(^4\text{He}) = r_{\text{str}}^2(^4\text{He}) + \left(r_p^2 + \frac{3}{4m_p^2} \right) + r_n^2 \Rightarrow$$

$$r_p^2 + r_n^2 = (0.607 \pm 0.010) \text{ fm}^2$$

$$r_p = (0.844 \pm 0.007) \text{ fm}$$

preliminary (own determination of r_n)

Proton charge radius



Isoscalar charge radius of ${}^3\text{H}$, ${}^3\text{He}$

Arseniy Filin, Vadim Baru, EE, Hermann Krebs, Daniel Möller, Andreas Nogga, Patrick Reinert, in preparation

Predicted value of the isoscalar 3N charge radius $r_C^{\text{isoscalar}} = \sqrt{\frac{1}{3}(r_C^{3\text{H}})^2 + \frac{2}{3}(r_C^{3\text{He}})^2}$

$$r_C^{\text{isoscalar}} = (1.9060 \pm 0.0026) \text{ fm}$$

preliminary (own determination of r_n)

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Arseniy Filin, Vadim Baru, EE, Hermann Krebs, Daniel Möller, Andreas Nogga, Patrick Reinert, in preparation

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$$r_C^{\text{isoscalar}} = (1.9060 \pm 0.0026) \text{ fm}$$

preliminary (own determination of r_n)

Experimental value: $r_{C, \text{exp}}^{\text{isoscalar}} = (1.9010 \pm 0.0260) \text{ fm}$

error dominated by the ${}^3\text{H}$ datum

⇒ our prediction is 10x more precise than the current experimental value

The ongoing T-REX experiment in Mainz [Pohl et al.] aims at measuring the ${}^3\text{H}$ charge radius within $\pm 0.0002 \text{ fm}$ (i.e., 400x more precise) ⇒ the isoscalar radius will be known within $\pm 0.0009 \text{ fm}$

⇒ precision test of nuclear chiral EFT

${}^3\text{He}$

$r = 1.97007(94) \text{ fm}$
0.05 %

CREMA 2023
arXiv:2305.11679

${}^3\text{H}$

$r = 1.7\text{xxx}(\quad) \text{ fm}$
T-REX Mainz
[in progress]

$r = 1.7550(860) \text{ fm}$
Amroun et al. 1994

Scaling of our results

Arseniy Filin, Vadim Baru, EE, Hermann Krebs, Daniel Möller, Andreas Nogga, Patrick Reinert, in preparation

$$r_C^2(^2H) = 1r_p^2 + 1r_{DF}^2 + 1r_n^2 + 3.8746 - 0.0161M_1 - 0.0047M_2$$

$$r_C^2(^3H) = 1r_p^2 + 1r_{DF}^2 + 2r_n^2 + 2.5148 - 0.0326M_1 - 0.0089M_2 - 0.0238M_3 + 0.0158M_4$$

$$2r_C^2(^3He) = 2r_p^2 + 2r_{DF}^2 + 1r_n^2 + 6.2821 - 0.0321M_1 - 0.0088M_2 - 0.0230M_3 - 0.0151M_4$$

$$2r_C^2(^4He) = 2r_p^2 + 2r_{DF}^2 + 2r_n^2 + 4.1117 - 0.0765M_1 - 0.0189M_2 - 0.0609M_3$$

nucleon size

structure radii

Slope of charge FF
assuming $F(0) = Z$

$$Z \times r_C^2$$

Cutoff 450 MeV

all numbers in fm²

All LECs in dimensionless natural units

M4 is isovector contact interaction similar to isoscalar M3

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Arseniy Filin, Vadim Baru, EE, Hermann Krebs, Daniel Möller, Andreas Nogga, Patrick Reinert, in preparation

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proportional to
the number of protons
proportional to
the number of neutrons
proportional to
the number of pn pairs
proportional to
the number of
(pp pairs + nn
pairs)
proportional to
the number of
(pp pairs - nn
pairs)

$Z \times r_C^2$
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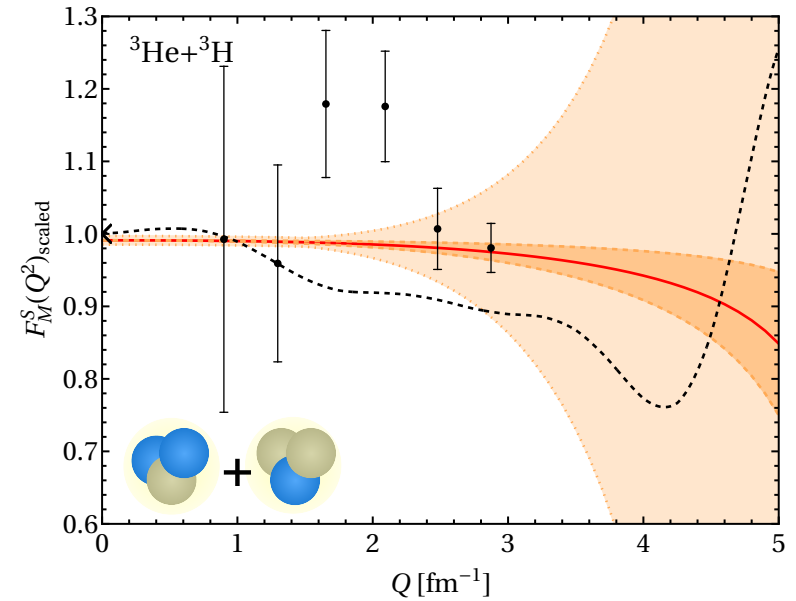
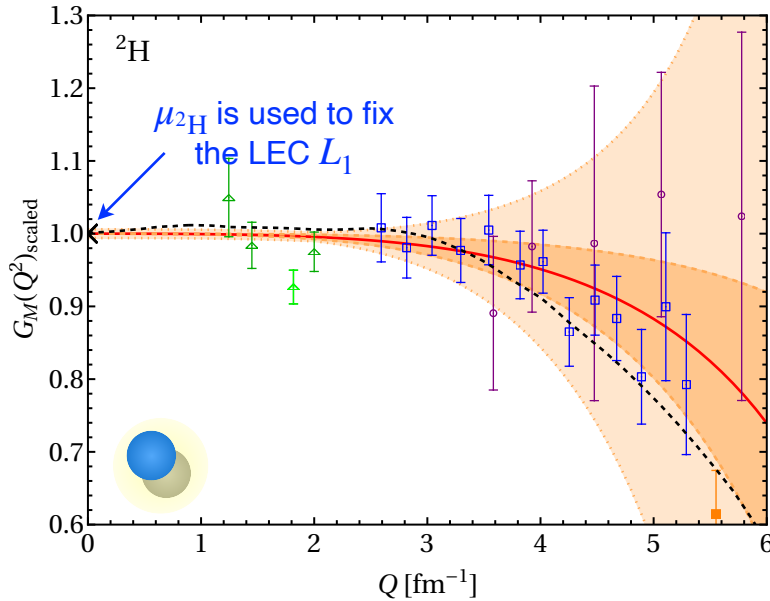
all numbers in fm²

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M4 is isovector contact interaction similar to isoscalar M3

Magnetic FFs of ^2H and $^3\text{H}(e)$

Daniel Möller, PhD thesis, Bochum, 2024; Daniel Möller et al., in preparation



	$\mu_{^3\text{H}(e)}^{\text{iso-s}}$ [in μ_N]	$r_{^2\text{H}}^2$ [fm^2]	$r_{^3\text{H}(e), \text{iso-s}}^2$ [fm^2]
Theory (N ⁴ LO)	0.4219(25) _{tr}	4.384(26) _{tr} (138) _{nucl}	2.280(14) _{tr} (128) _{nucl}
Experiment	0.425669	4.29(7) Sick '01	2.1(2.4) Sick '14; Amroun et al. '94

Summary and outlook

- Charge & quadrupole FFs of ${}^2\text{H}$ are in good shape (N⁴LO, high-precision)
- Other systems and processes are limited to N²LO accuracy due to unavailability of (consistently regularized) many-body forces & exchange currents
 - ⇒ **symmetry-preserving gradient flow regularization** Krebs, EE '24
- Correlations between BEs and radii can be employed to obtain precise results for the charge FFs of ${}^4\text{He}$ & ${}^3\text{H}(\text{e})_{\text{isoscalar}}$ already at this stage Arseniy Filin et al., in progress
- ${}^4\text{He}$: Nuclear effects under control ⇒ new source of information about 1N radii
- ${}^3\text{He}/{}^3\text{H}$: prediction for the isoscalar 3N charge radius is more precise than exp

Thank you for your attention

Error analysis for the radii (ongoing)

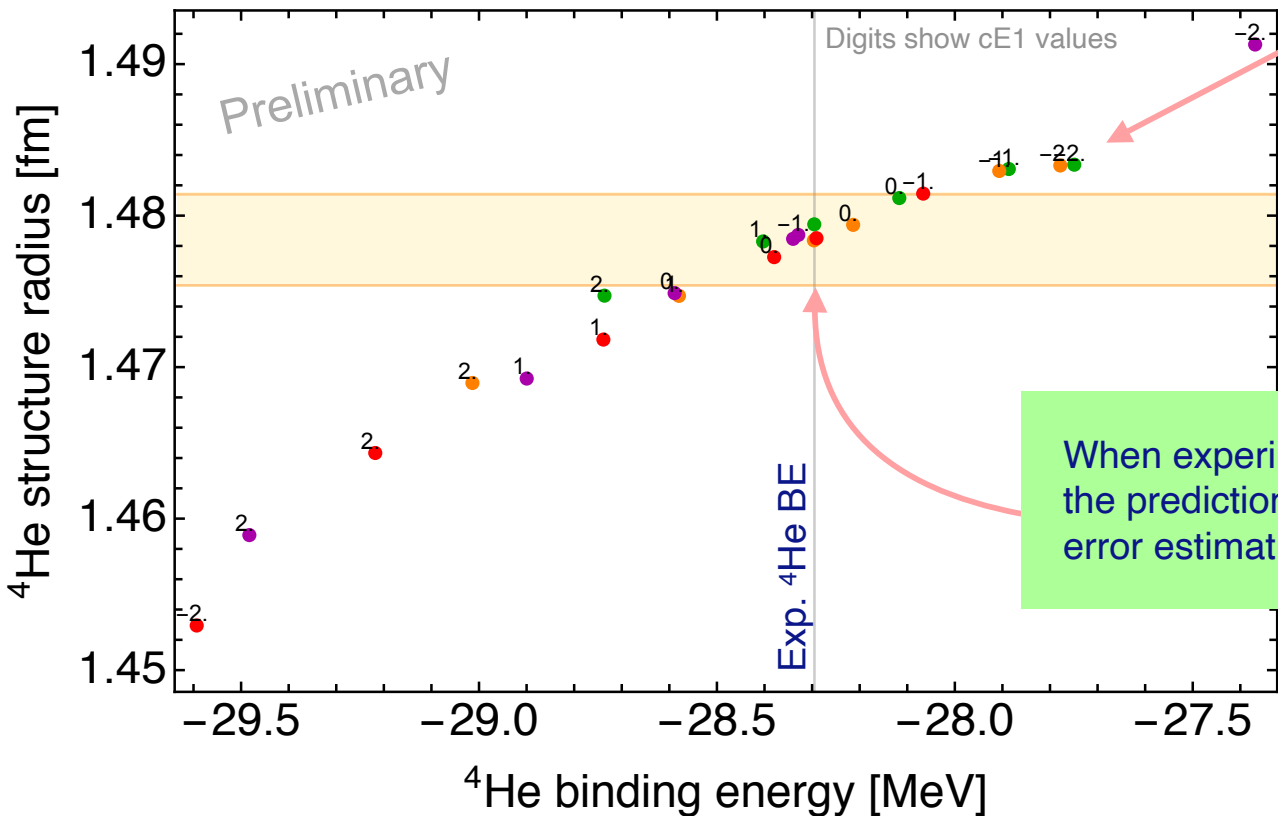
Type of error	$\delta r_{A=3, \text{str}}^2$	$\delta r_{4\text{He}, \text{str}}^2$
Statistical uncertainties of LECs determined from πN, NN, $p^2\text{H}$, $e^2\text{H}$ and $e^4\text{He}$ data		
(a) NN LECs extracted from the database ²³ of mutually consistent np and pp data	14	6
(b) πN LECs from matching χEFT to the Roy-Steiner analysis of πN scattering ⁷⁹	4	3
(c) LEC c_D in the 3NF extracted from $p^2\text{H}$ differential cross section data of ⁵²	4	6
(d) LECs M_i in $\hat{\rho}_{2\text{N}}$ extracted from ^2H ^{86,87} and ^4He ⁶⁹⁻⁷⁵ form factor data	21	32
Total statistical uncertainty	26	33
Theory uncertainties		
(e) Truncation of the χEFT expansion	69	83
(f) Parametrizations of the nucleon form factors	XXX	30
Systematic uncertainties		
(g) Approximate treatment of relativistic effects	5	18

all numbers very preliminary...

Implicit treatment of higher-order 3NFs

^4He structure radius vs BE

^4He binding energy and r_{str} are strongly correlated!



When experimental BE is reproduced the prediction for r_{str} is consistent with our error estimation (fits inside the orange band)

$\Lambda = 400$ MeV
 $\Lambda = 450$ MeV
 $\Lambda = 500$ MeV
 $\Lambda = 550$ MeV