

ELPH workshop C033, Dec. 06-07, 2022

$N - \phi$ interaction from lattice QCD

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Dec. 07, 2022

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 - why $N\phi$
 - why lattice QCD
- Theoretical framework
 - Lüscher's finite volume method
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- Summary and outlook

$N\phi$ interaction

➤ Fundamental inputs for studying

R. S. Hayano and T. Hatsuda, Rev. Mod. Phys. **82**, 2949 (2010)

- meson properties in nuclear matter
- chiral symmetry and its restoration

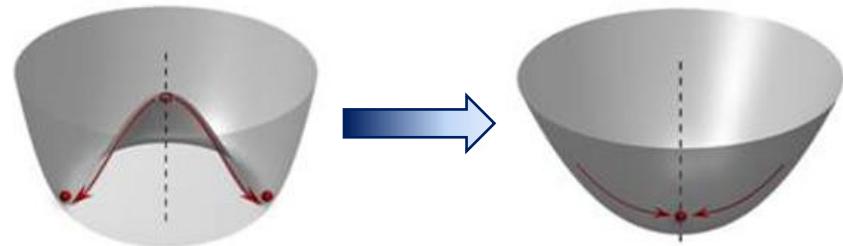
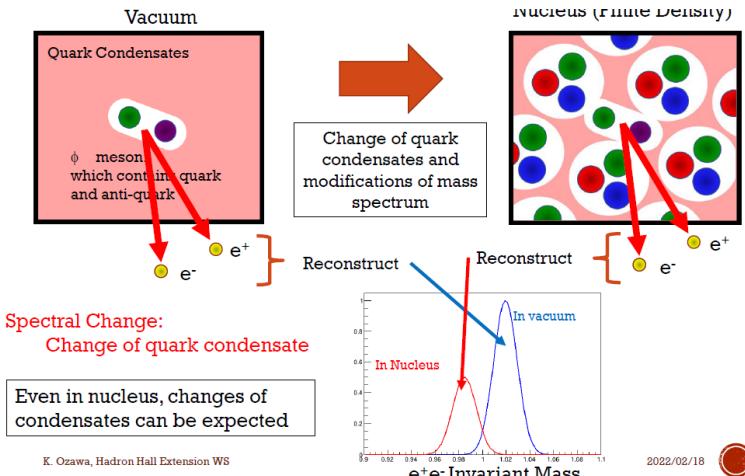


Fig. taken from K. Ozawa

➤ Theoretical studies

P. Gubler and D. Satow, Prog. Part. Nucl. Phys. **106**, 1 (2019); T. Hatsuda and S. H. Lee, Phys. Rev. C **46**, R34 (1992)
J. Castella and G. Krein, Phys. Rev. D **98**, 014029 (2018)

- ϕ in-medium property $\leftrightarrow \bar{s}s$ condensate
- open questions on $N\phi$: microscopic origin, the role of dynamical π

Experiments on $N\phi$ interaction

➤ Production and absorption experiments

- Spring8-LEPS
- KEK-PS E325
- JLab-CLAS
- COSY-ANKE

T. Ishikawa *et al.*, Phys. Lett. B **608**, 215 (2005)

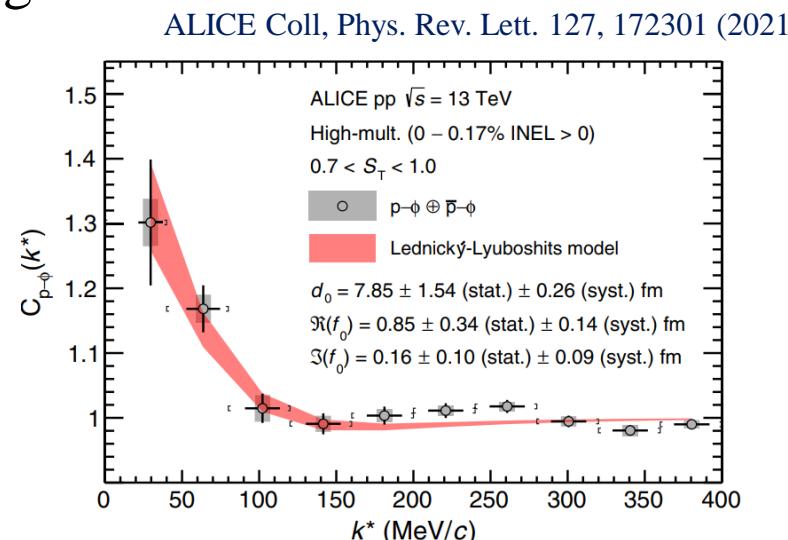
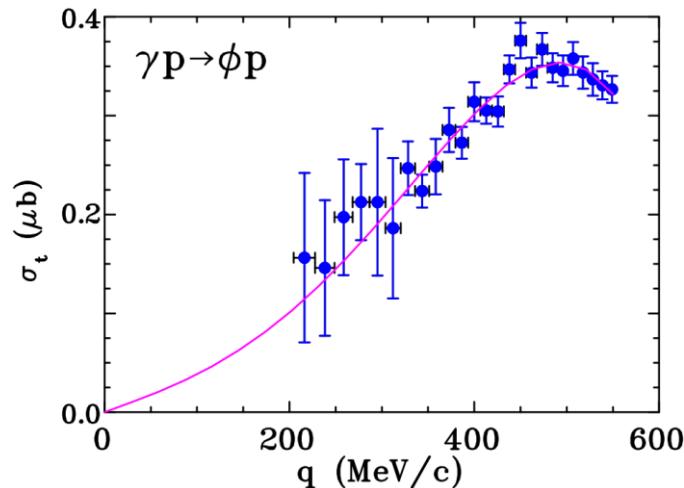
R. Muto *et al.* (KEK-PS-E325 Coll.), Phys. Rev. Lett. **98**, 042501 (2007)

M. Wood *et al.* (CLAS Coll.), Phys. Rev. Lett. **105**, 112301 (2010)

A. Polyanskiy *et al.*, Phys. Lett. B **695**, 74 (2011)

➤ Recent measurements on scattering length

I. Strakovsky, L. Pentchev, and A. Titov, Phys. Rev. C 101, 045201 (2020)



Reference	Spin-averaged a_0 [fm]
Strakovsky <i>et al.</i> 2020	$\pm 0.063(10)$
ALICE 2021	$-0.85(34)_{\text{stat.}}(14)_{\text{syst.}}$

➤ A first principle study on $N\phi$ interaction is highly needed

QCD at different scales

- Strong interaction strength strongly depends on energy scale (Q)

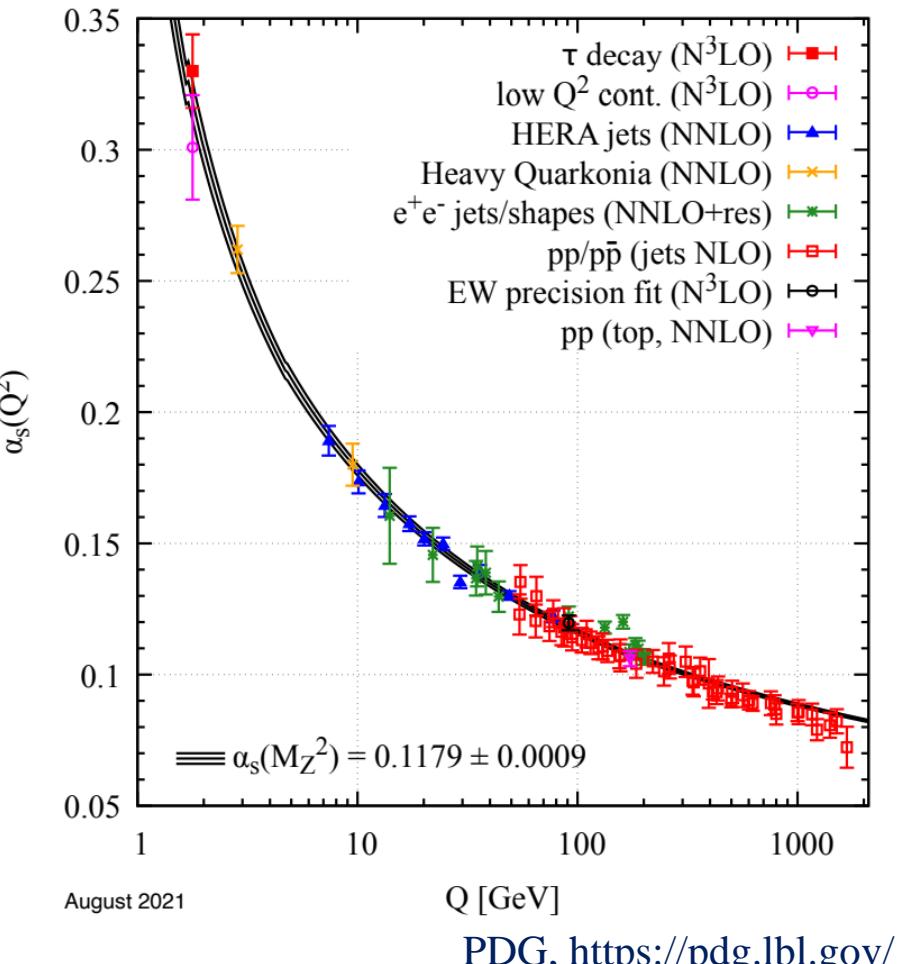
- High Q ($>$ few GeV) region, perturbative

- ✓ asymptotic freedom
- ✓ weakly interacting quarks and gluons
- ✓ small coupling
- ✓ pen and paper



- Low Q (< 1 GeV) region, non perturbative

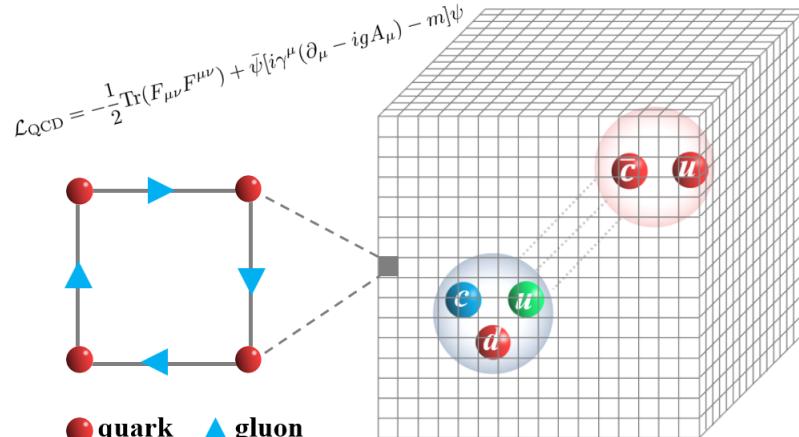
- ✓ chiral symmetry spontaneous breaking
- ✓ strongly interacting quarks and gluons
- ✓ large coupling
- ✓ lattice QCD



➤ Lattice discretization



K.G. Wilson
(1974)



- quarks: lattice site, $q(x)$
- gluons: link between lattice sites, $U_\mu(x) = e^{iagA_\mu(x)}$

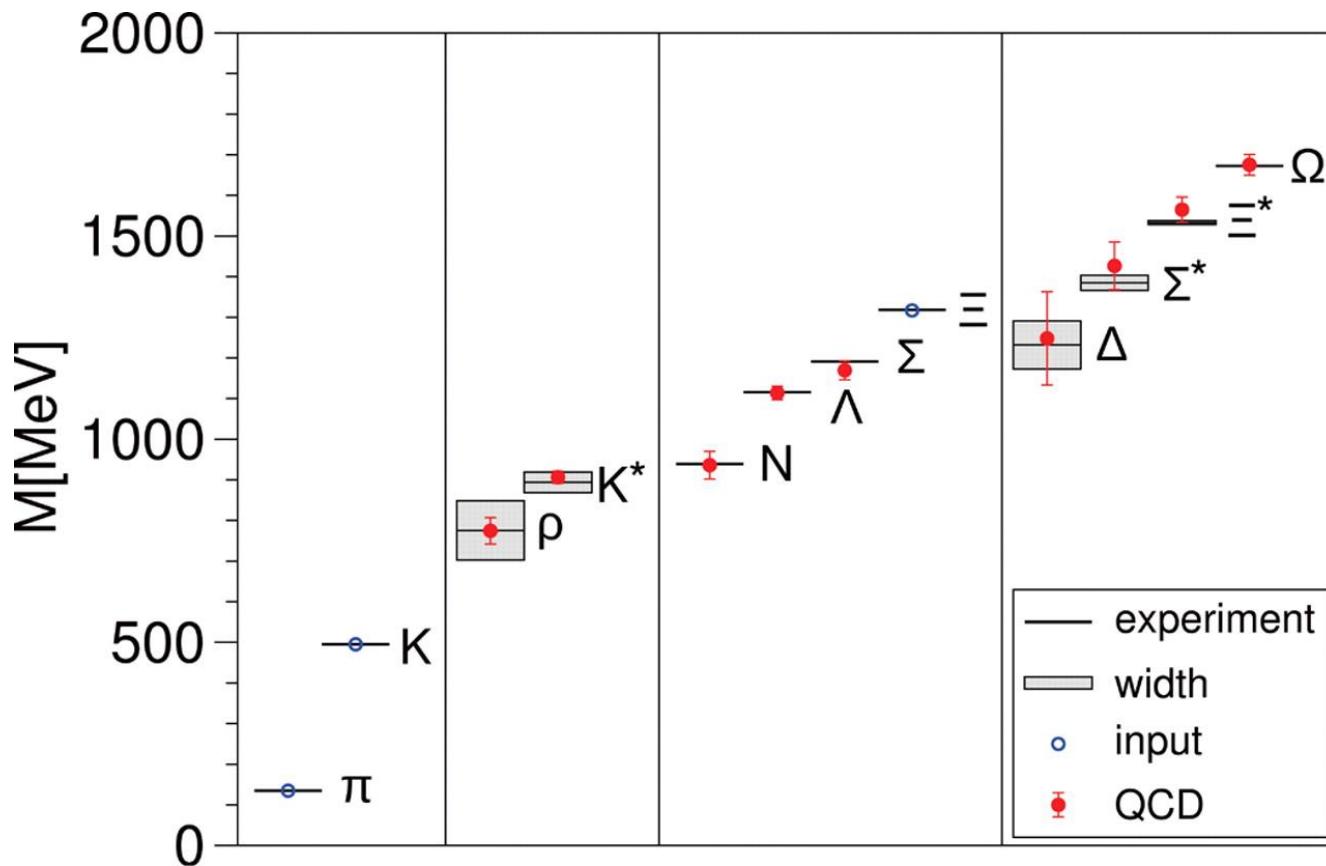
➤ Euclidean path integral with QCD action S_E

$$\langle \hat{O} \rangle = \frac{1}{Z_E} \int [\mathcal{D}\bar{q}] [\mathcal{D}q] [\mathcal{D}U] O e^{-S_E}$$

- dof: $\sim 100^4 \times 8 \sim 10^9$

➤ Monte Carlo simulation: gauge configuration with probability e^{-S_E}

Single hadron mass from LQCD

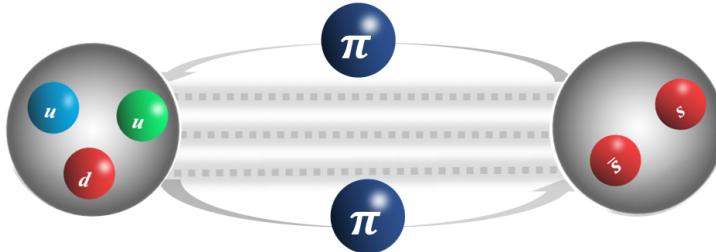


S. Durr *et al.* [BMW Coll.] Science 322, 1224 (2008)

- Successful for single hadron system

Lattice calculation of $N\phi$ interaction

- In this talk: First realistic LQCD results on $N\phi$ system
 - interaction potential from lattice QCD
 - interaction mechanism and the role of dynamical pion
 - scattering properties
 - detailed systematic error analysis



PHYSICAL REVIEW D **106**, 074507 (2022)

Attractive $N\phi$ interaction and two-pion tail from lattice QCD near physical point

Yan Lyu^{1,2,*}, Takumi Doi^{2,†}, Tetsuo Hatsuda^{2,‡}, Yoichi Ikeda^{3,§}, Jie Meng^{10,1,4,||}, Kenji Sasaki,^{3,¶} and Takuya Sugiura^{2,**}

Finite volume method

M. Lüscher, Nucl. Phys. B 354, 531 (1991)

- QM in 1d (PBC $\psi(0) = \psi(L) = 0$)

- Free theory

$$0 \quad \quad \quad L$$

$$\psi(x) \sim \sin(kx), \quad k = \frac{n\pi}{L}$$

- Interacting theory

$$0 \quad \quad \quad R \quad \quad \quad L$$

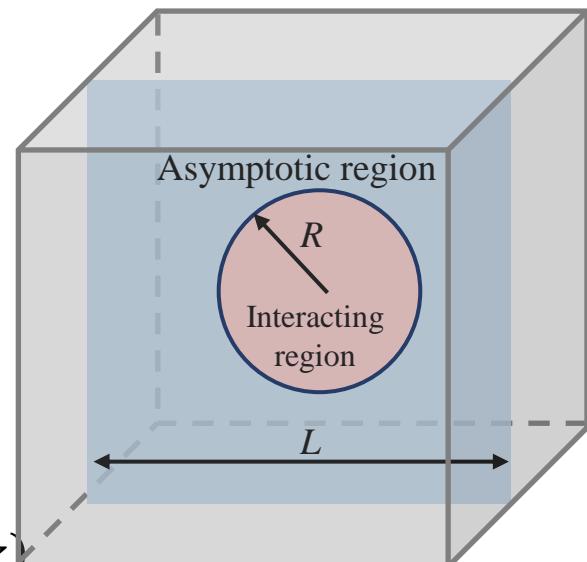
interacting region asymptotic region

$$\psi(x > R) \sim \sin(kx + \delta(k)), \quad k + \delta(k) = \frac{n\pi}{L}$$

- Lüscher's formula (3d box)

$$k \cot \delta_0(k) = \frac{2}{\sqrt{\pi}L} Z_{00}(1, q^2)$$

$$Z_{00}(s, q^2) = \frac{1}{4\sqrt{\pi}} \sum_{\vec{n} \in \mathbb{Z}^3} \frac{1}{(\vec{n}^2 - q^2)^s}, \quad q = \frac{kL}{2\pi}$$



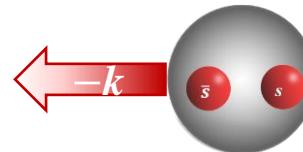
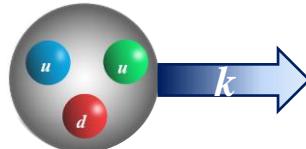
- FV momentum k from LQCD \rightarrow scattering phase shift $\delta(k)$

HAL QCD method

N. Ishii, S. Aoki and T. Hatsuda, Phys. Rev. Lett. **99**, 022001 (2007)

➤ Nambu-Bethe-Salpeter (NBS) amplitude

$$\psi^k(\mathbf{r}, t) = \langle 0 | N(\mathbf{r}, 0) \phi(\mathbf{0}, 0) | N(\mathbf{k}) \phi(-\mathbf{k}); E \rangle$$



$$N_\alpha = (u^T C \gamma_5 d) u_\alpha$$

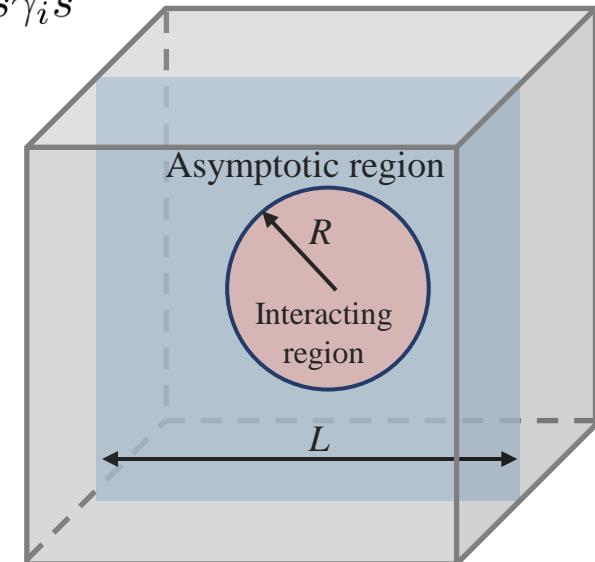
$$\phi_i = \bar{s} \gamma_i s$$

- asymptotic region

- ✓ Helmholtz eq. $(\nabla^2 + k^2)\psi(r) \simeq 0$
- ✓ Scattering phase shifts $\psi^k(r) \simeq A \frac{\sin(kr + \delta(k))}{kr}$

- “interacting” region

$$(\nabla^2 + k^2)\psi^k(\mathbf{r}) = m \times \int d^3 r' \mathbf{U}(\mathbf{r}, \mathbf{r}') \psi^k(\mathbf{r}')$$



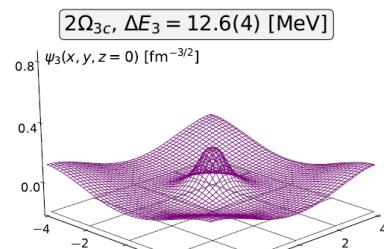
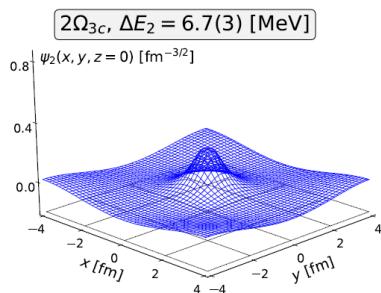
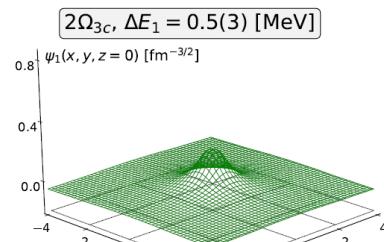
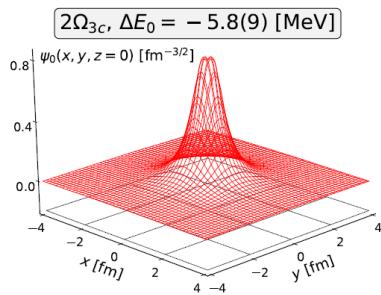
- ✓ non-local energy-independent kernel $U(\mathbf{r}, \mathbf{r}')$
- ✓ derivative expansion $U(\mathbf{r}, \mathbf{r}') = V(\mathbf{r}) \delta(\mathbf{r} - \mathbf{r}') + \sum V_i(\mathbf{r}) \nabla^i \delta(\mathbf{r} - \mathbf{r}')$

Check on convergence: Iritani *et al.*, JHEP 03, 007 (2019); YL *et al.*, PRD 105, 074512 (2022)

FV v.s. HAL

➤ If elastic states are available

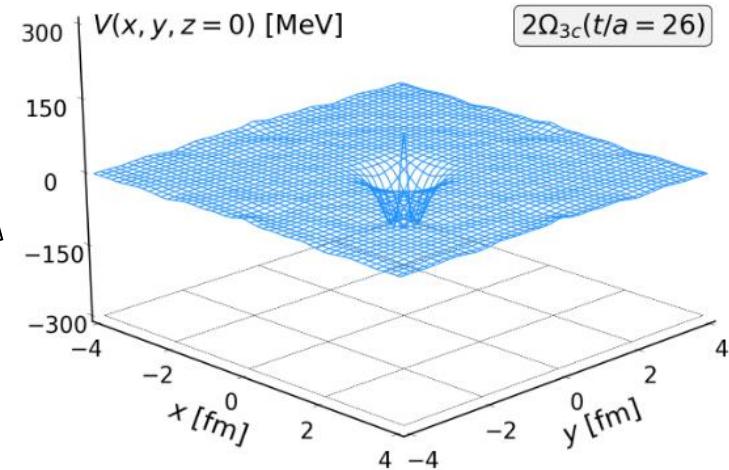
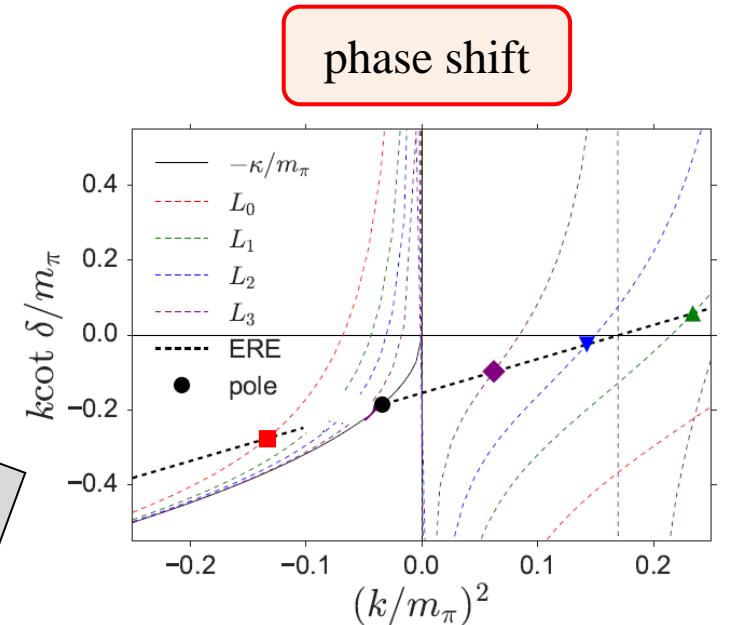
- FV: $\Delta E_{0,1,2,3\dots} \rightarrow$ phase shifts
- HAL: $\psi_{0,1,2,3\dots} \rightarrow$ potential



eigenstates

FV

HAL

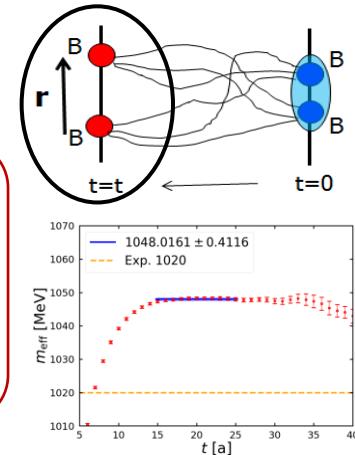


potential

Challenge in real calculation

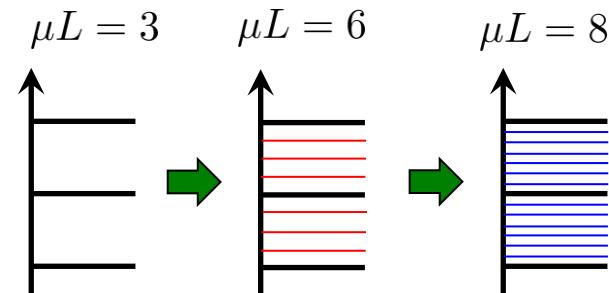
➤ Ground state saturation

$$\begin{aligned}
 R(\mathbf{r}, t) &= \langle \hat{O}(\mathbf{r}, t) \hat{O}^\dagger(0) \rangle = \sum_n \langle 0 | \hat{O}(\mathbf{r}, t) | n \rangle \langle n | \hat{O}^\dagger(0) | 0 \rangle \\
 &= \sum_n a_n \psi_n(\mathbf{r}) e^{-E_n t} \xrightarrow{t \gg \frac{1}{E_1 - E_0}} a_0 \psi_0(\mathbf{r}) e^{-E_0 t}
 \end{aligned}$$



➤ More and more dense states for large box and/or heavy system

$$t \gg \frac{1}{E_1 - E_0} \sim \mu \left(\frac{L}{2\pi} \right)^2$$



➤ Signal-to-noise problem for baryonic system

$$\frac{\text{Signal}}{\text{Noise}} = \frac{\langle N(t) \bar{N} \rangle}{\sqrt{\langle |N(t) \bar{N}|^2 \rangle}} = \frac{e^{-m_N t}}{e^{-3/2 m_\pi t}} = e^{-(m_N - 3/2 m_\pi)t}$$

- signal exponentially decrease as t increasing

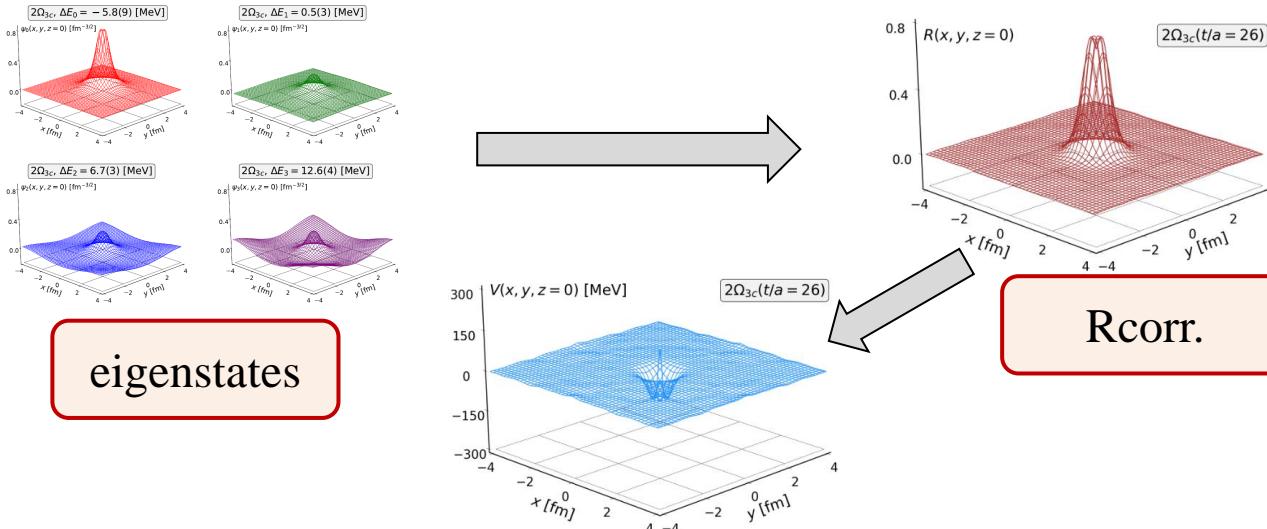
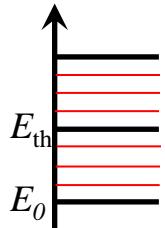
G. P. Lepage, in From Actions to Answers: Proceedings of the TASI 1989

Time-dependent HAL QCD method

N. Ishii, *et al.* (HAL QCD Collaboration), Phys. Lett. B **712**, 437 (2012)

- Define R-correlator as a summation over all elastic states

$$R(\mathbf{r}, t) = \sum_n a_n \psi_n(\mathbf{r}) e^{-E_n t} \xrightarrow{t \gg \frac{1}{E_{\text{th}} - E_0}} \sum_{n < N_{\text{th}}} a_n \psi_n(\mathbf{r}) e^{-E_n t}$$



- Master equation $\sim k^2/(2\mu)$

potential

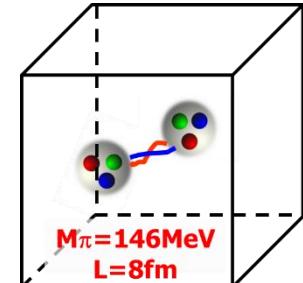
$$\left(\frac{1 + 3\delta^2}{8\mu} \frac{\partial^2}{\partial t^2} - \frac{\partial}{\partial t} + \frac{\nabla^2}{2\mu} \right) R(\mathbf{r}, t) = \int d\mathbf{r}' \mathbf{U}(\mathbf{r}, \mathbf{r}') R(\mathbf{r}', t)$$

- $\delta = \frac{m_N - m_\phi}{m_N + m_\phi}; \mu = \frac{m_N m_\phi}{m_N + m_\phi}$

Lattice setup

- (2+1)-flavor configuration near the physical point
 - Iwasaki gauge action and $O(a)$ -improved Wilson quark action

$L \times T$	a [fm]	La [fm]	m_π [MeV]	m_K [MeV]
$96^3 \times 96$	0.0846(7)	8.1	146	525



K.-I. Ishikawa, *et al.* (PACS Collaboration), *Proc. Sci.*, LATTICE2015 (2016) 075

- Relevant hadron masses

Hadron	Lattice Mass [MeV]	Physical Mass[MeV]
N	954.0(2.9)	938.9
ϕ	1048.0(4)	1019.5

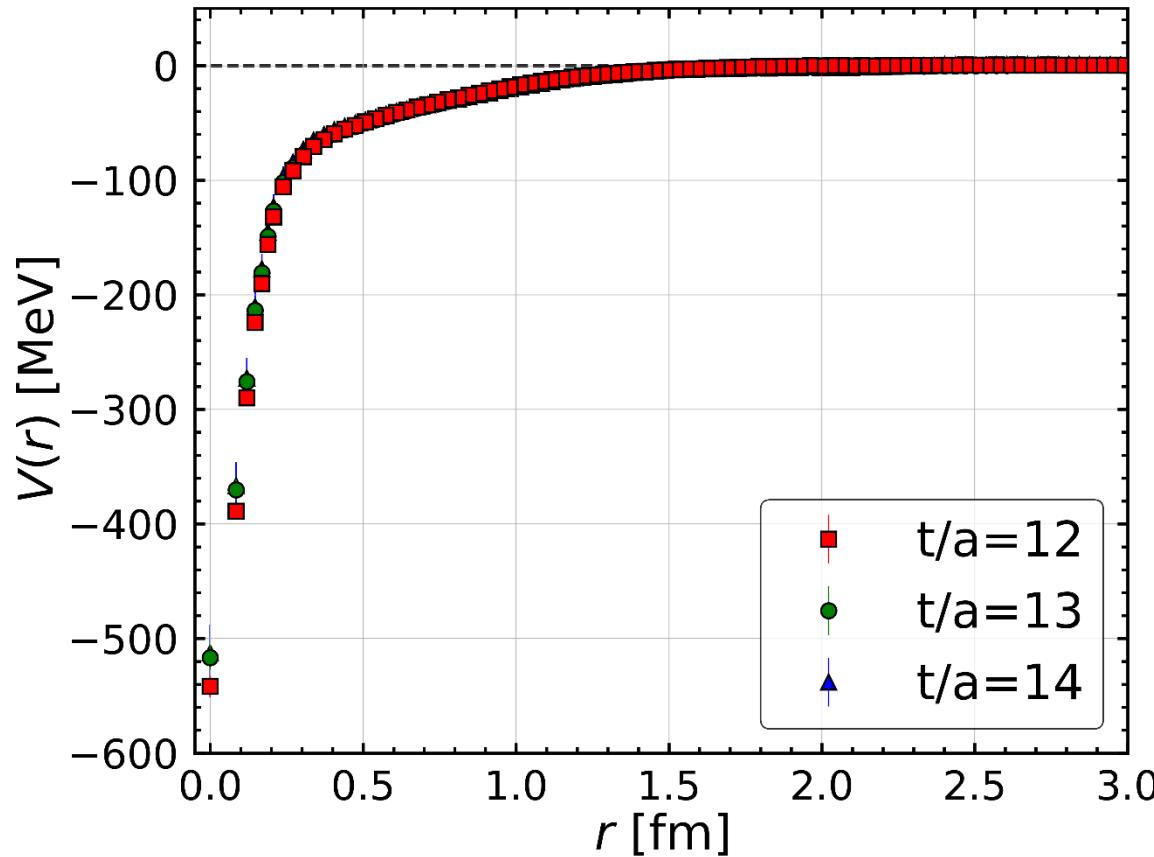
- Note that $\phi \rightarrow K\bar{K}$ is prohibited due to $m_\phi < 2m_K$
- Statistics

$$200_{\text{conf.}} \times 80_{\text{src}} \times 4_{\text{rot.}} \times 2_{\text{b.f.}} = 128000$$



Potential

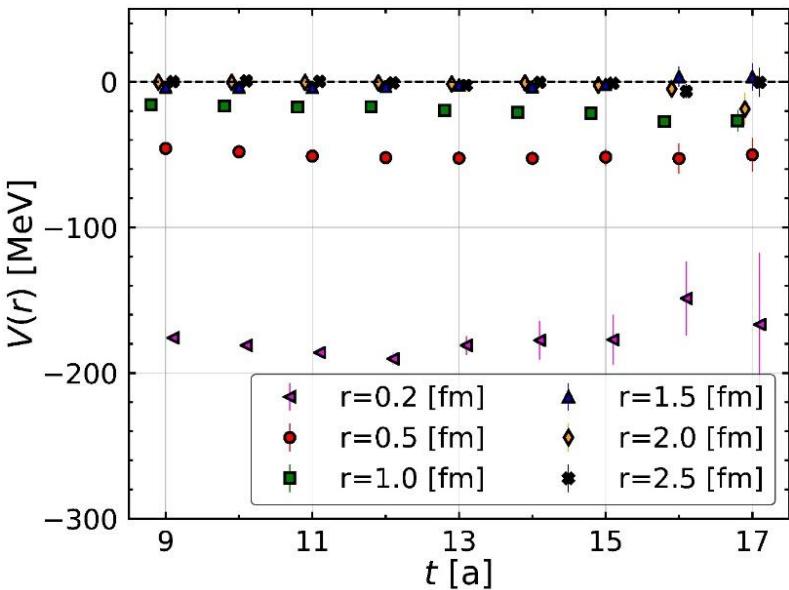
- $N\phi(^4S_{3/2})$ potential at Euclidean time 12, 13 and 14



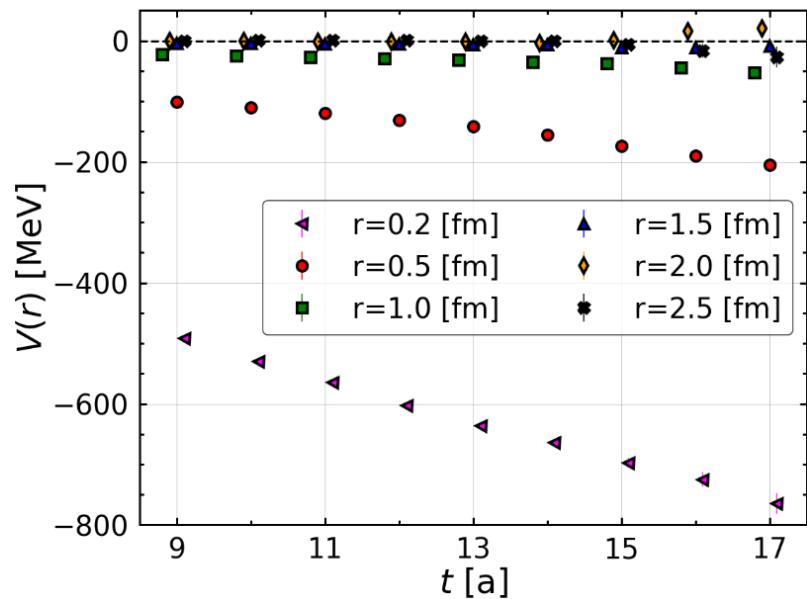
- attractive core, Pauli exclusion does not operate due to no common quarks
- long-ranged attractive tail, hints of pion dynamics
- weak t dependence

Potential in wider range of t

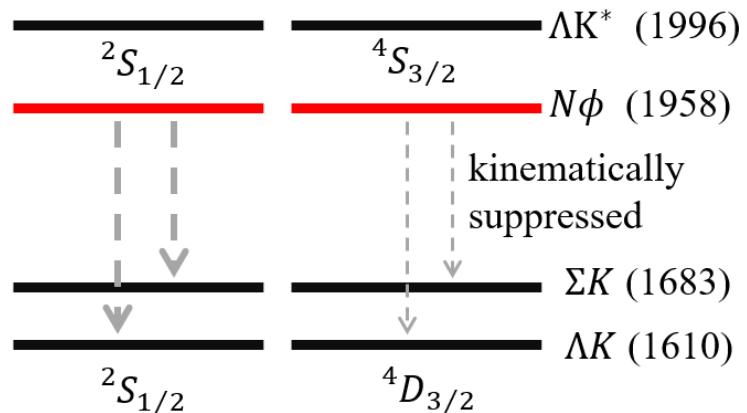
➤ $^4S_{3/2}$



➤ $^2S_{1/2}$

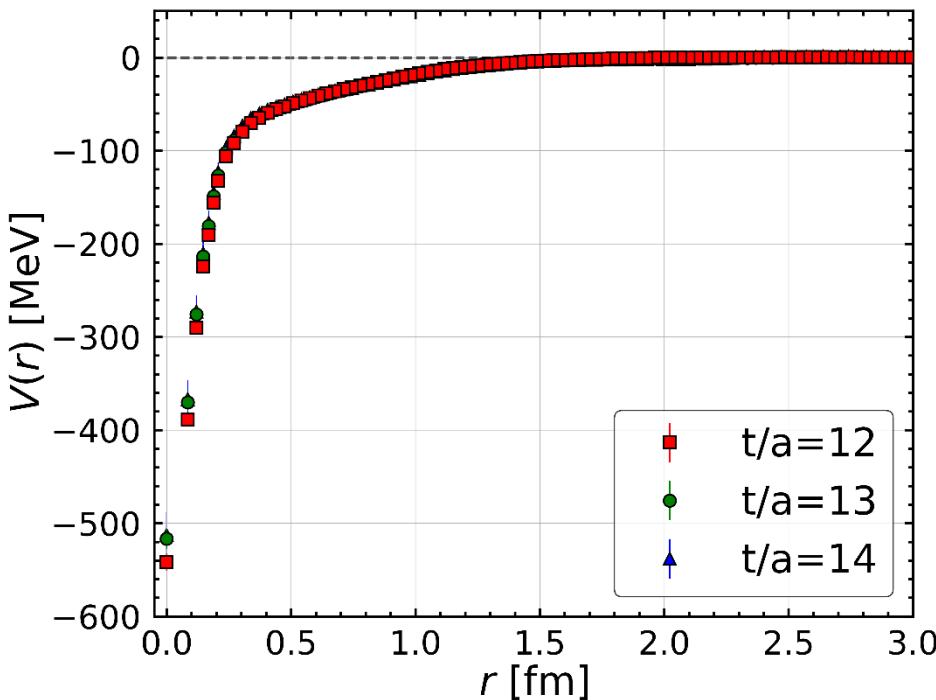


- correlation function is dominated by the elastic states in $^4S_{3/2}$
- this is in sharp contrast to $^2S_{1/2}$ case, where potential show clear t -dependence

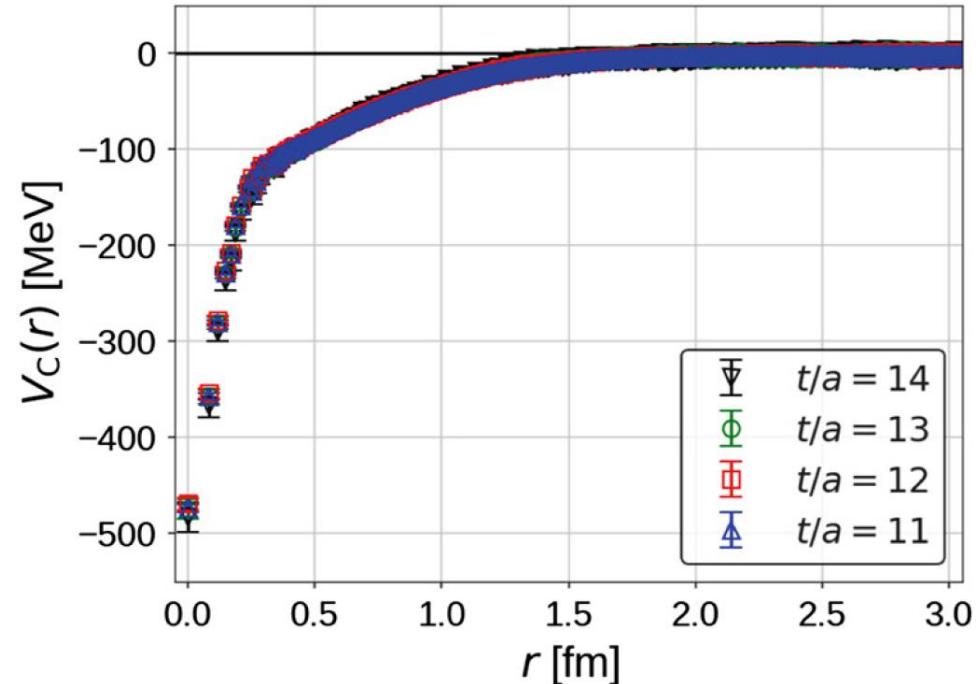


$N\phi$ v.s. $N\Omega$

➤ $N\phi(^4S_{3/2})$ potential



➤ $N\Omega(^5S_2)$ potential



T. Iritani, et al. (HAL QCD Coll.), Phys. Lett. B 792, 284 (2019)

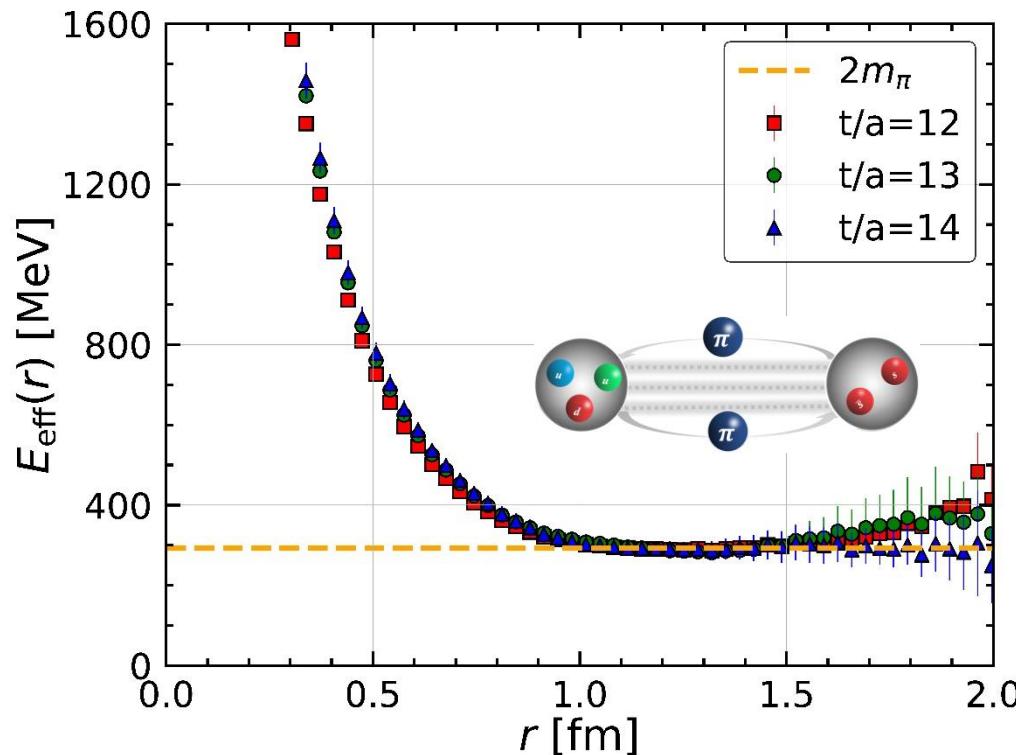
- universality of pion dynamics (?)

Two-pion-exchange tail

- Interaction between nucleon and heavy quarkonium at large distances

J. Castella and G. Krein, Phys. Rev. D **98**, 014029 (2018); H. Fujii and D. Kharzeev, Phys. Rev. D **60**, 114039 (1999)
G. Bhanot and M. E. Peskin, Nucl. Phys. B **156**, 391(1979)

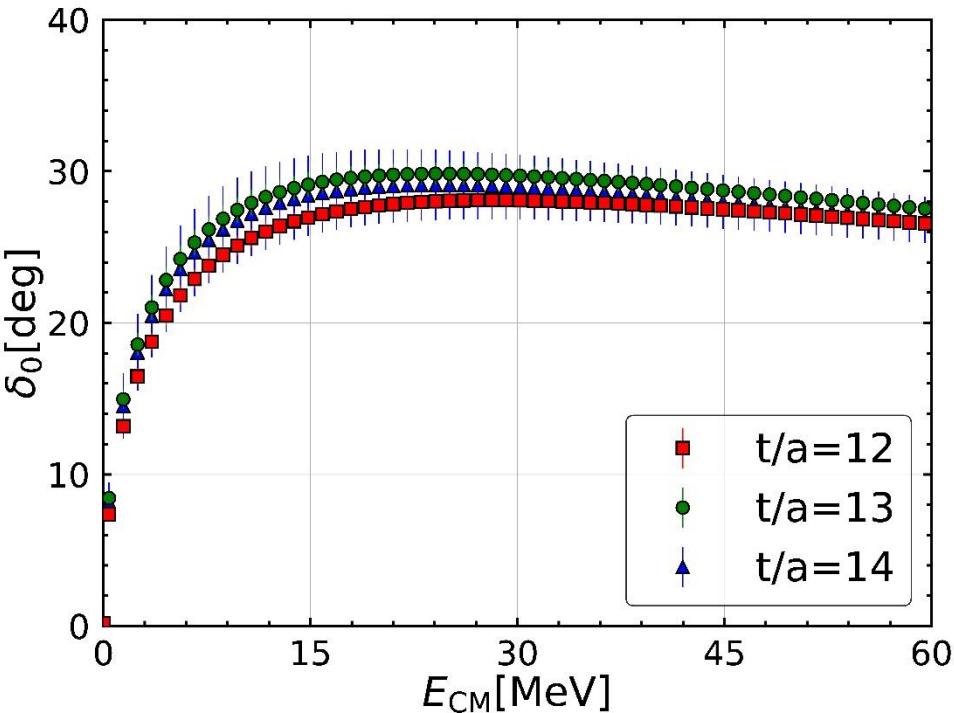
$$V(r) = -\alpha \frac{e^{-2m_\pi r}}{r^2} \quad \longrightarrow \quad E_{\text{eff}}(r) = -\frac{\ln[-V(r)r^2/\alpha]}{r}$$



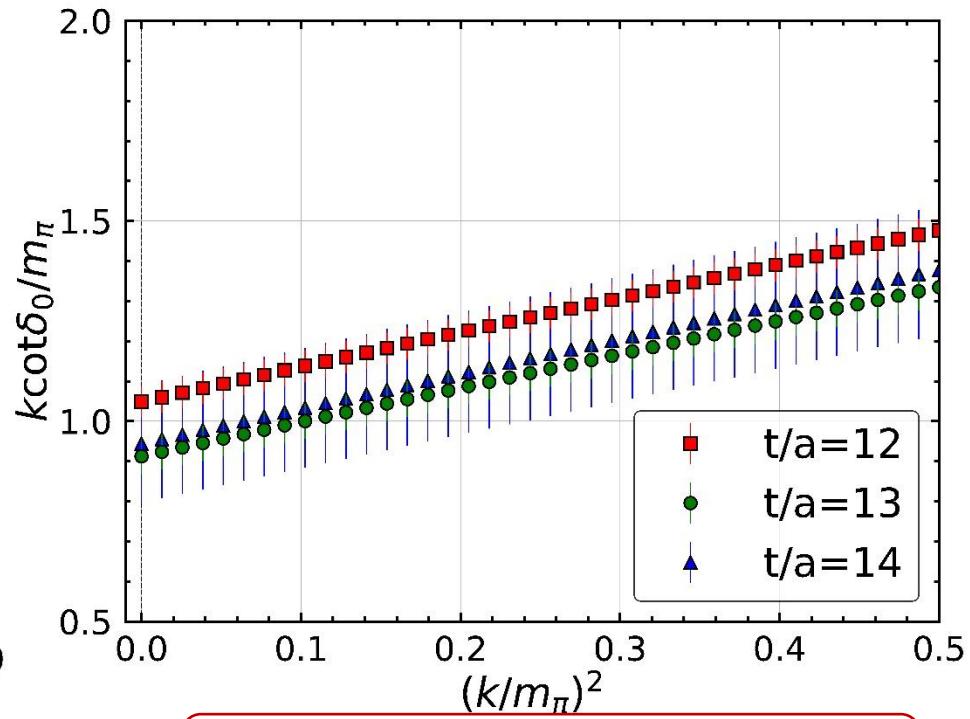
- long-range ($r > 1$ fm) potential is indeed dominated by the TPE

Physical observables

➤ Scattering phase shifts



➤ Effective range expansion



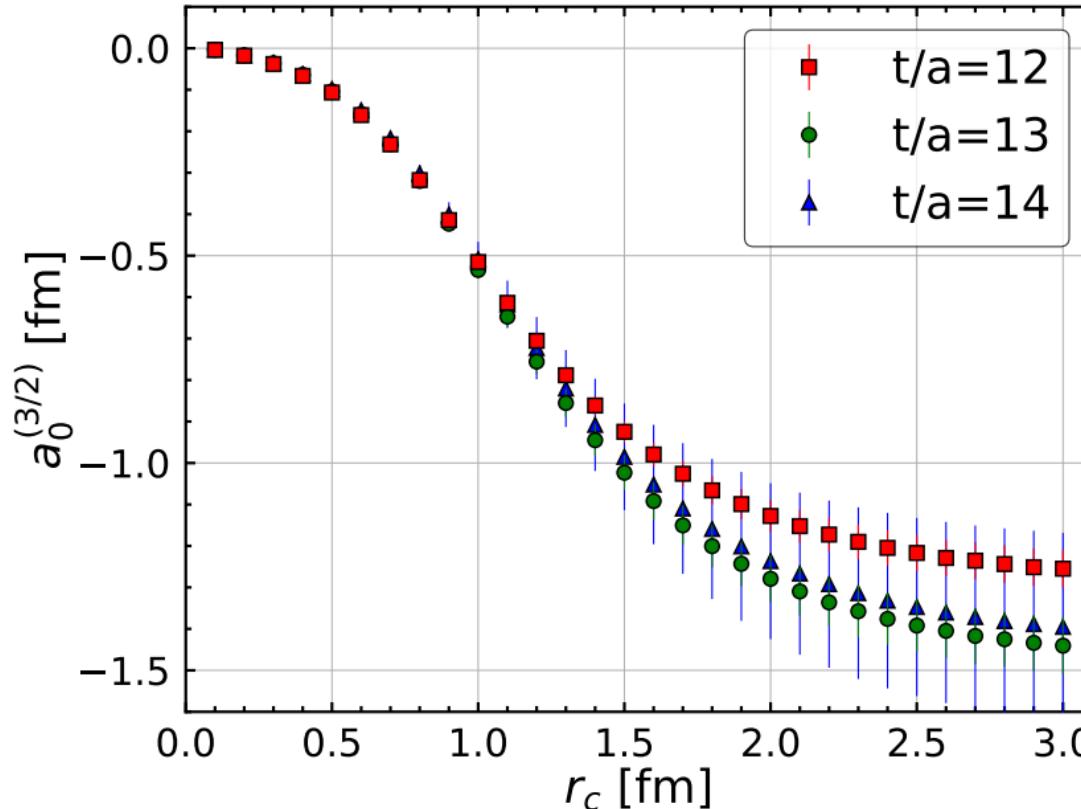
$$k \cot \delta_0 = -\frac{1}{a_0} + \frac{1}{2} r_{\text{eff}} k^2 + O(k^2)$$

	m_π [MeV]	$a_0^{(3/2)}$ [fm]	$r_{\text{eff}}^{(3/2)}$ [fm]
consistent w/ ALICE Exp.	146.4	$-1.43(23)_{\text{stat.}} \left({}^{+36}_{-06} \right)_{\text{syst.}}$	$2.36(10)_{\text{stat.}} \left({}^{+02}_{-48} \right)_{\text{syst.}}$

Long-range potential's contribution to scattering length

- Scattering length from potential with long-range cutoff

$$V(r; r_c) = \begin{cases} V(r), & r \leq r_c \\ 0, & r > r_c \end{cases}$$



- long-range potential significantly contributes to scattering length

Other systematic errors

- Slightly heavy quark masses
 - TPE becomes slightly long-ranged with weaker strength $m_\pi(\text{Lat.}) \rightarrow m_\pi(\text{Phy.})$
 - hadron masses decrease $m_{\phi,N}(\text{Lat.}) \rightarrow m_{\phi,N}(\text{Phy.})$
 - towards physical point
physical $m_{\phi,N}$, change m_π to 138.0 MeV for TPE tail

m_π [MeV]	$a_0^{(3/2)}$ [fm]	$r_{\text{eff}}^{(3/2)}$ [fm]
146.4	$-1.43(23)_{\text{stat.}} \left({}^{+36}_{-06} \right)_{\text{syst.}}$	$2.36(10)_{\text{stat.}} \left({}^{+02}_{-48} \right)_{\text{syst.}}$
138.0	$\simeq -1.25$	$\simeq 2.49$

- Finite volume
 - expected to be as small as $\exp(-2m_\pi L/2) \simeq 0.3\%$
- Finite cutoff
 - non-perturbative $O(a)$ -improvement for u, d, s quark, $O\left((a\Lambda_{\text{QCD}})^2\right) \sim O(1)\%$
 - even completely cut $V(r)$ at $r < 0.1$ fm $\rightarrow 2\%$

Summary and outlook

- **Summary:** $N\phi(^4S_{3/2})$ interaction is studied from LQCD for the first time w/ nearly physical quark masses and large volume
 - Extract potential from spacetime correlation function
 - Potential: attractive; TPE tail
 - Scattering properties
- **Outlook:**
 - $N\phi(^2S_{1/2})$ interaction from a coupled-channel analysis $\Lambda K - \Sigma K - p\phi$
 - Physical point simulations
 - $N - \bar{c}c$ interaction

Thanks for your attention

